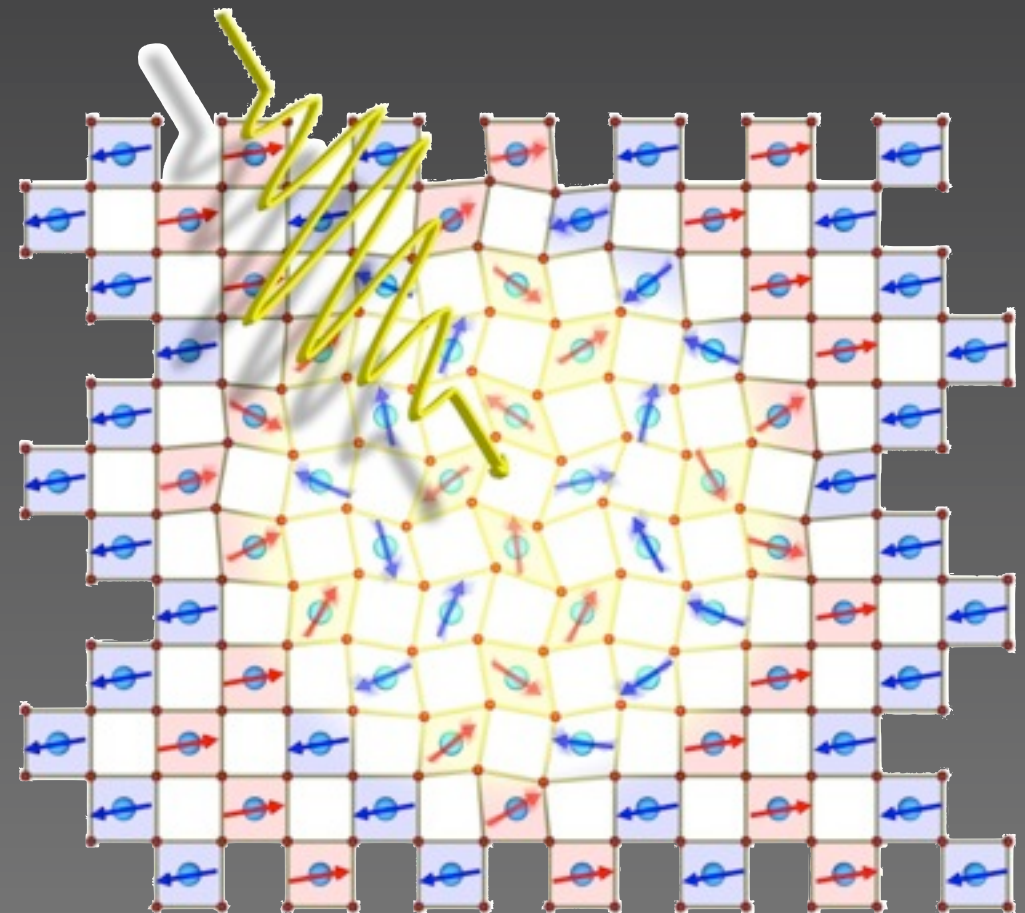
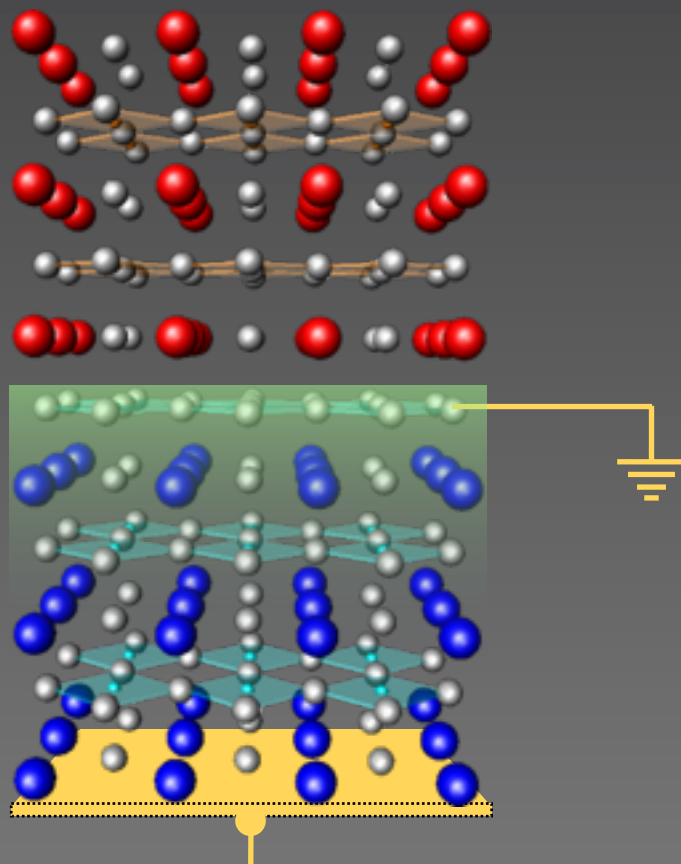


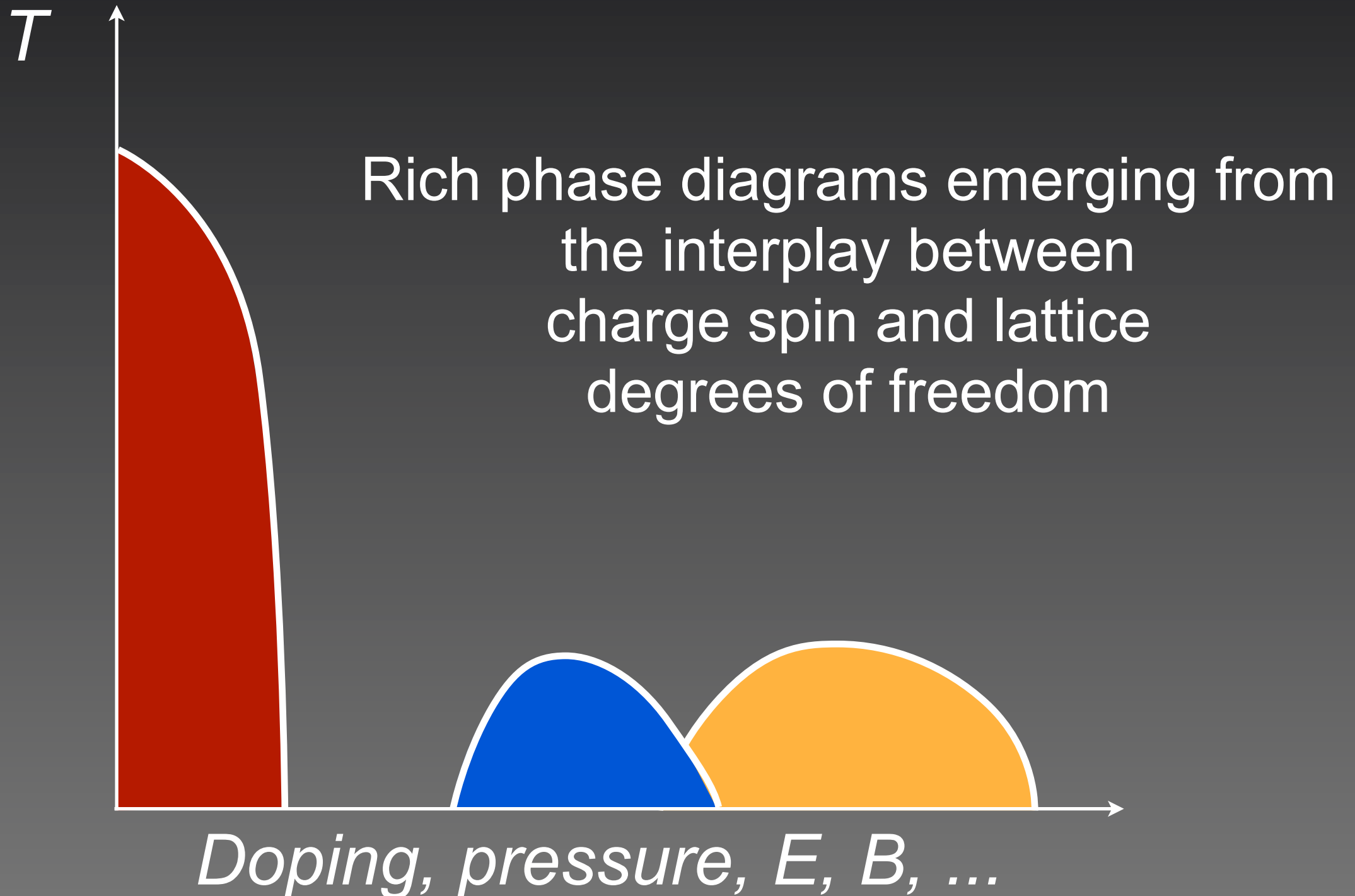
Electronic phase control in complex oxides heterostructures

Andrea Caviglia

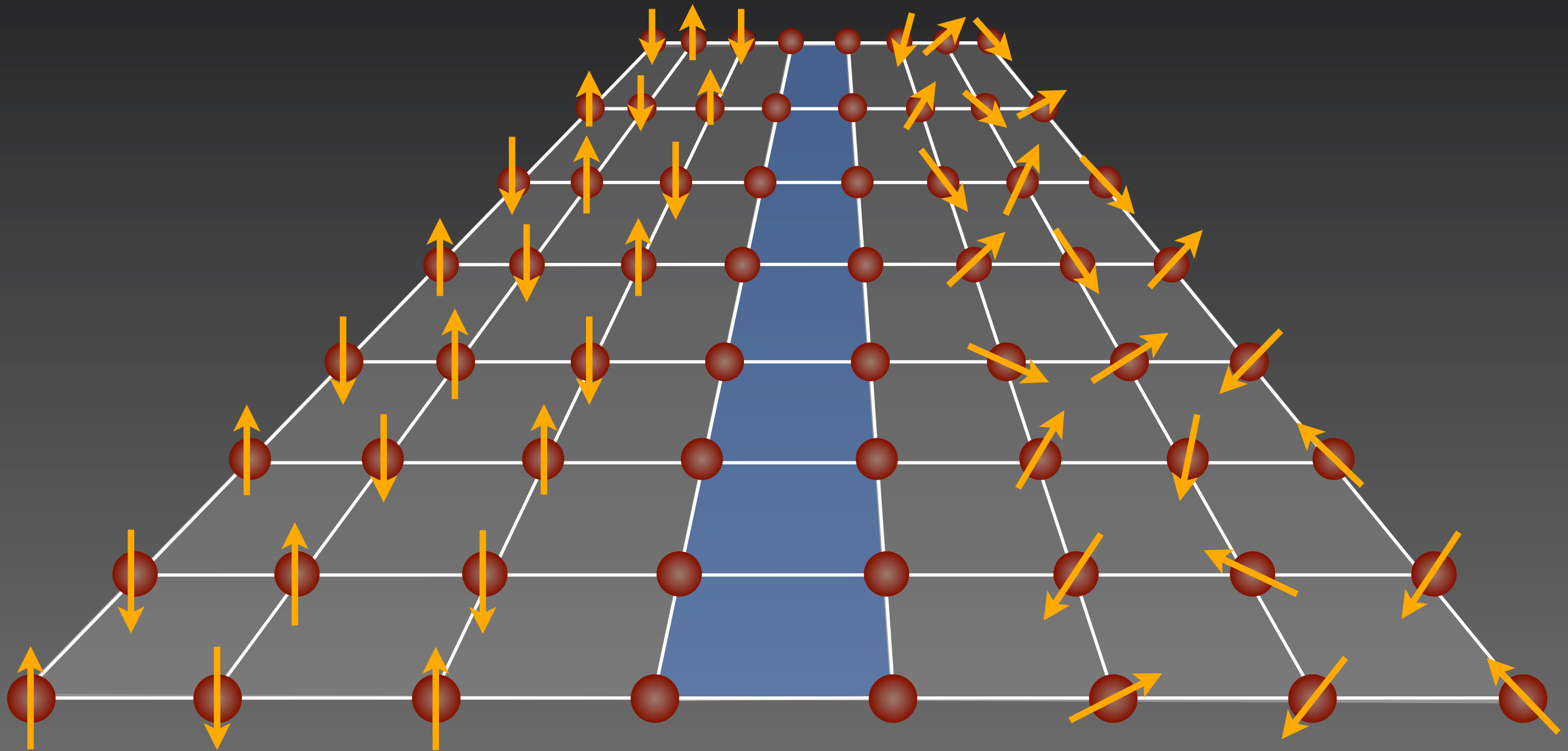
Max Planck Department for Structural Dynamics
Center for Free Electron Laser Science
University of Hamburg



Competing ground states in quantum materials



Novel electronic phases at their interfaces



Phase control at interfaces

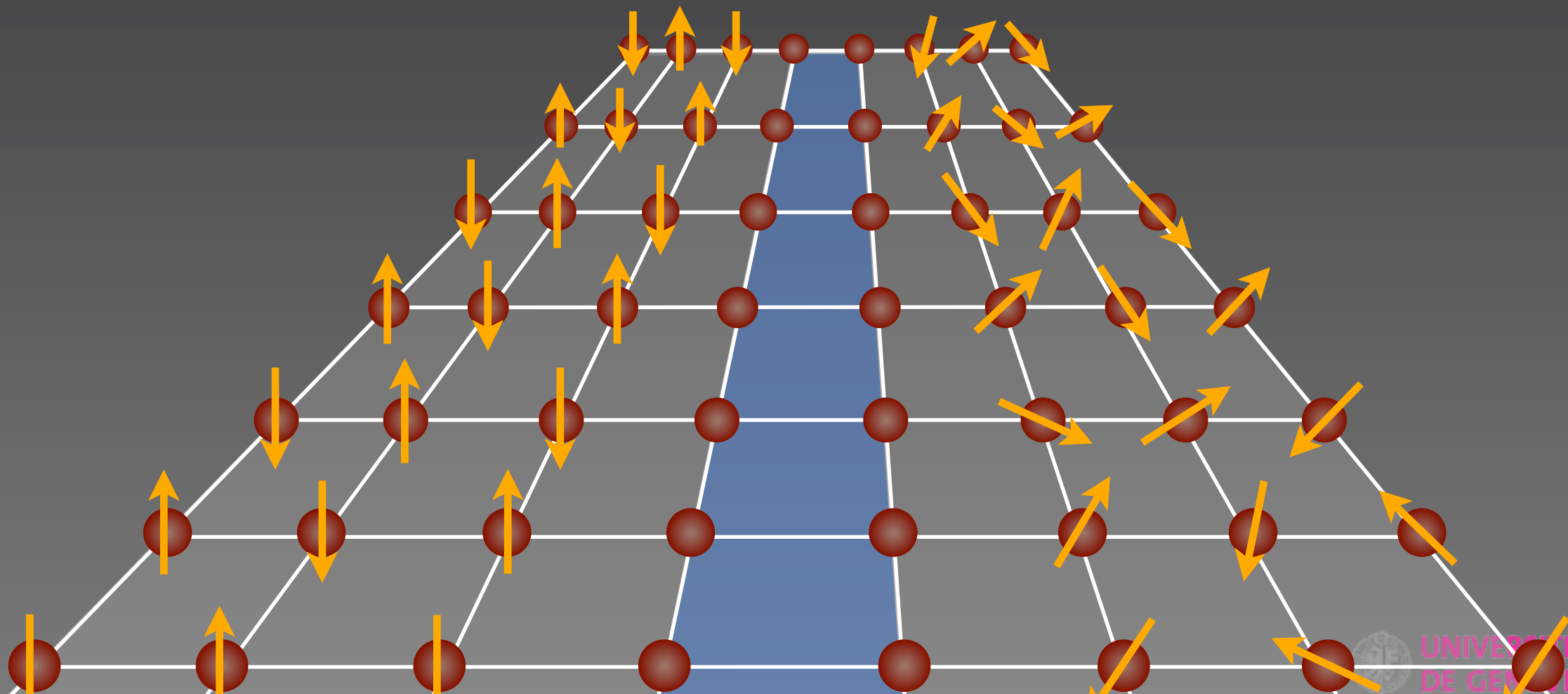
Breaking of inversion symmetry

Low dimensionality

Structural distortions

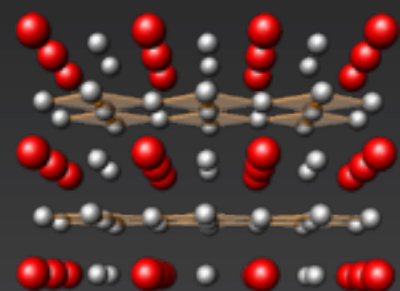
Modulation doping

Polar discontinuities

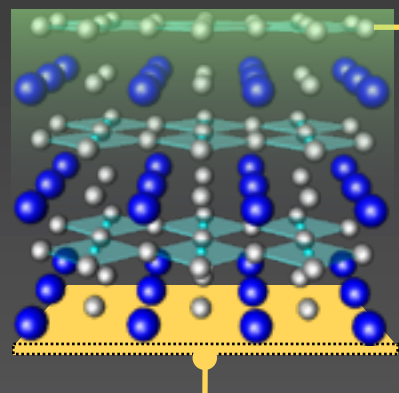


Two experimental approaches

Electrostatic field-effect



static 10 kV/cm fields
 $n_{2D} \sim 10^{13} - 10^{14} \text{ cm}^{-2}$
modulation of doping
and interactions



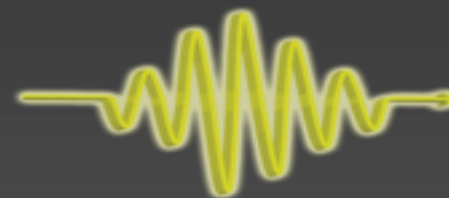
electronic properties
investigated by
dc transport

$\text{LaAlO}_3/\text{SrTiO}_3$

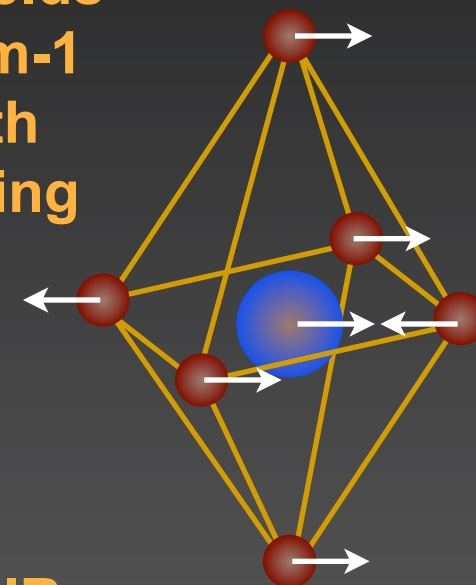
Superconductor - Insulator
Quantum Phase Transition

Ultrafast lattice excitation

100 fs pulsed 10 MV/cm fields
16 μm wavelength, 600 cm^{-1}
modulation of bandwidth
non-linear phonon coupling



electronic properties
investigated by
pump and probe THz and IR
spectroscopy



$\text{NdNiO}_3/\text{LaAlO}_3$

Insulator - Metal
Non-equilibrium Phase Transition

Outline

$\text{LaAlO}_3/\text{SrTiO}_3$

Two-dimensional superconductivity

Superconductor to insulator quantum phase transition

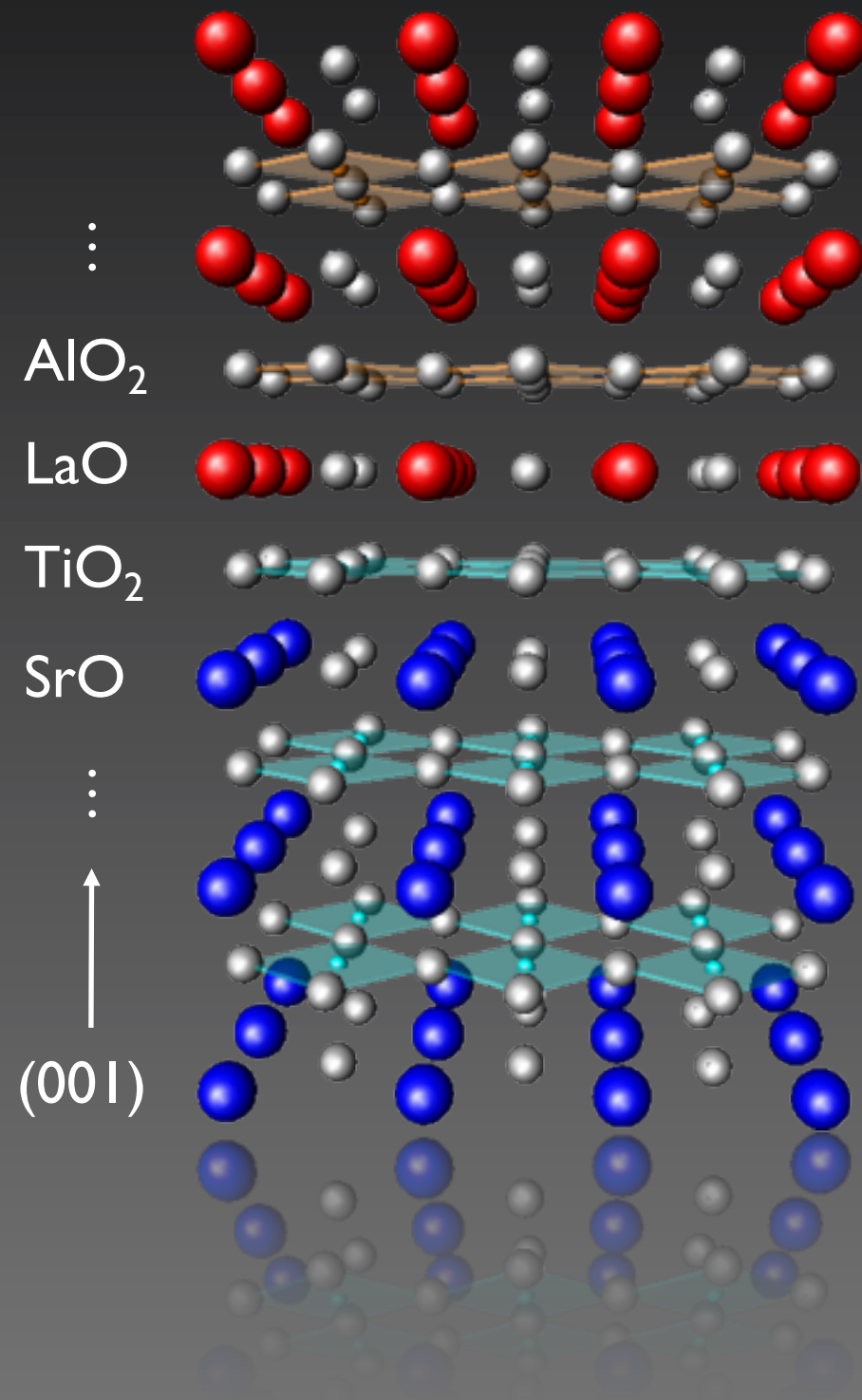
Spin-orbit interaction

$\text{NdNiO}_3/\text{LaAlO}_3$

Metal Insulator transition

Pump and probe spectroscopy

Vibrational excitation



LaAlO_3 :

band insulator

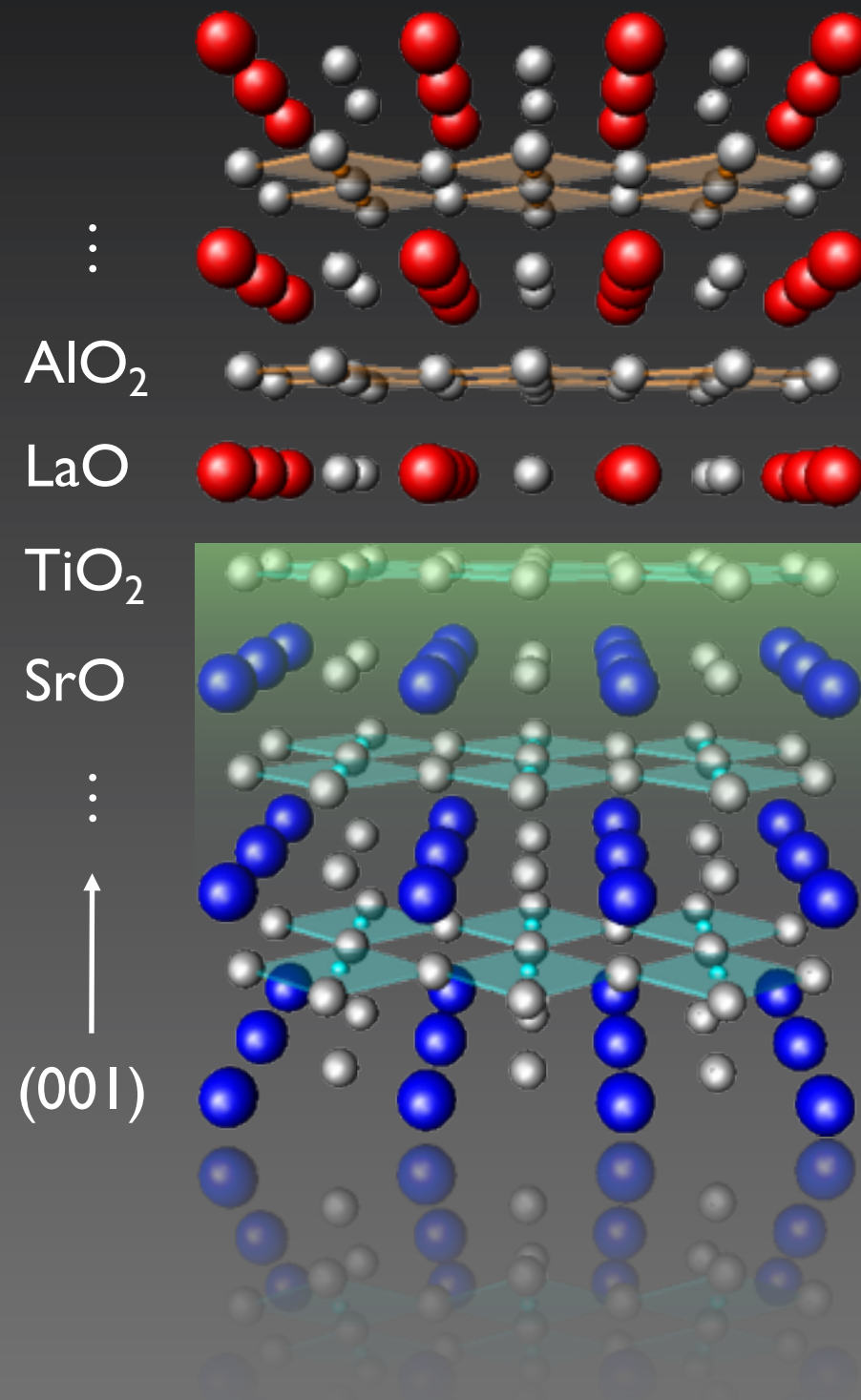
$$\Delta = 5.6 \text{ eV}, \quad \kappa = 24$$

SrTiO_3 :

band insulator

$$\Delta = 3.2 \text{ eV}, \quad \kappa(300 \text{ K}) = 300$$

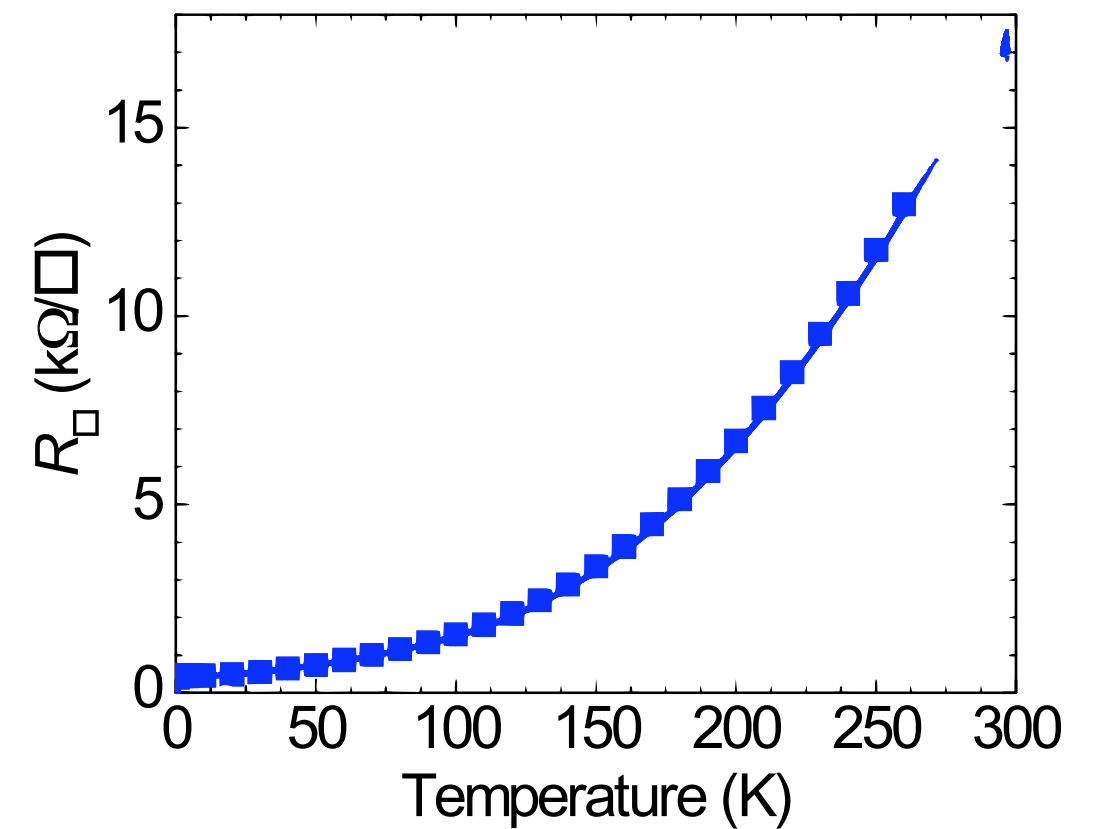
quantum paraelectric

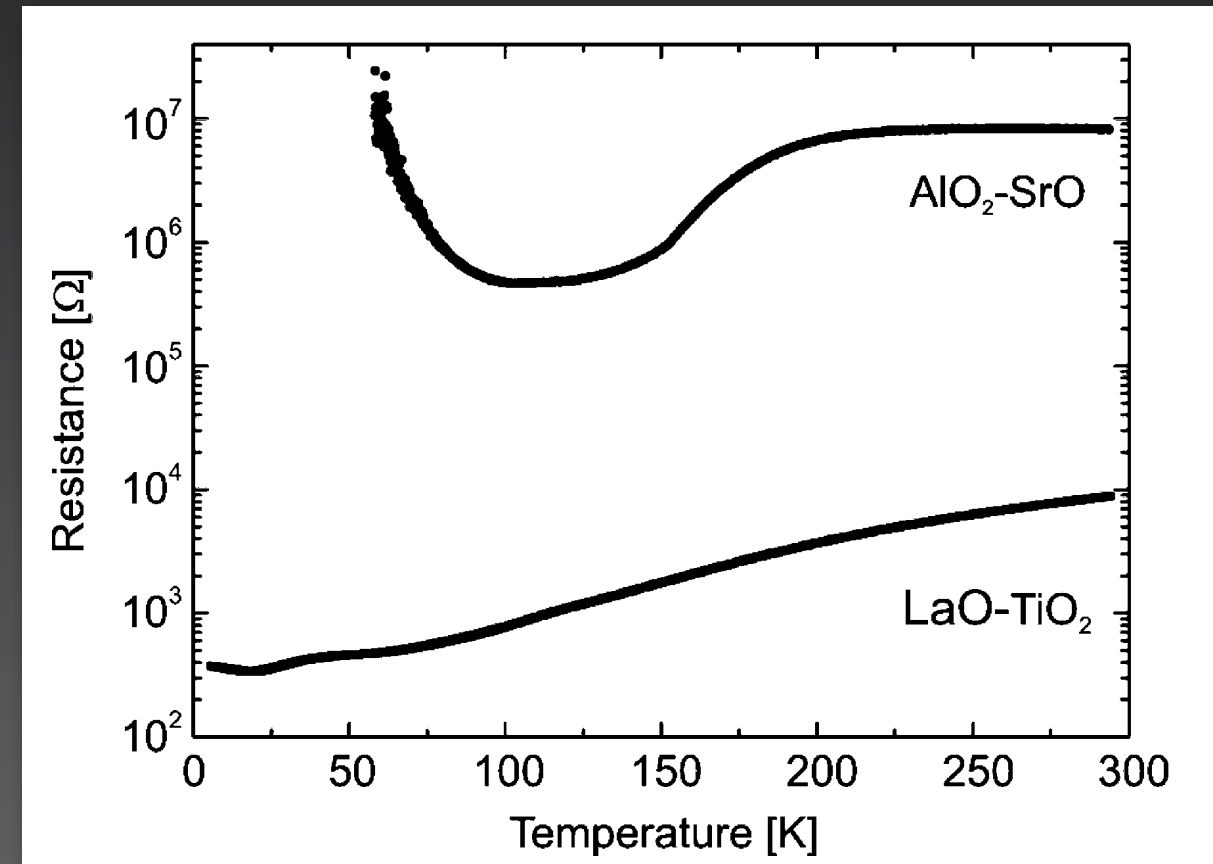
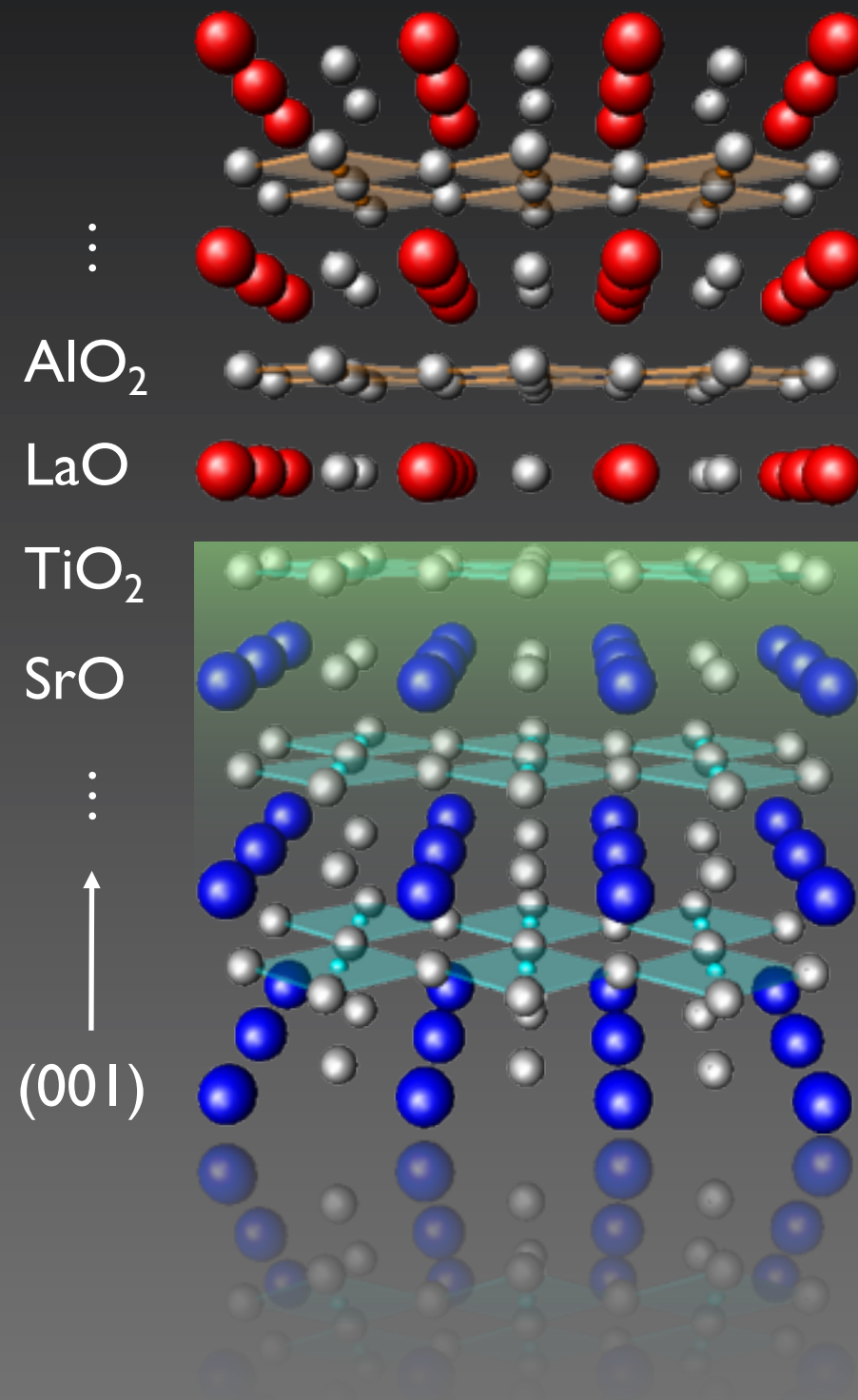


A high-mobility electron gas at the $\text{LaAlO}_3/\text{SrTiO}_3$ heterointerface

A. Ohtomo^{1,2,3} & H. Y. Hwang^{1,3,4}

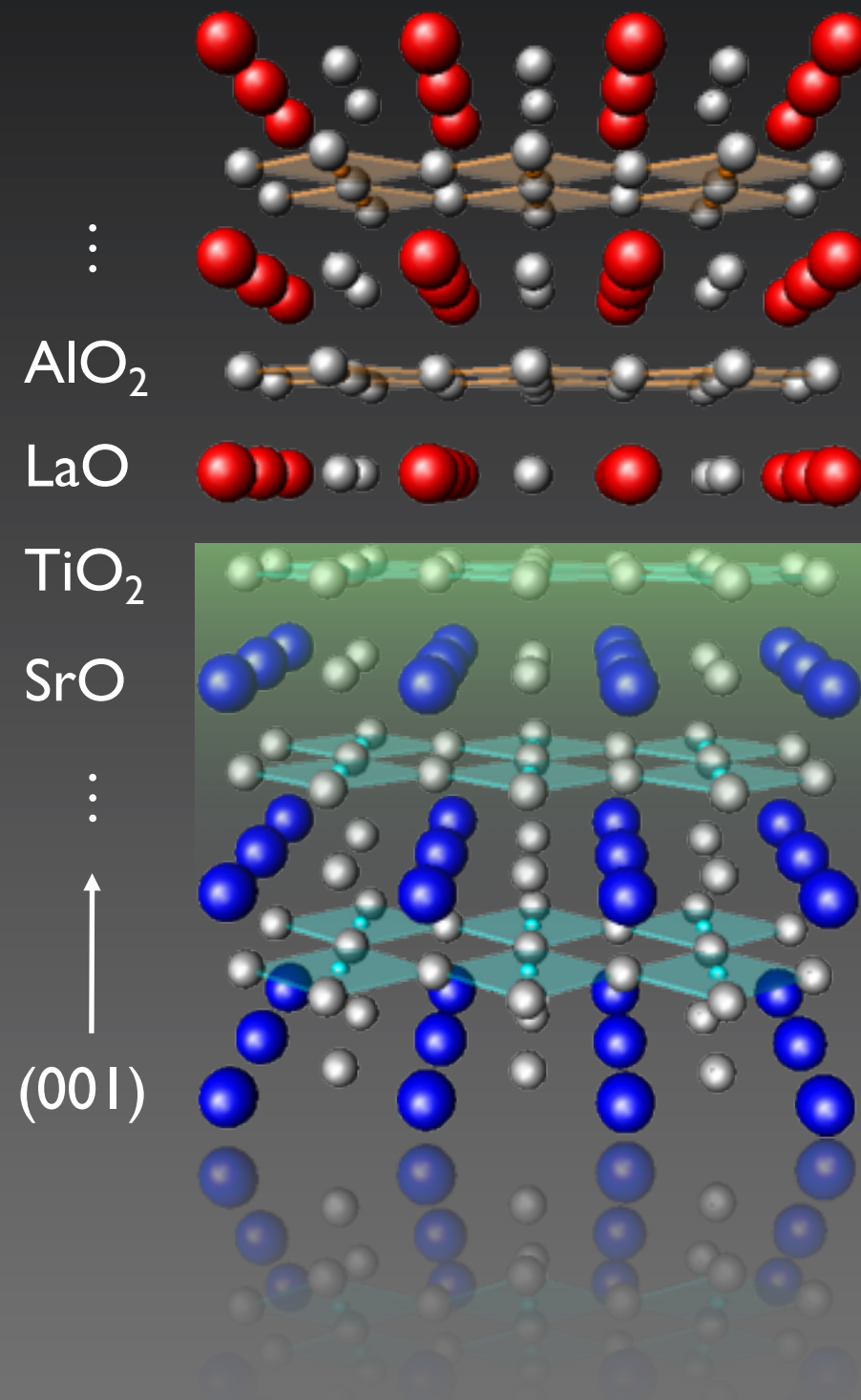
Nature **427**, 423 (2004)





A. Ohtomo and H.Y. Hwang *Nature* **427**, 423 (2004)
 M. Huijben *et al.* *Advanced materials* **21**, 1665 (2009)

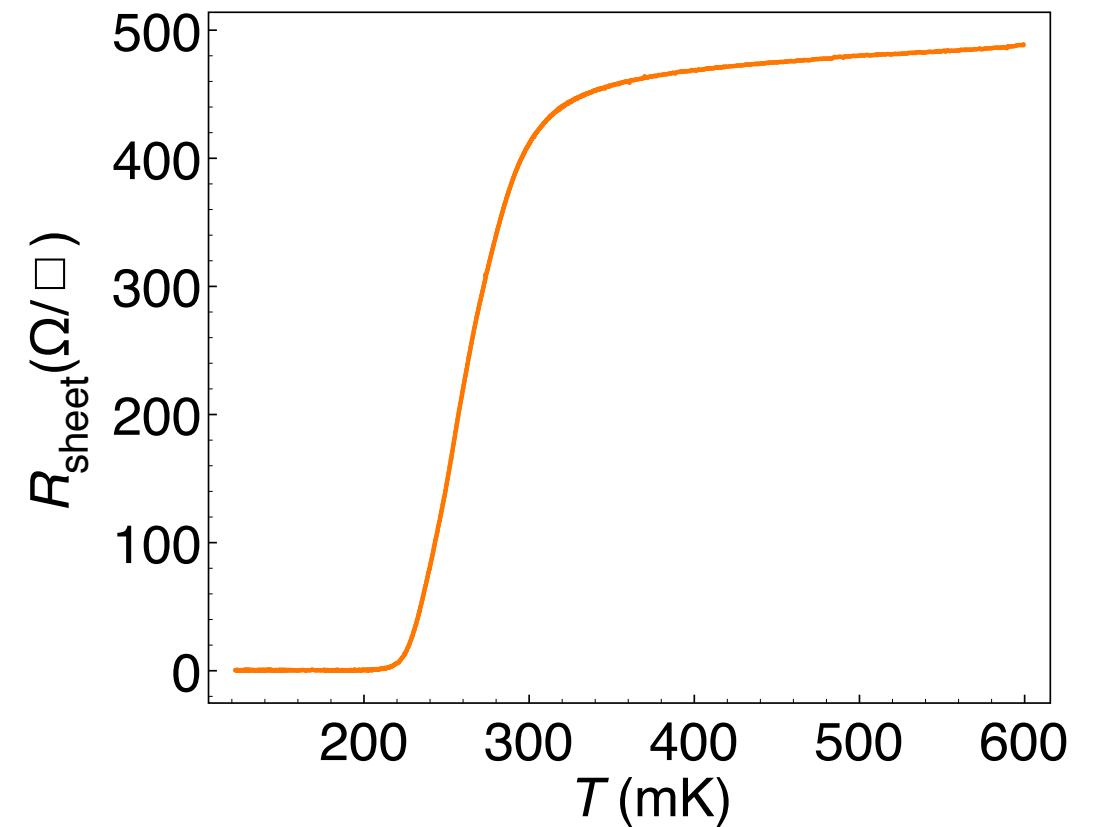
SC in LAO/STO



Superconducting Interfaces Between Insulating Oxides

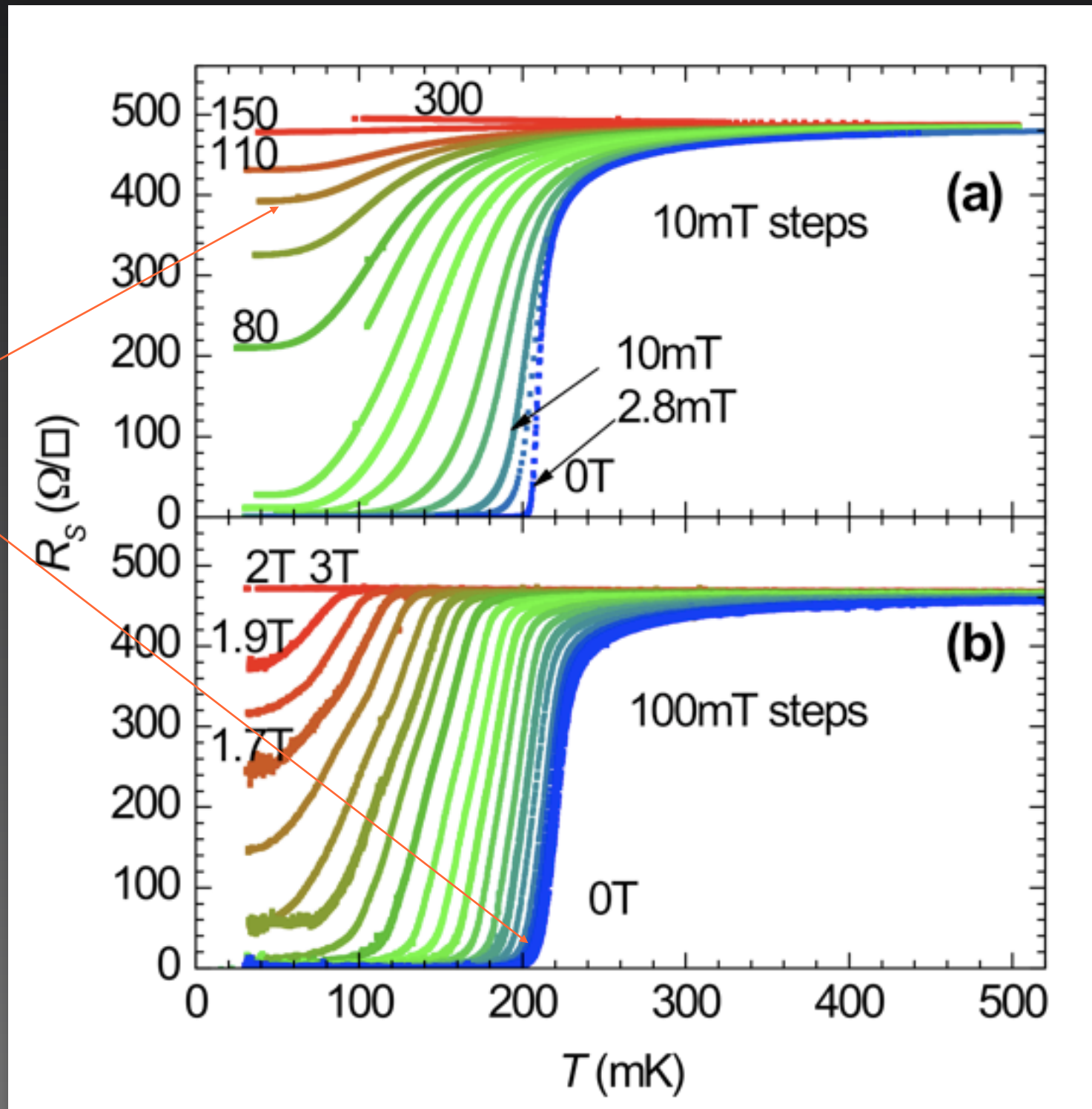
N. Reyren,¹ S. Thiel,² A. D. Caviglia,¹ L. Fitting Kourkoutis,³ G. Hammerl,² C. Richter,²
 C. W. Schneider,² T. Kopp,² A.-S. Rüetschi,¹ D. Jaccard,¹ M. Gabay,⁴ D. A. Müller,³
 J.-M. Triscone,¹ J. Mannhart^{2*}

Science **317**, 1196 (2007)



Transport anisotropy

100mT

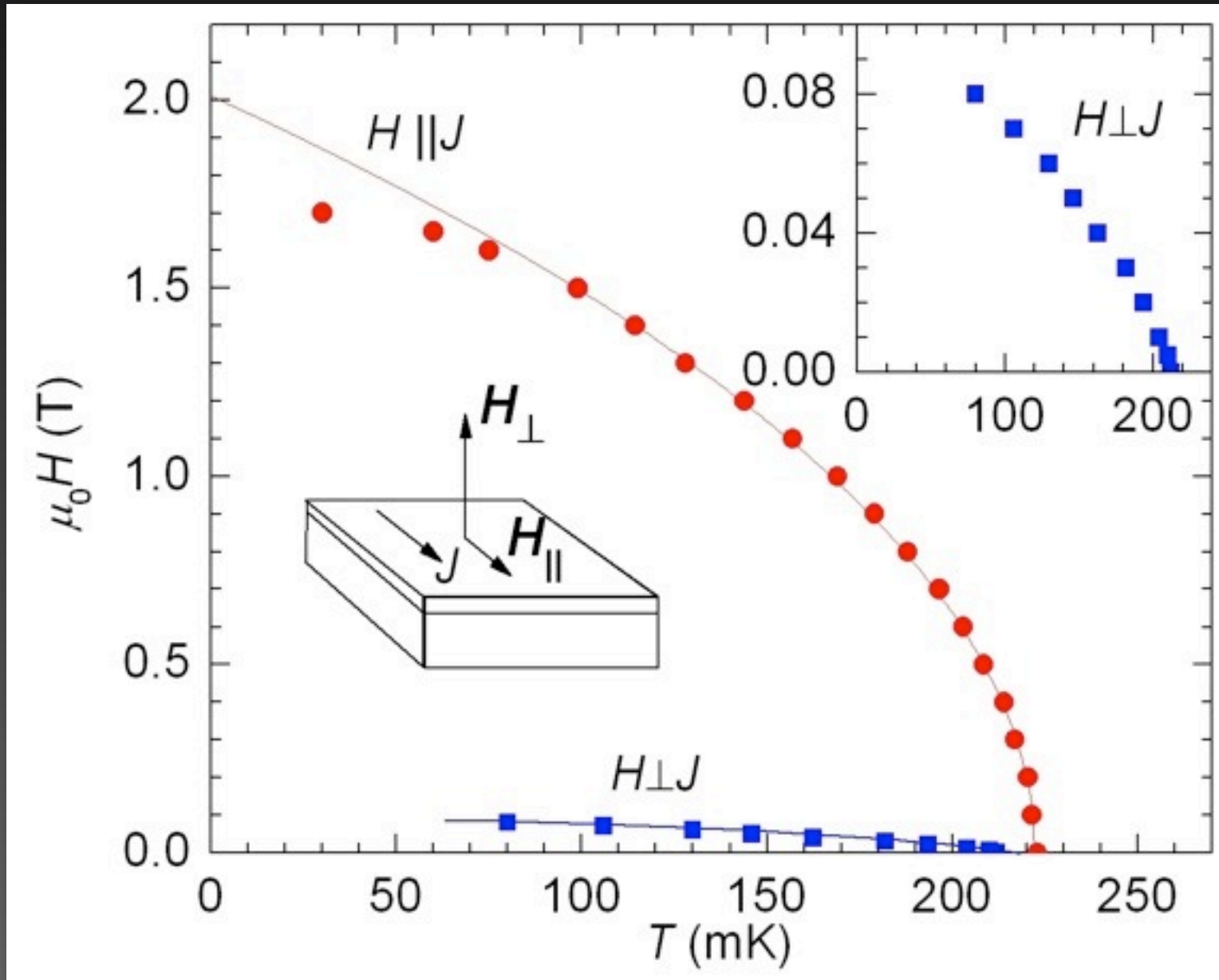


$H \perp$

$H \parallel$

N. Reyren, AC *et al*, *APL* 94, 112506 (2009)

Two-dimensional SC



$$\mu_0 H_{\perp}(T) = \frac{\Phi_0}{2\pi\xi_{\parallel}^2(T)}$$

$$\xi_{\parallel}(T=0) \sim 60 \text{ nm}$$

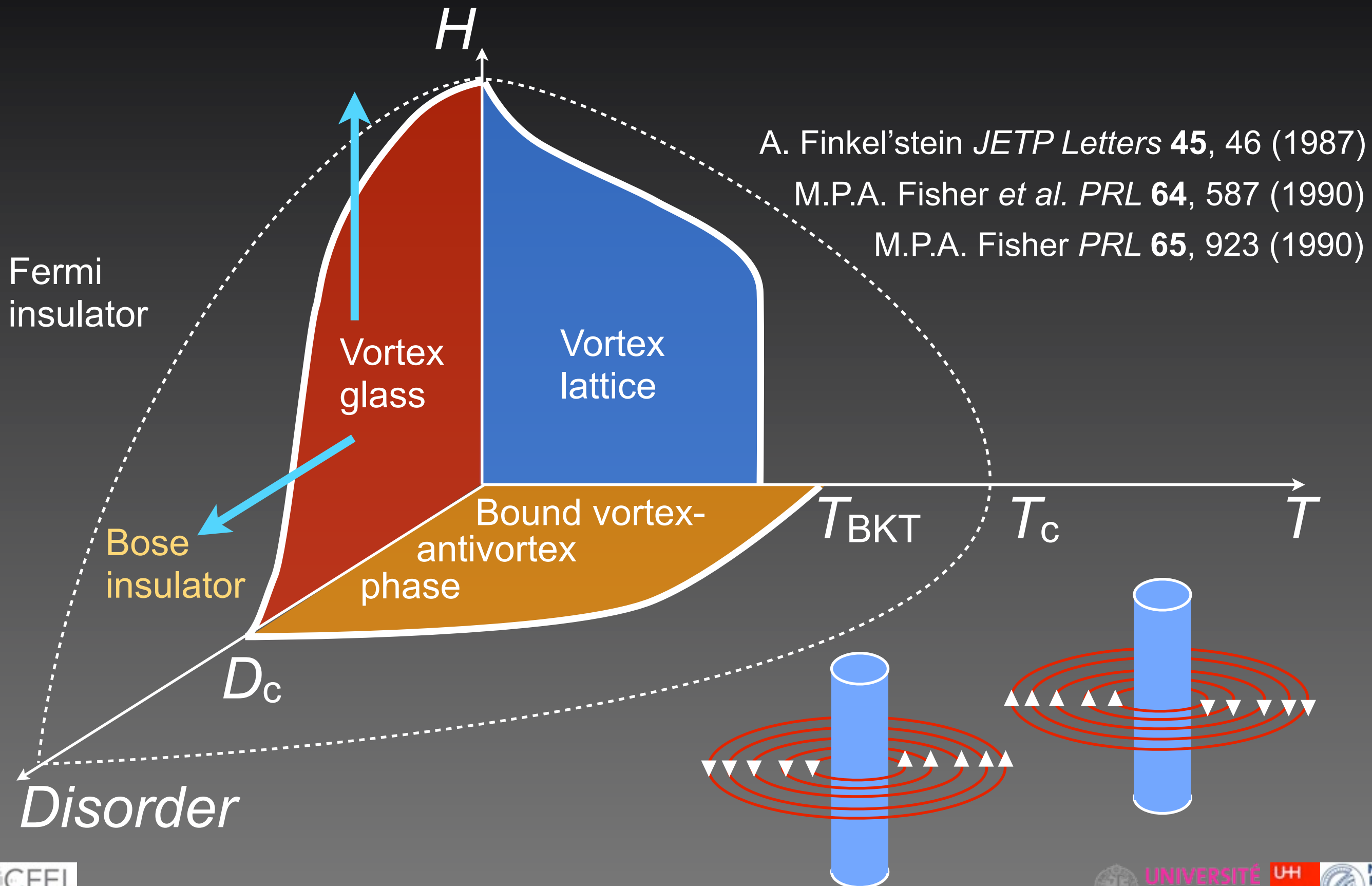
$$\mu_0 H_{\parallel}(T) = \frac{\sqrt{3}\Phi_0}{\pi d\xi_{\parallel}(T)}$$

$d \approx 10 \text{ nm}$

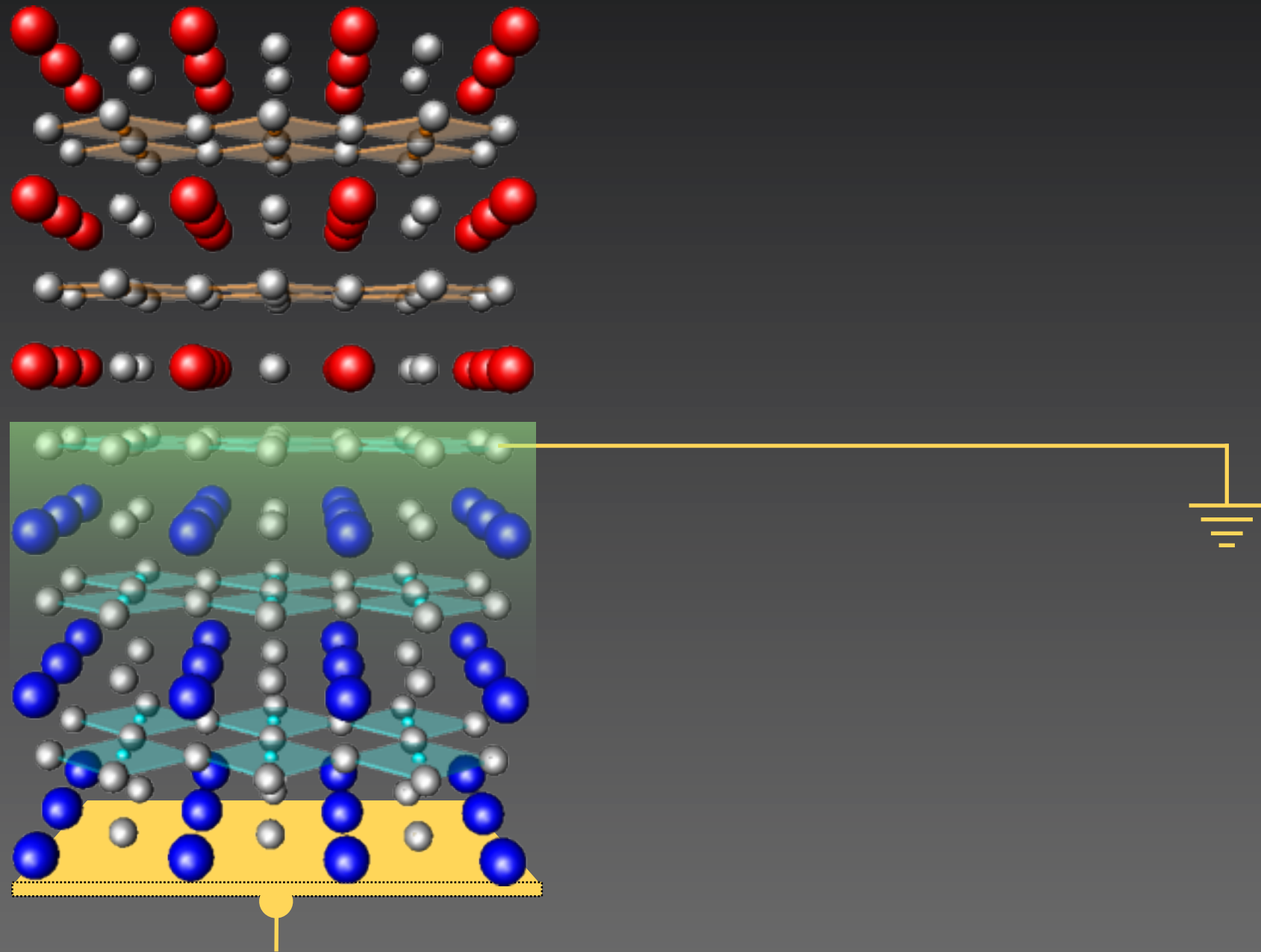
consistent with

- A. Dubroka *et al.*, *PRL* **104**, 156807 (2010)
- M. Basletic *et al.*, *Nature Materials* **7**, 621 (2008)
- O. Copie *et al.*, *PRL* **102**, 216804 (2009)
- M. Sing *et al.*, *PRL* **102**, 176805 (2009)

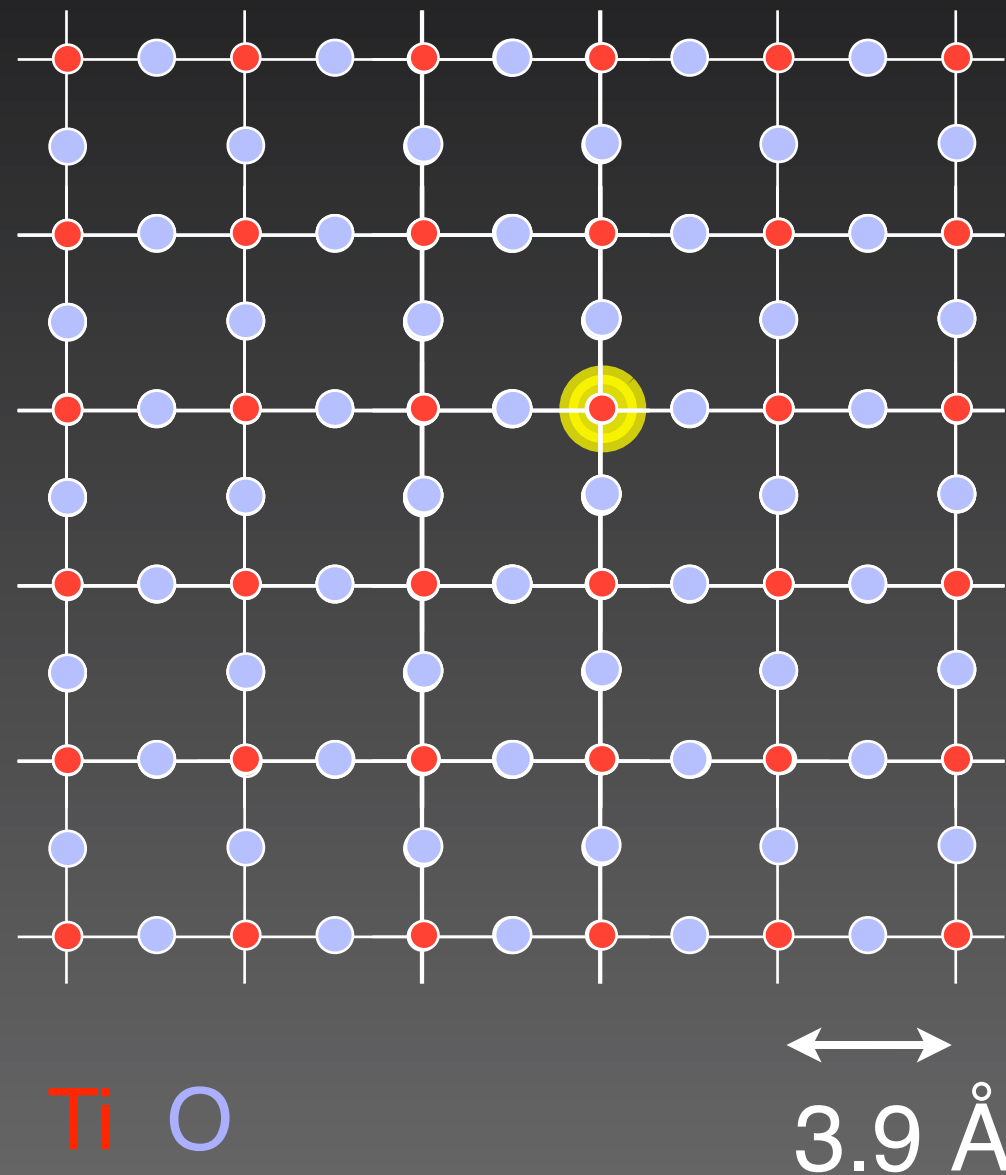
Phase diagram



Field effect experiments

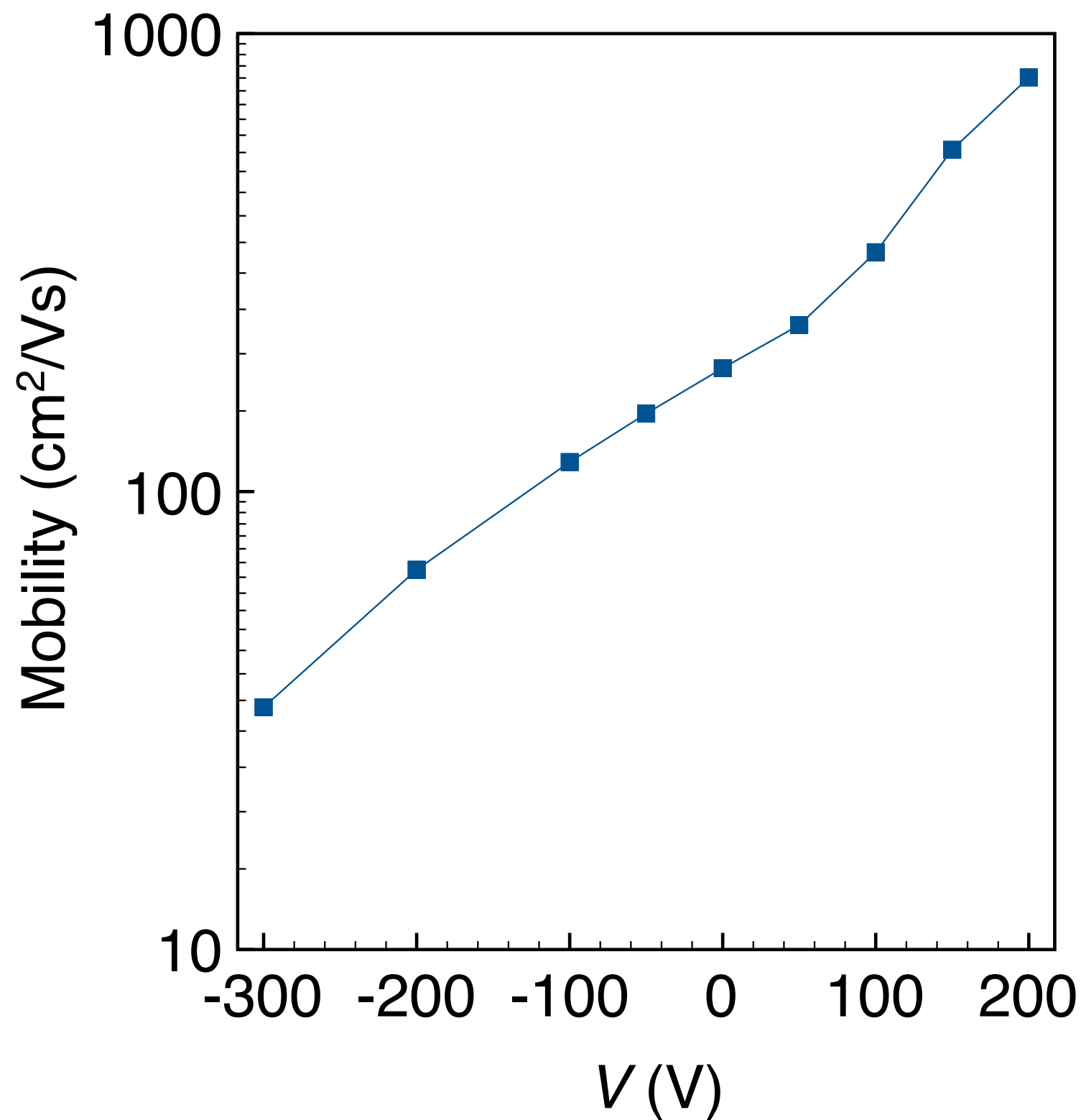


Low carrier density



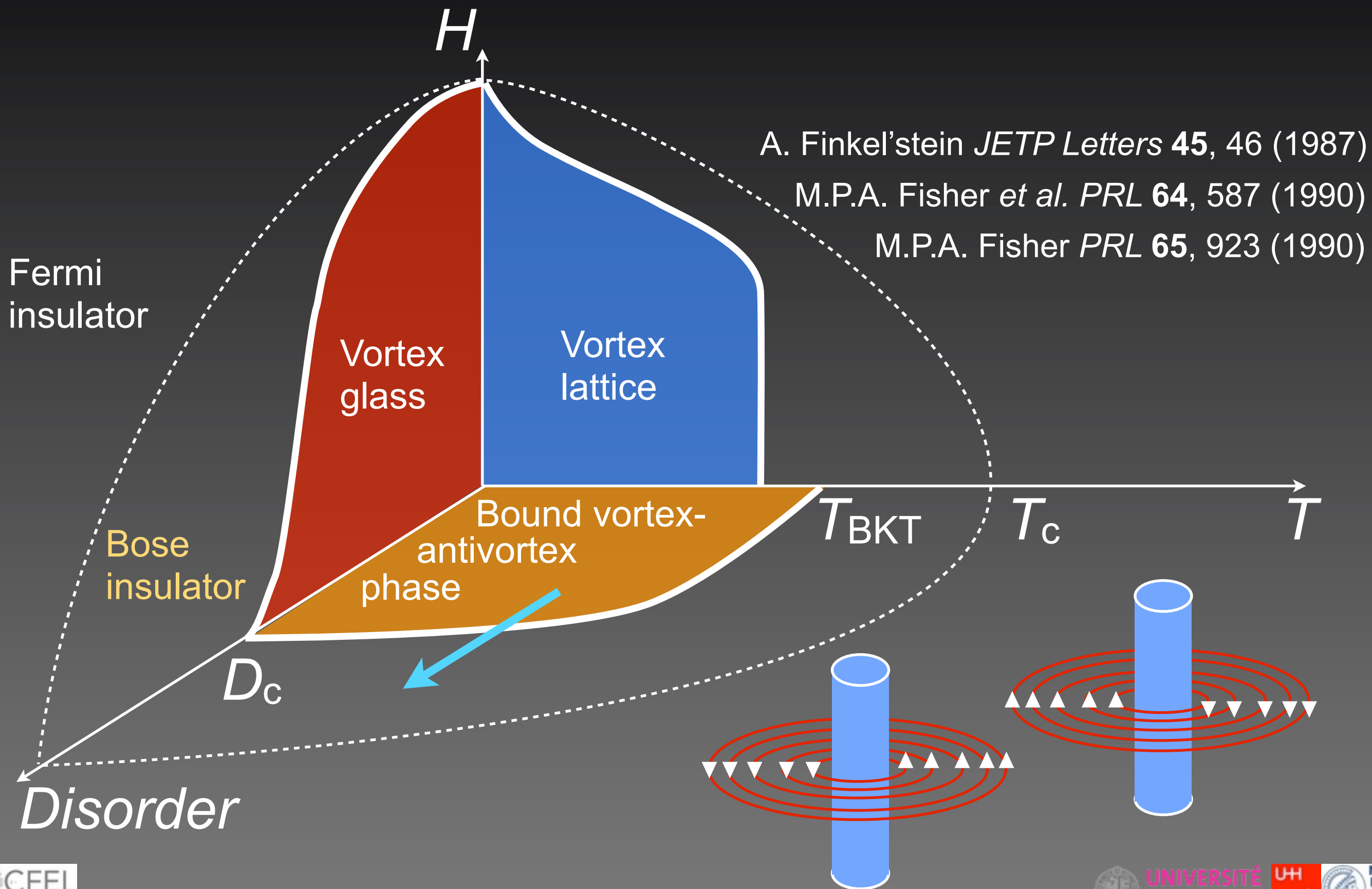
$$\sim 2-4 \times 10^{13} / \text{cm}^2$$

TiO_2 -plane

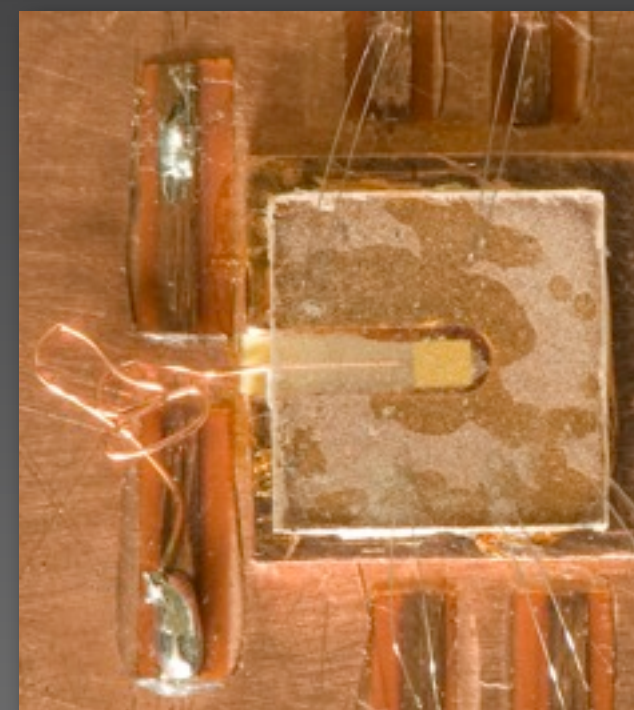
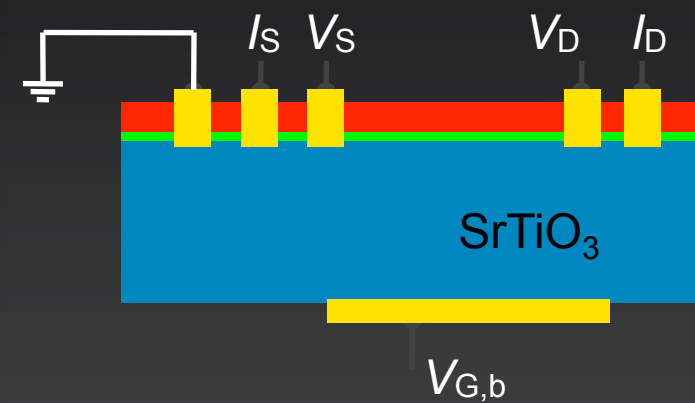
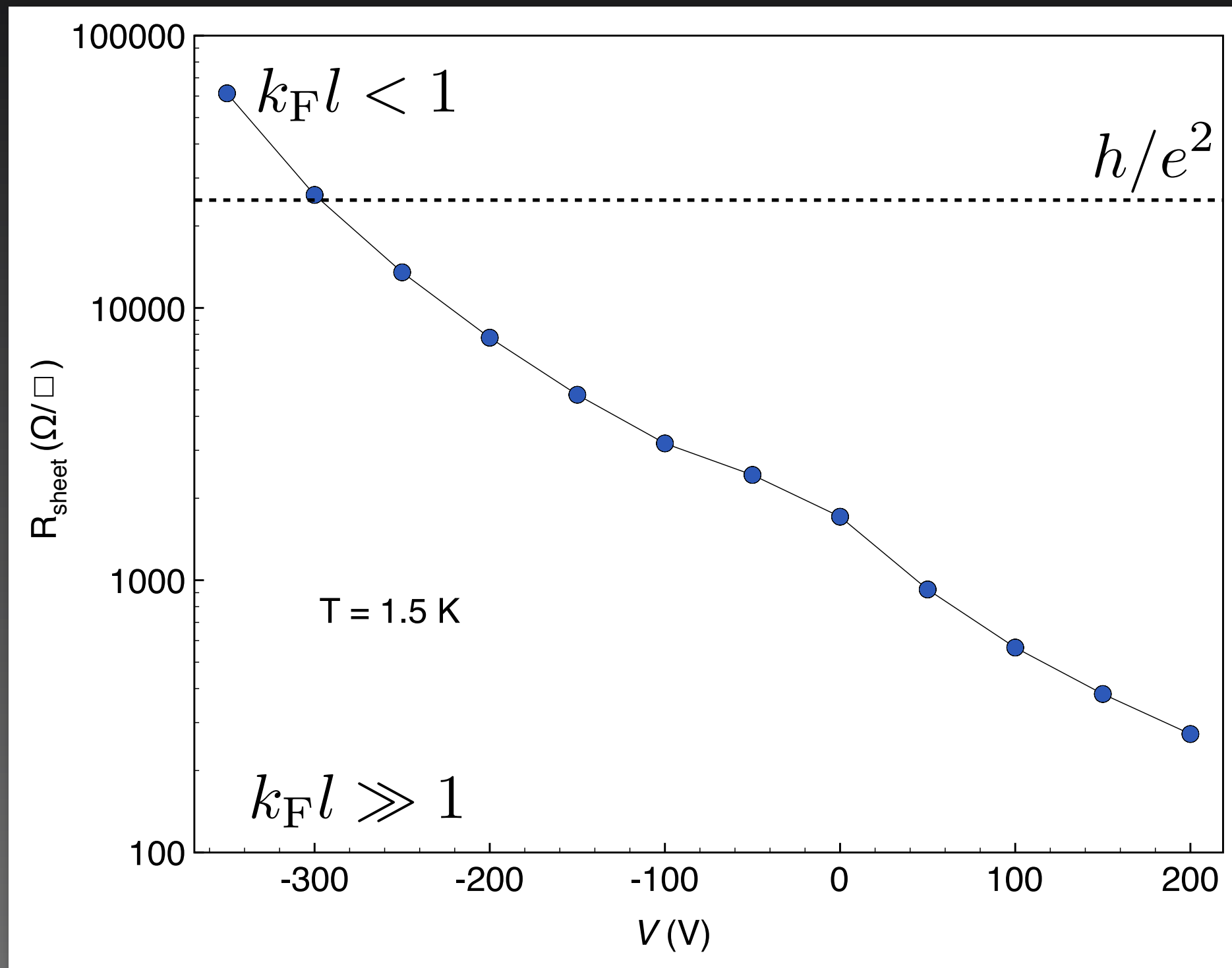


C. Bell *et al.*, *PRL* **103**, 226802 (2009)
A.D. Caviglia *et al.*, *PRL* **104**, 126803 (2010)

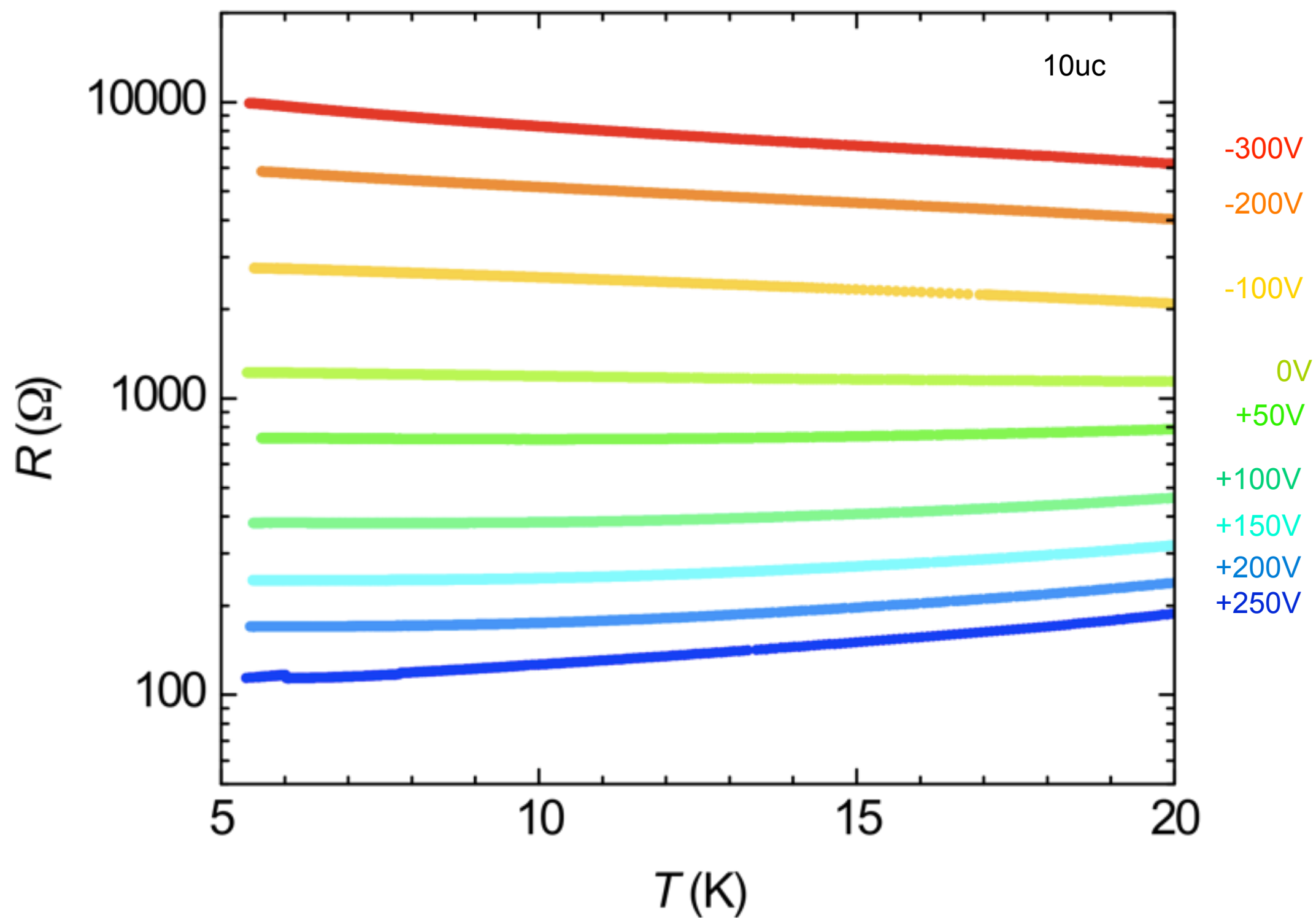
Phase diagram



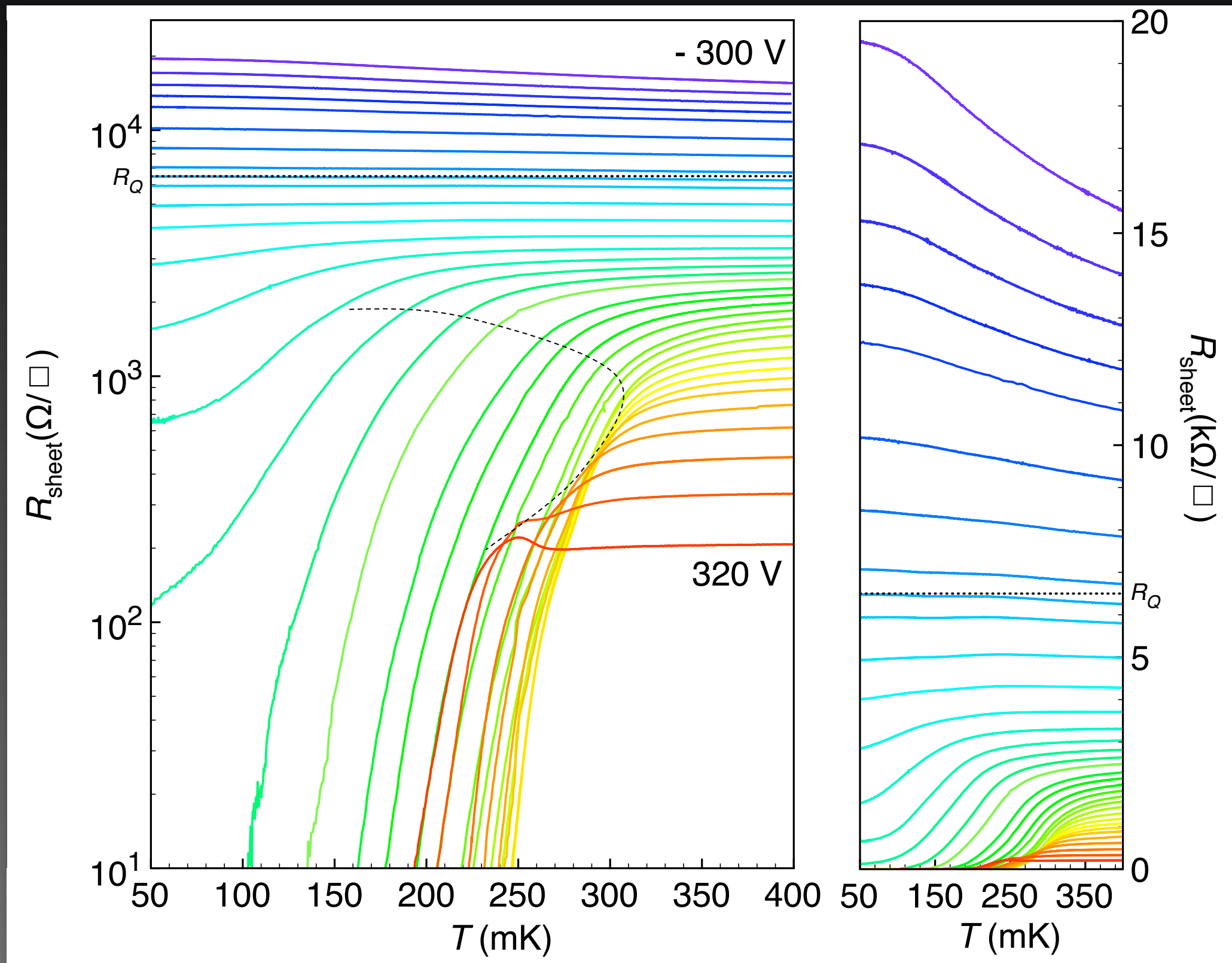
R_{sheet} modulation



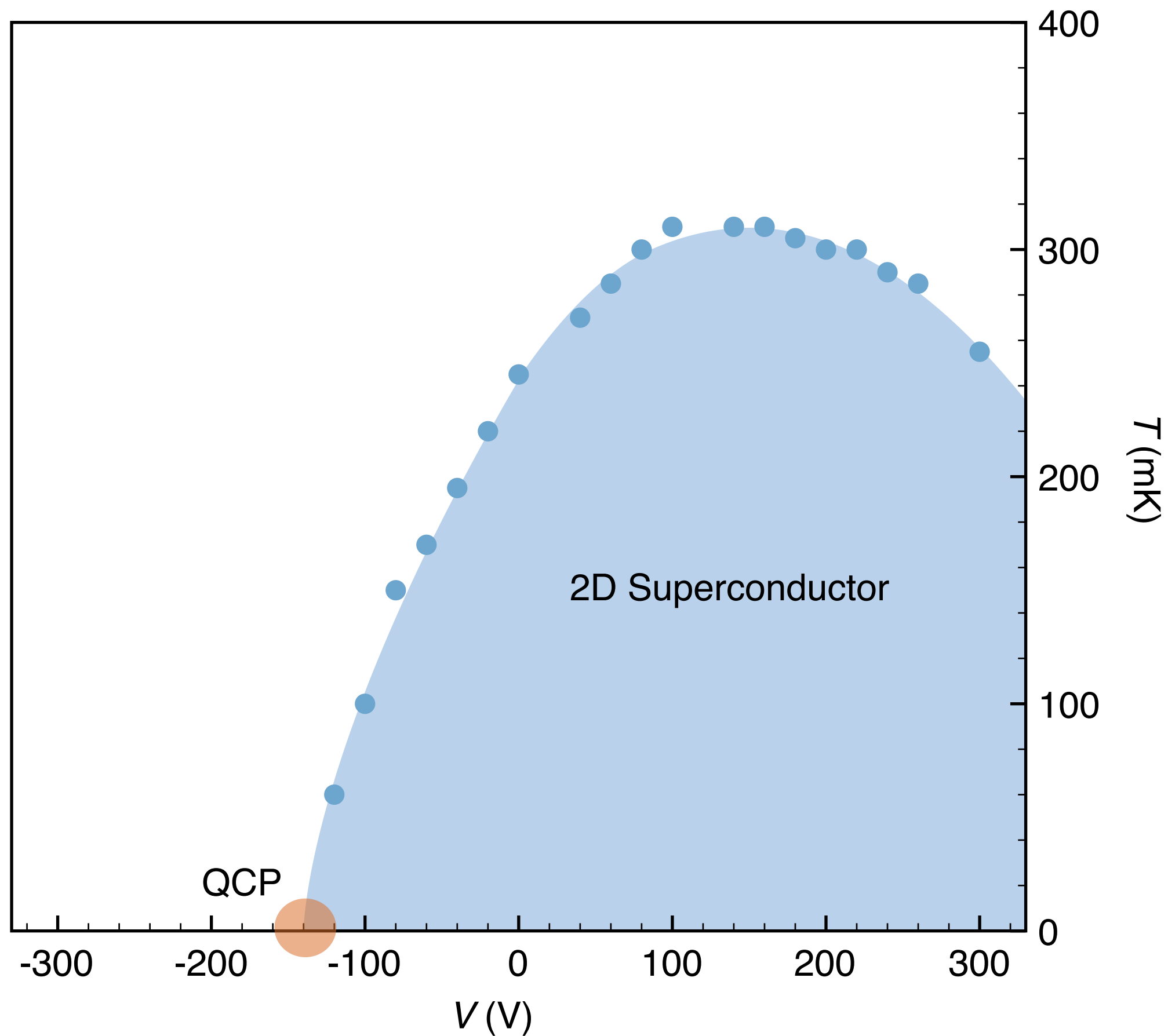
5 mm

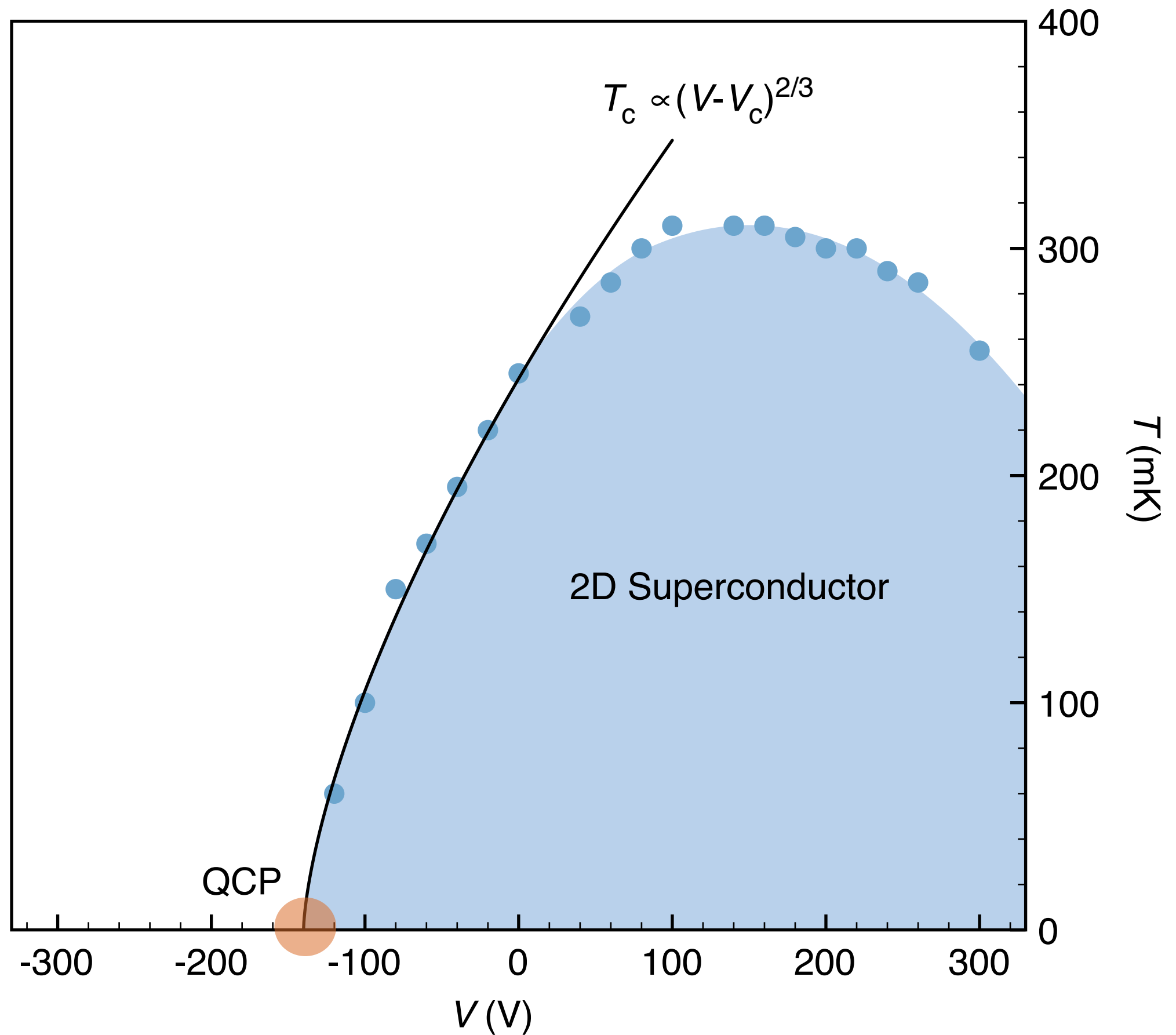


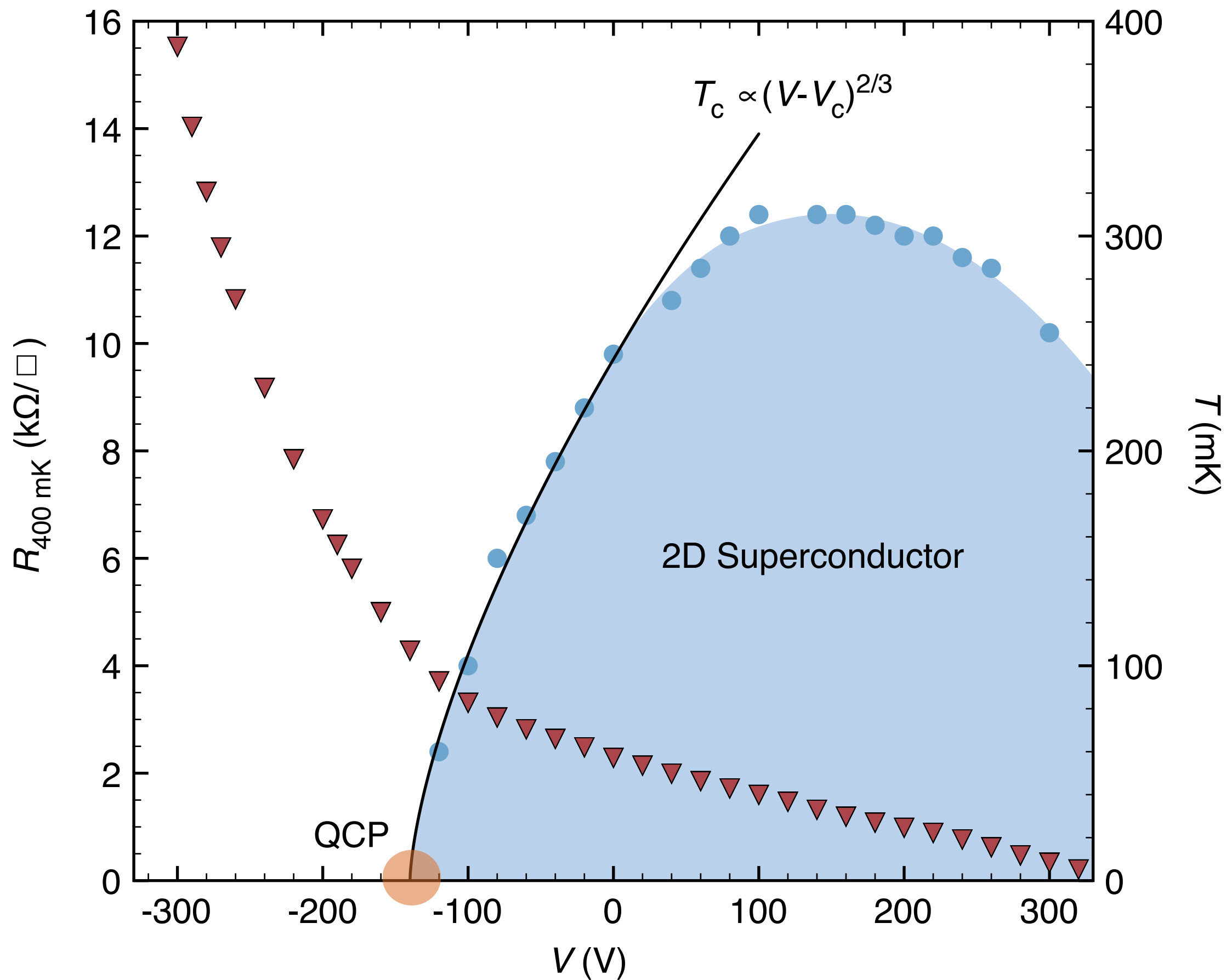
Modulation of SC

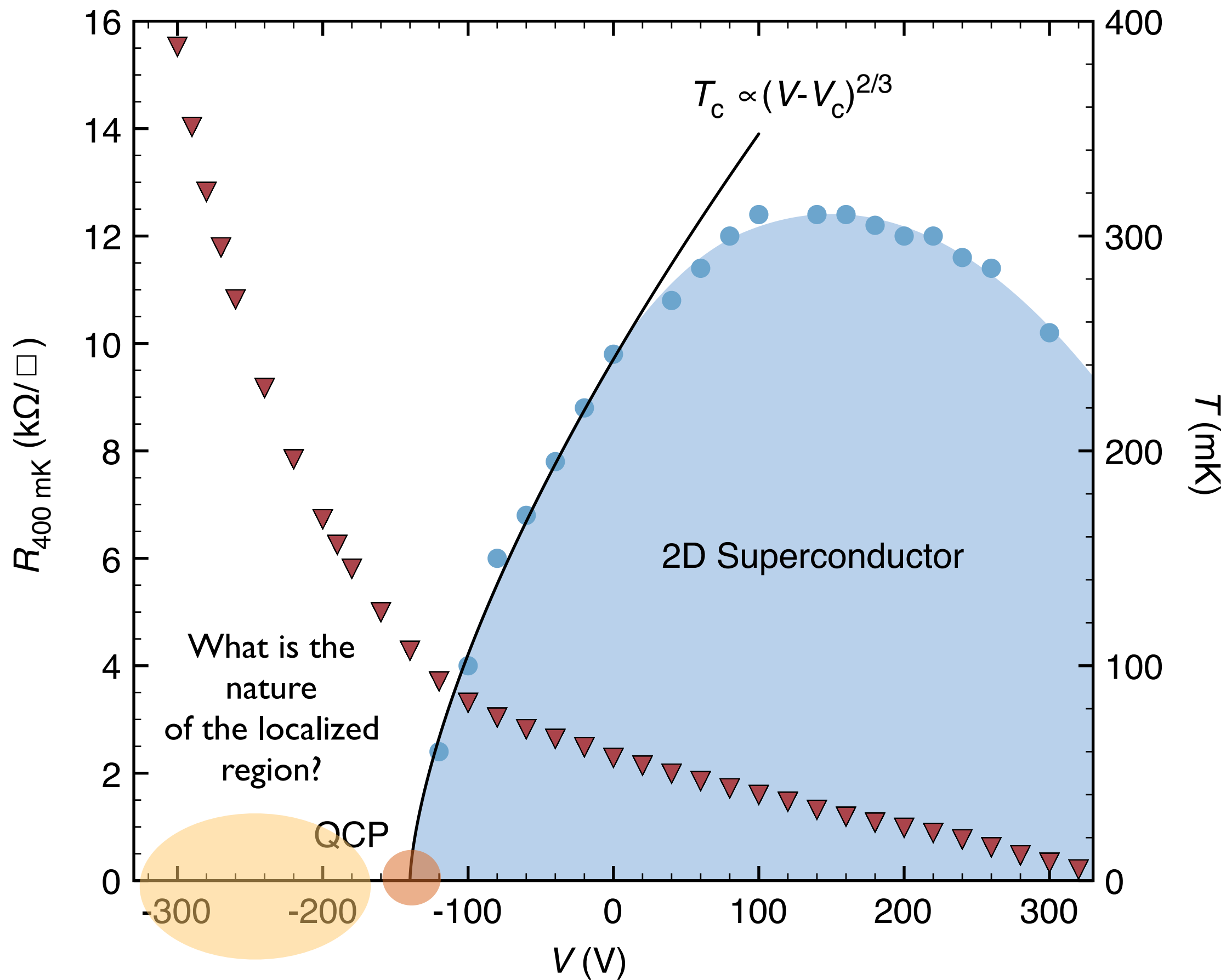


A.D. Caviglia et al, *Nature* 456, 625 (2008)

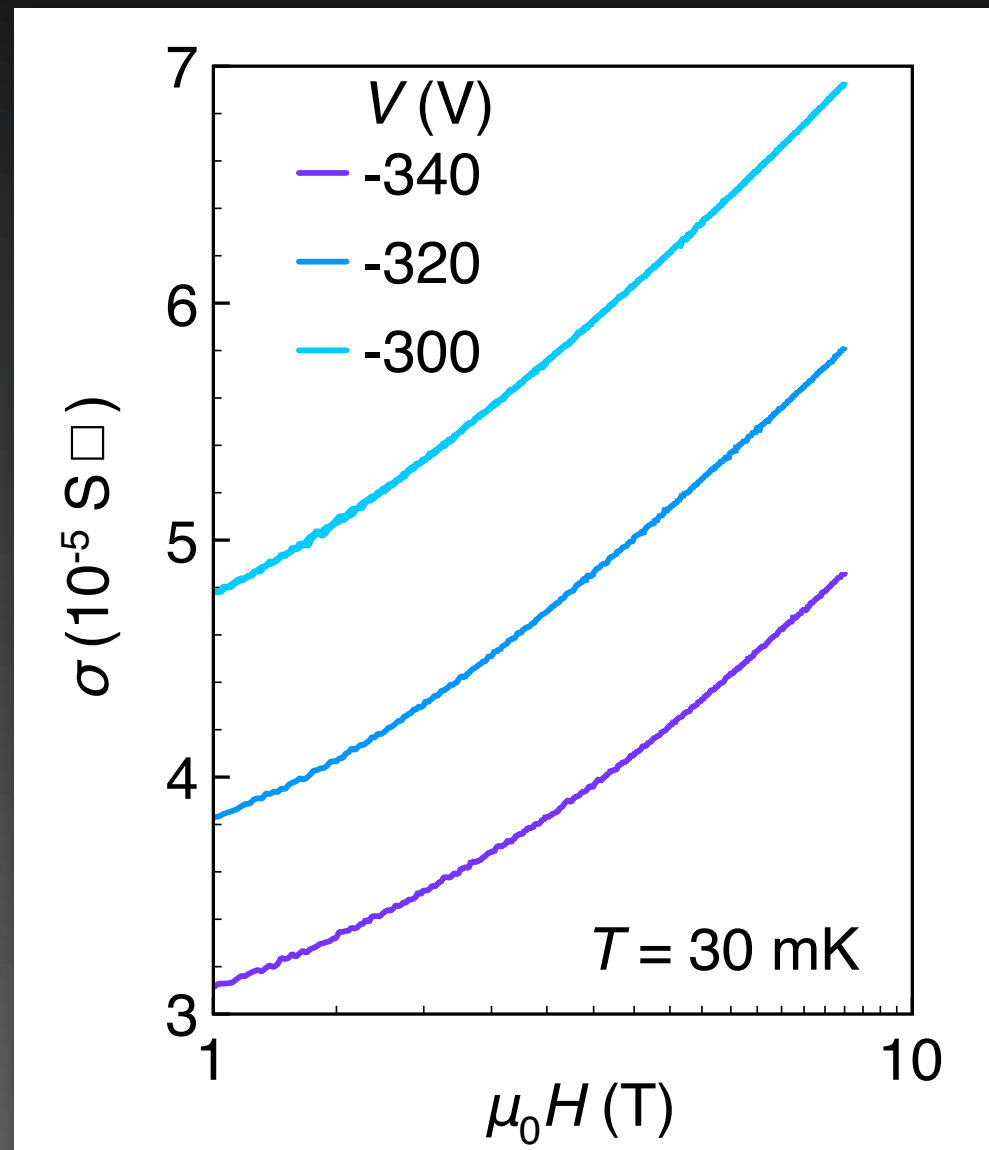




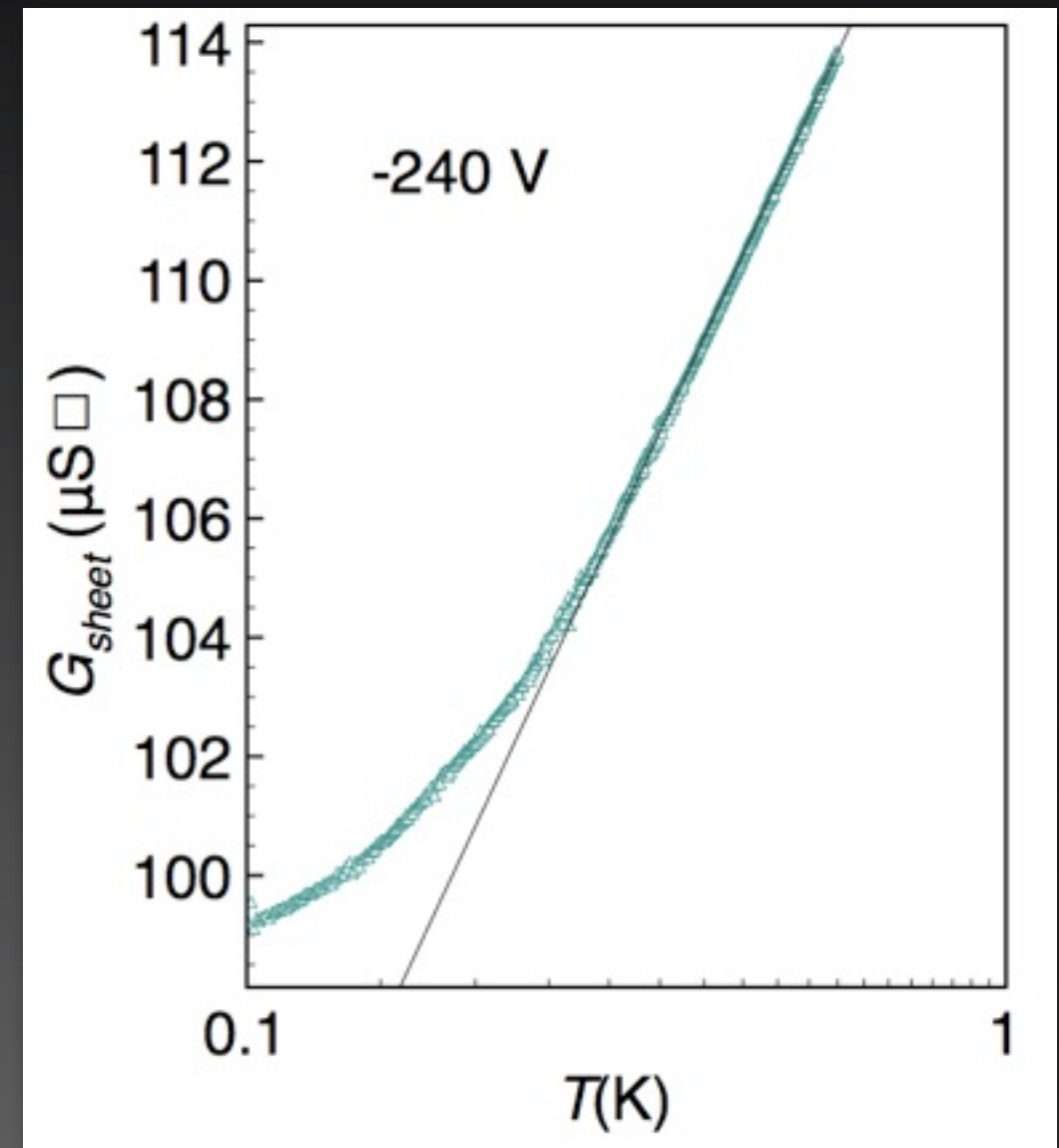




2D Weak localization

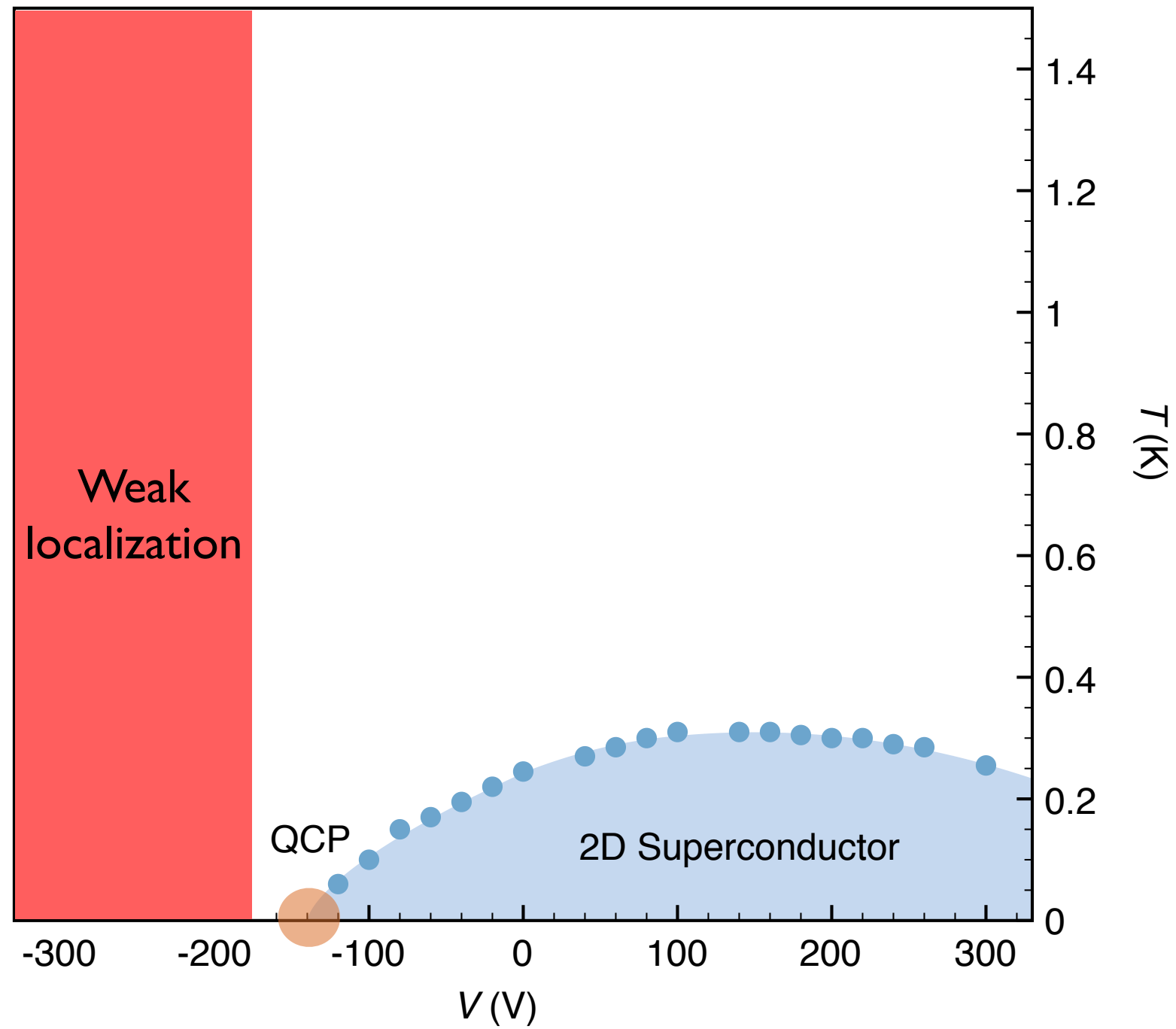


$$\frac{d\sigma}{d \ln H} = \frac{e^2}{\pi h} = 1.23 \times 10^{-5} S$$

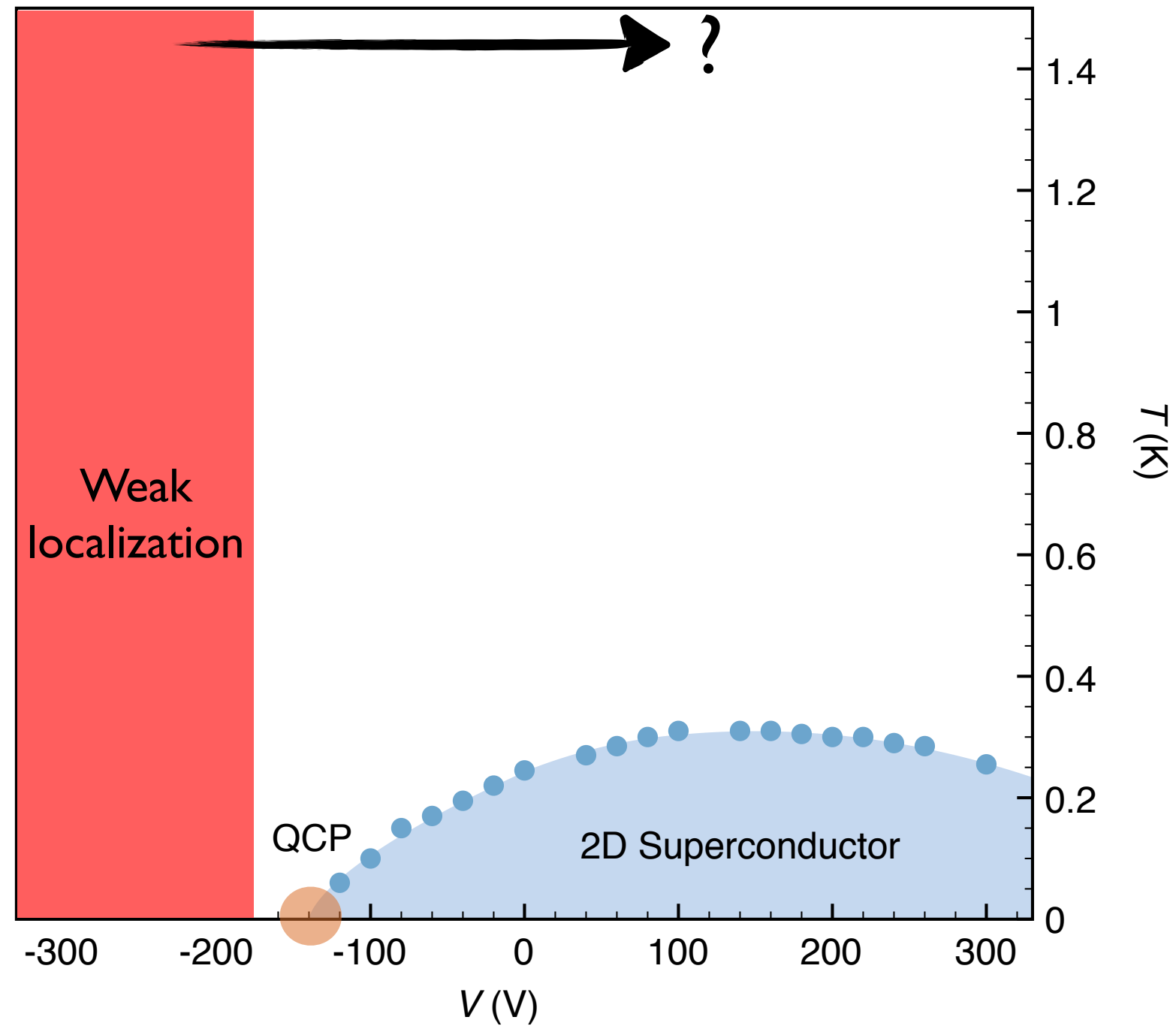


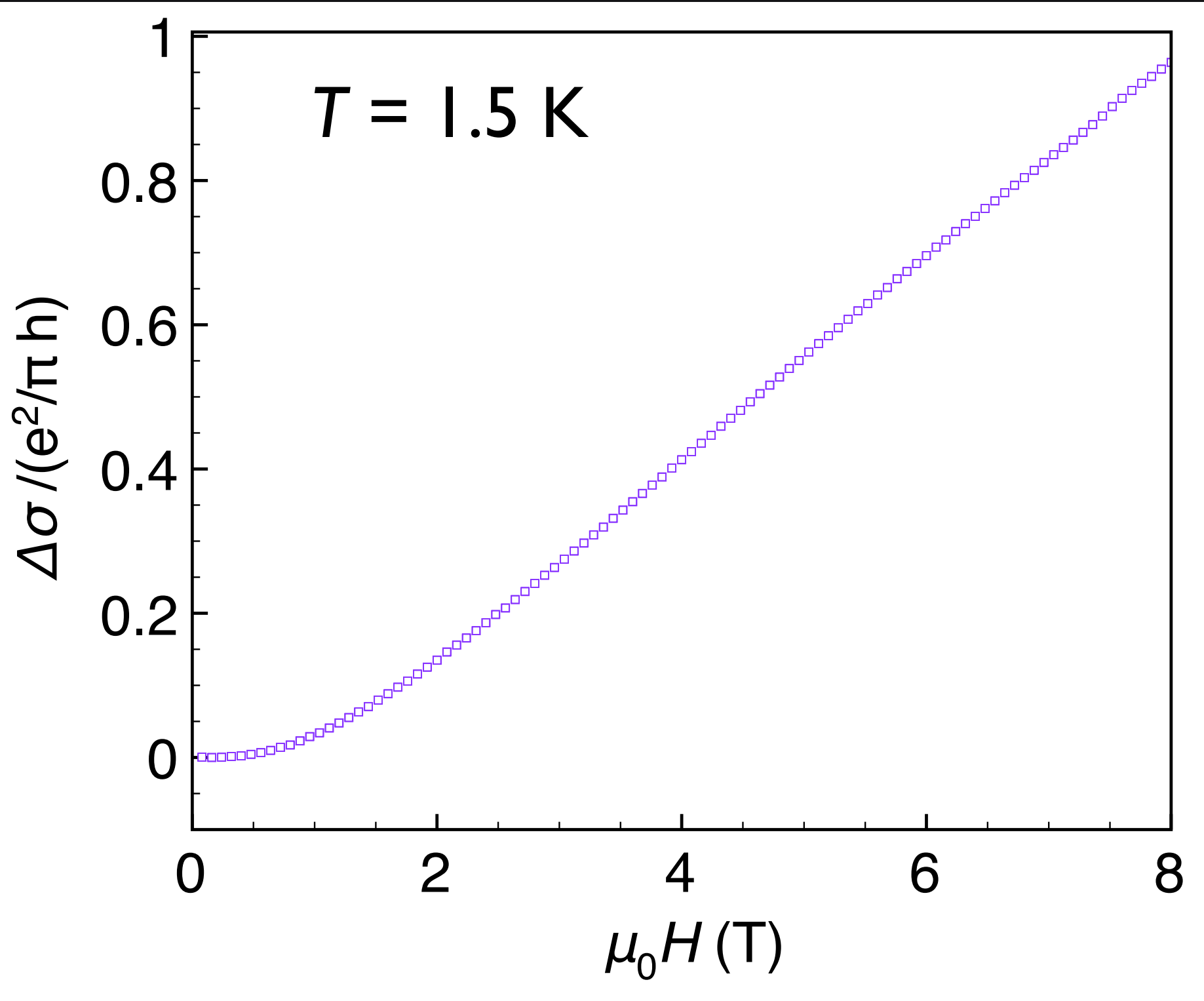
$$\frac{d\sigma}{d \ln T} = \frac{e^2}{\pi h} = 1.23 \times 10^{-5} S$$

Phase diagram

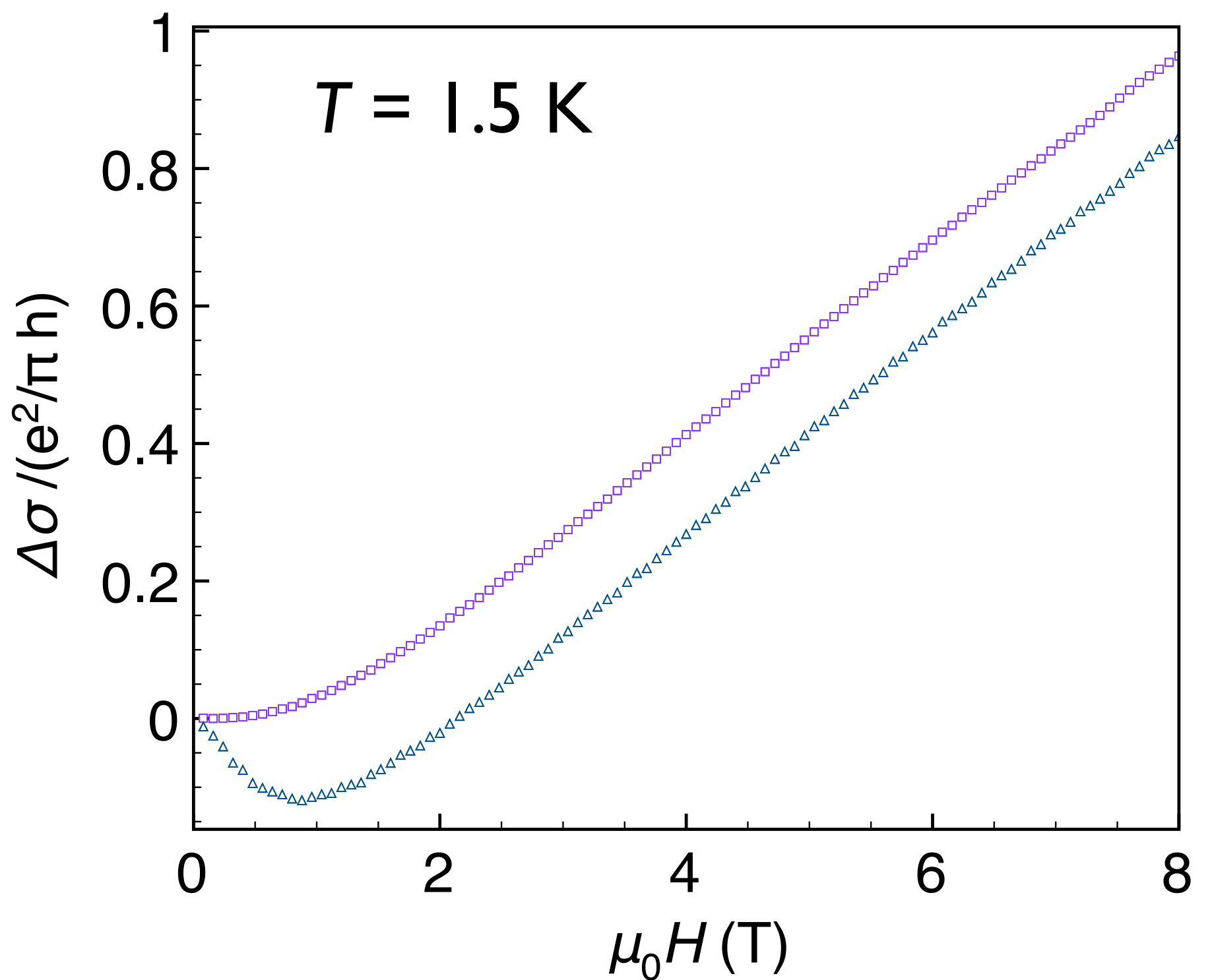


Phase diagram



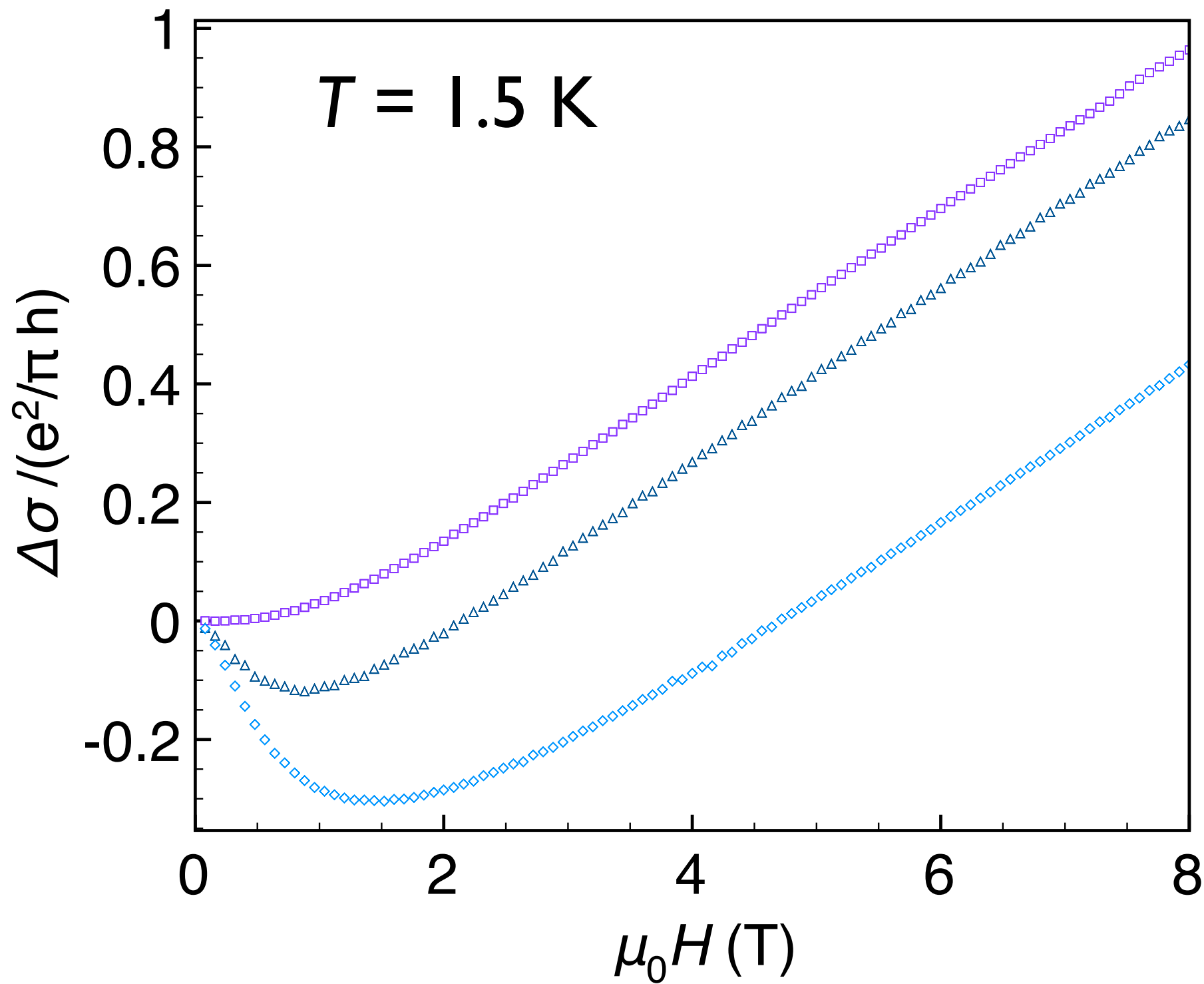


-300 V



-300 V

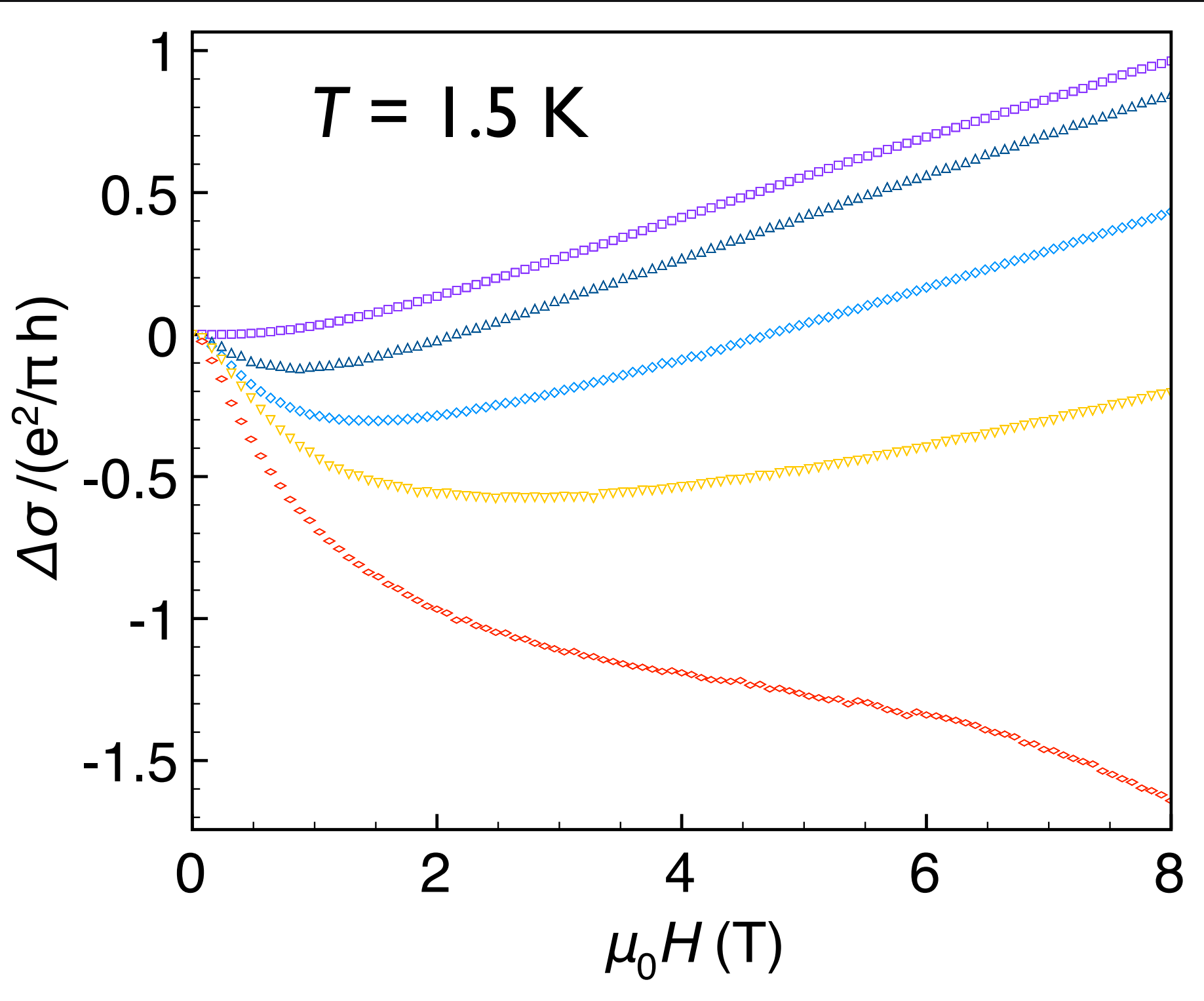
-100 V



-300 V

-100 V

-50 V



-300 V

-100 V

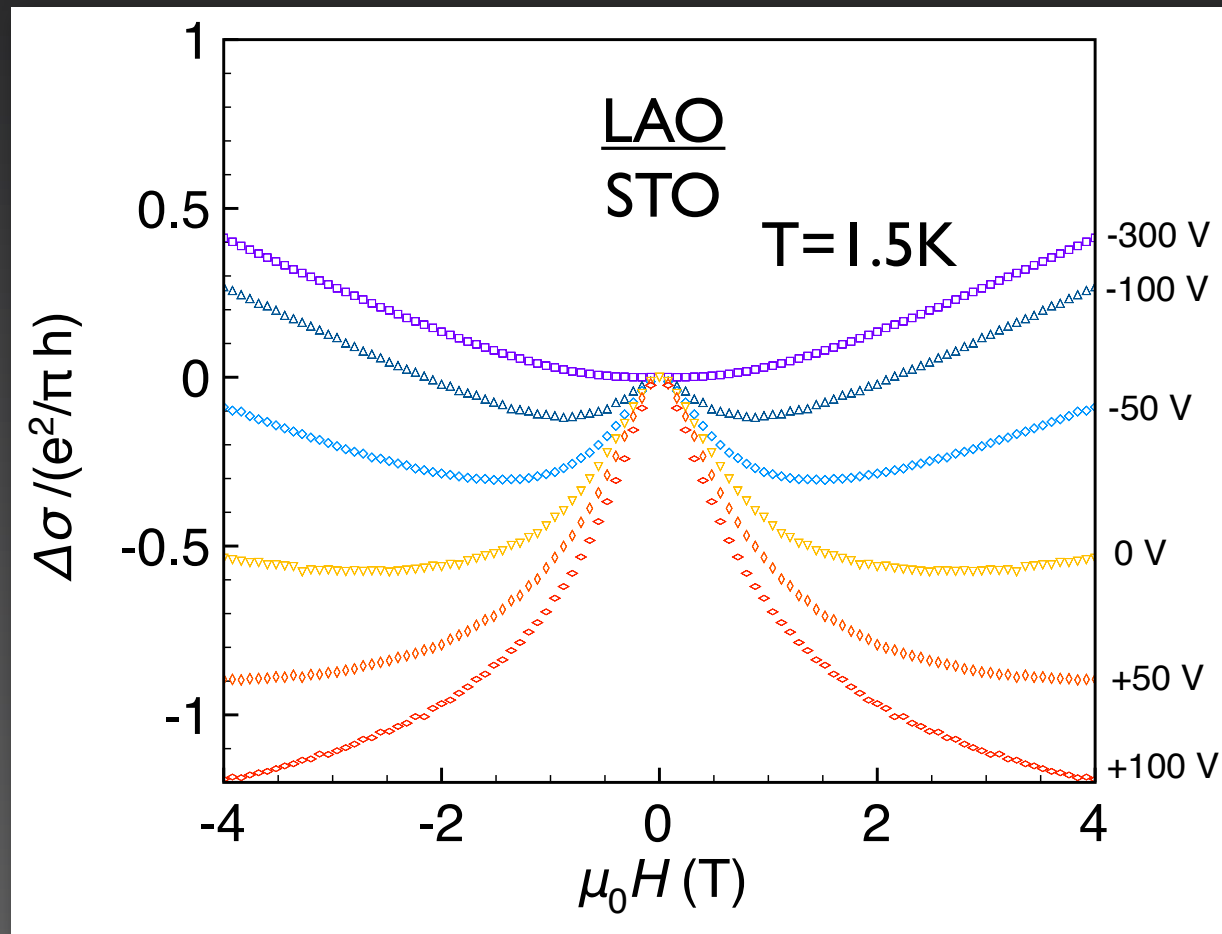
-50 V

0 V

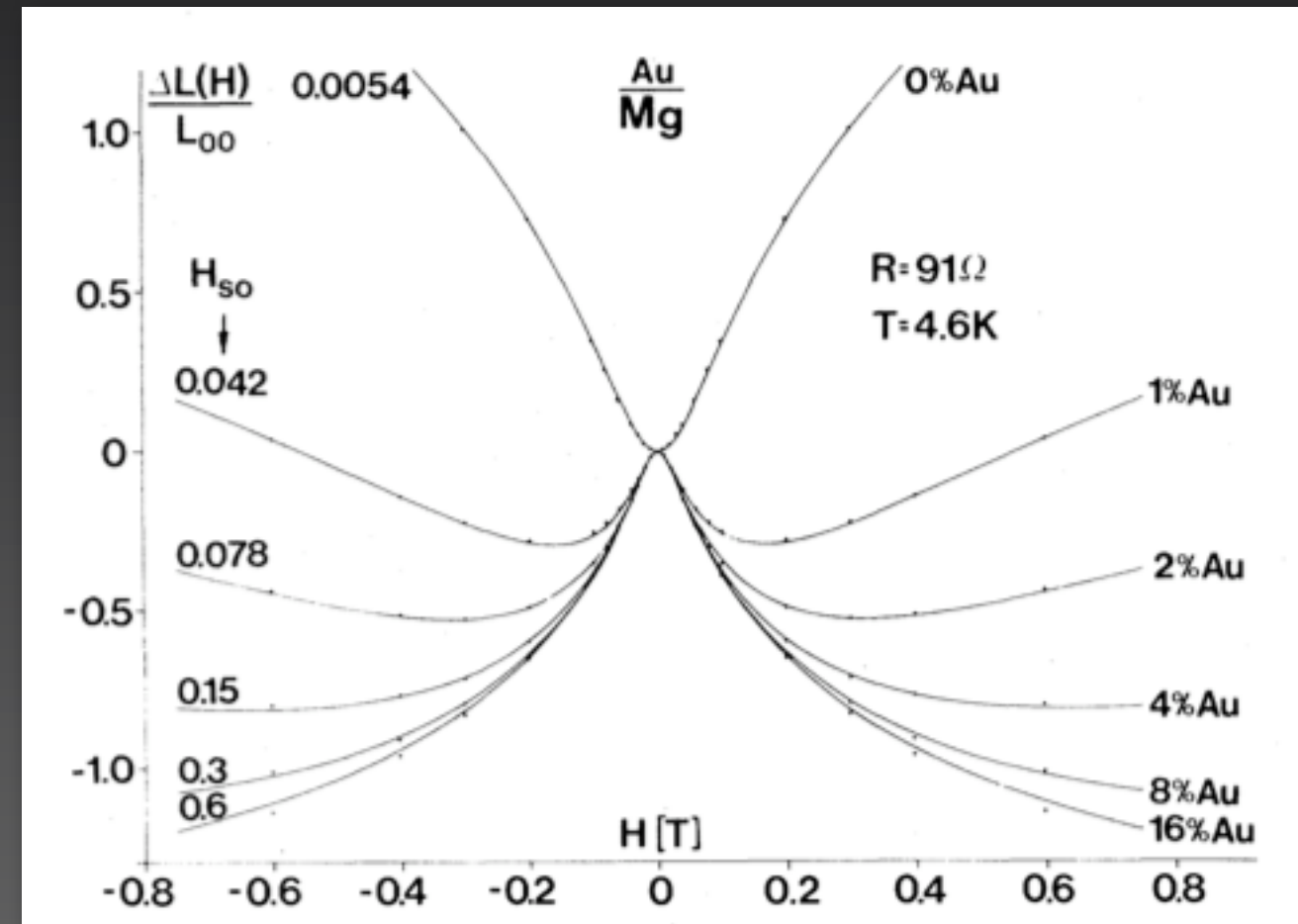
100 V

Weak antilocalization

Strong spin-orbit interaction



A.D. Caviglia et al., *PRL* **104**, 126803 (2010)



G. Bergman, *PRL* **48**, 1046 (1982)
in metallic thin films

J.B. Miller et al., *PRL* **90**, 076807 (2003)
in GaAs/AlGaAs 2DEG

$$\frac{\Delta\sigma(H)}{\sigma_0} = \Psi\left(\frac{H}{H_i + H_{so}}\right) + \frac{1}{2\sqrt{1-\gamma^2}}\Psi\left(\frac{H}{H_i + H_{so}\left(1 + \sqrt{1-\gamma^2}\right)}\right) \\ - \frac{1}{2\sqrt{1-\gamma^2}}\Psi\left(\frac{H}{H_i + H_{so}\left(1 - \sqrt{1-\gamma^2}\right)}\right)$$

Digamma function

$$\Psi(x) = \ln(x) + \psi\left(\frac{1}{2} + \frac{1}{x}\right)$$

$$\sigma_0 = e^2/\pi h = 1.2 \cdot 10^{-5} \text{ S}$$

Sadamichi Maekawa and Hidetoshi Fukuyama, *Journal of the Physical Society of Japan* **50**, 2516 (1981)

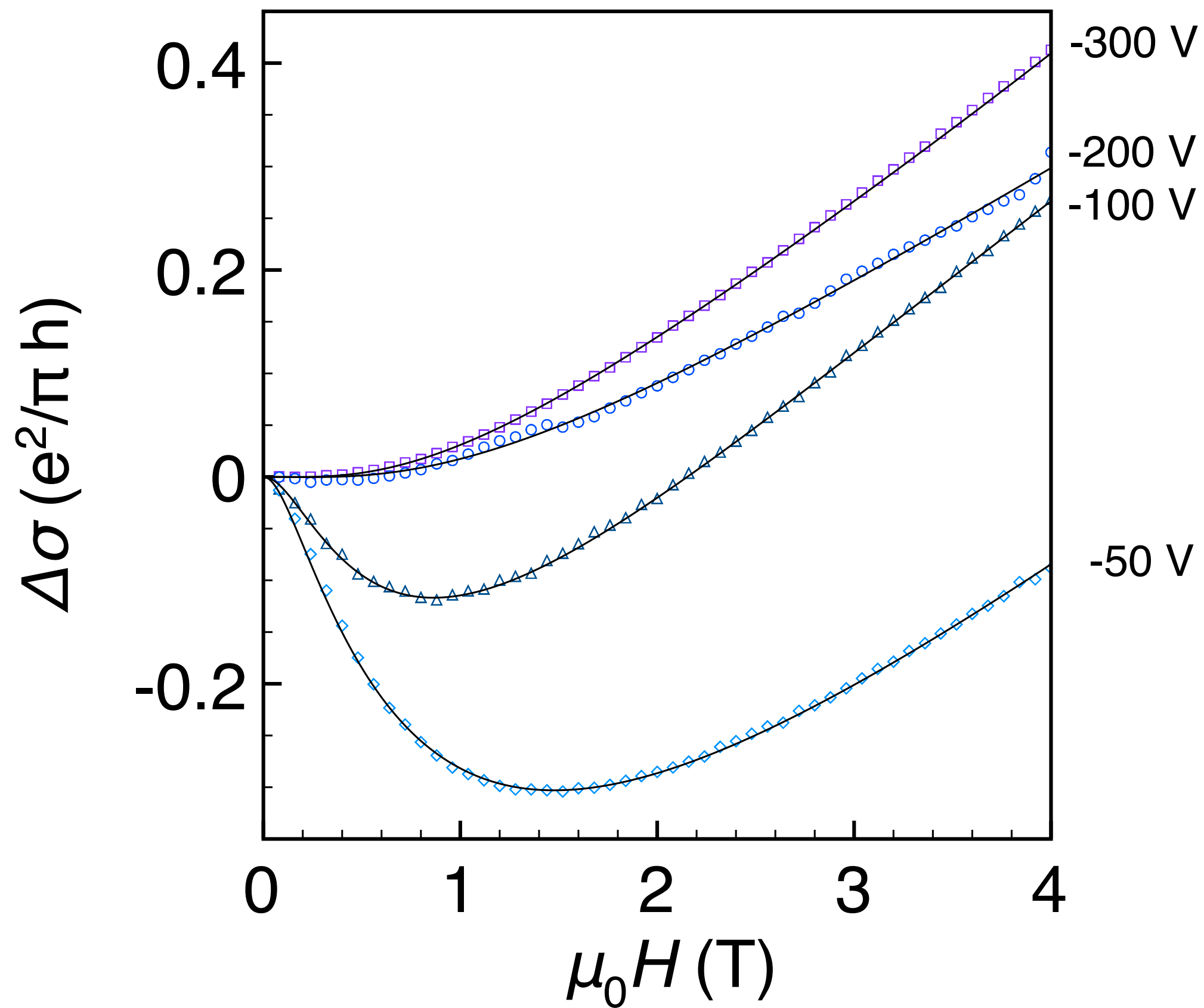
fitting parameters

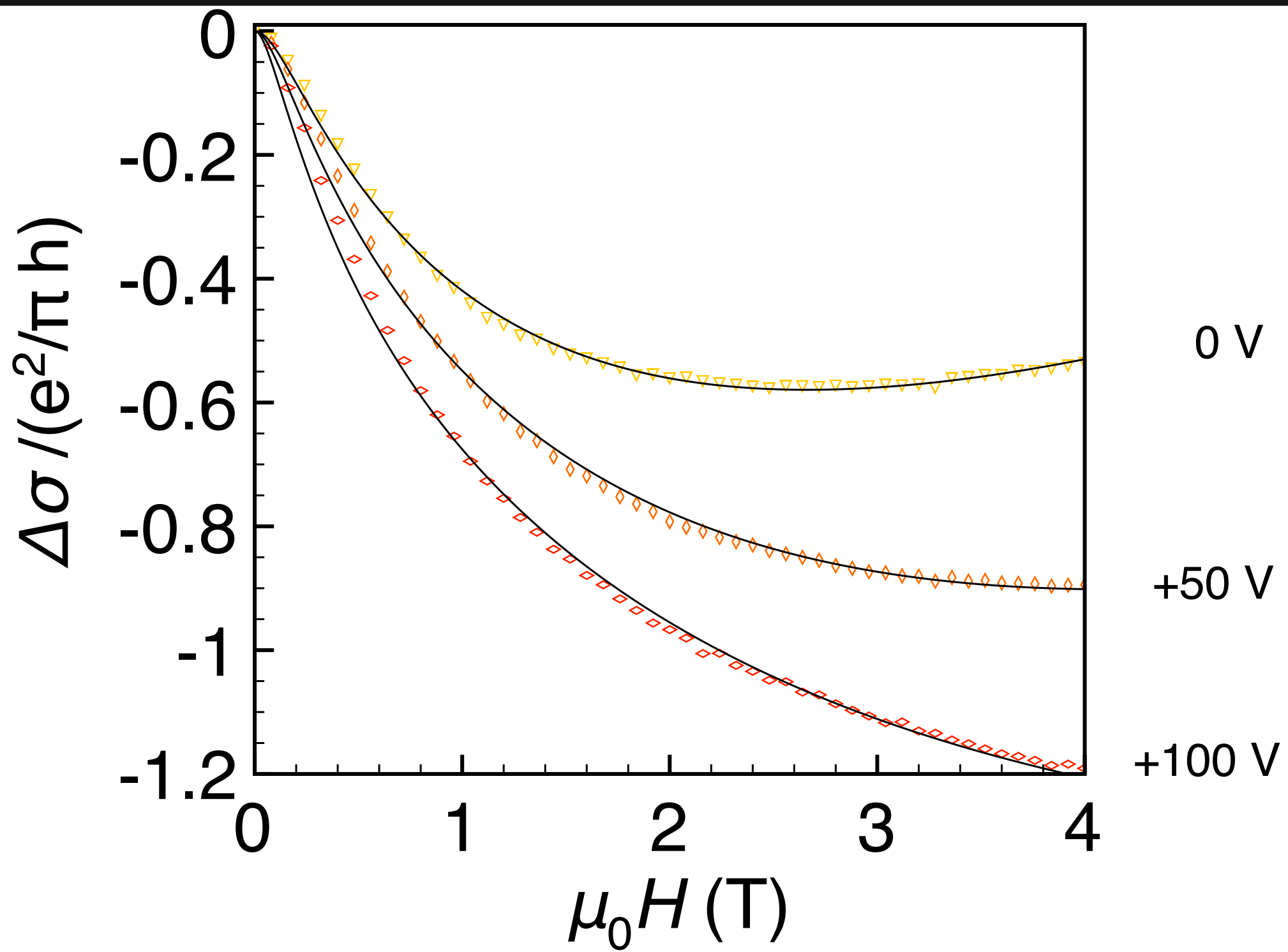
$$H_i = \hbar/4eD\tau_i$$

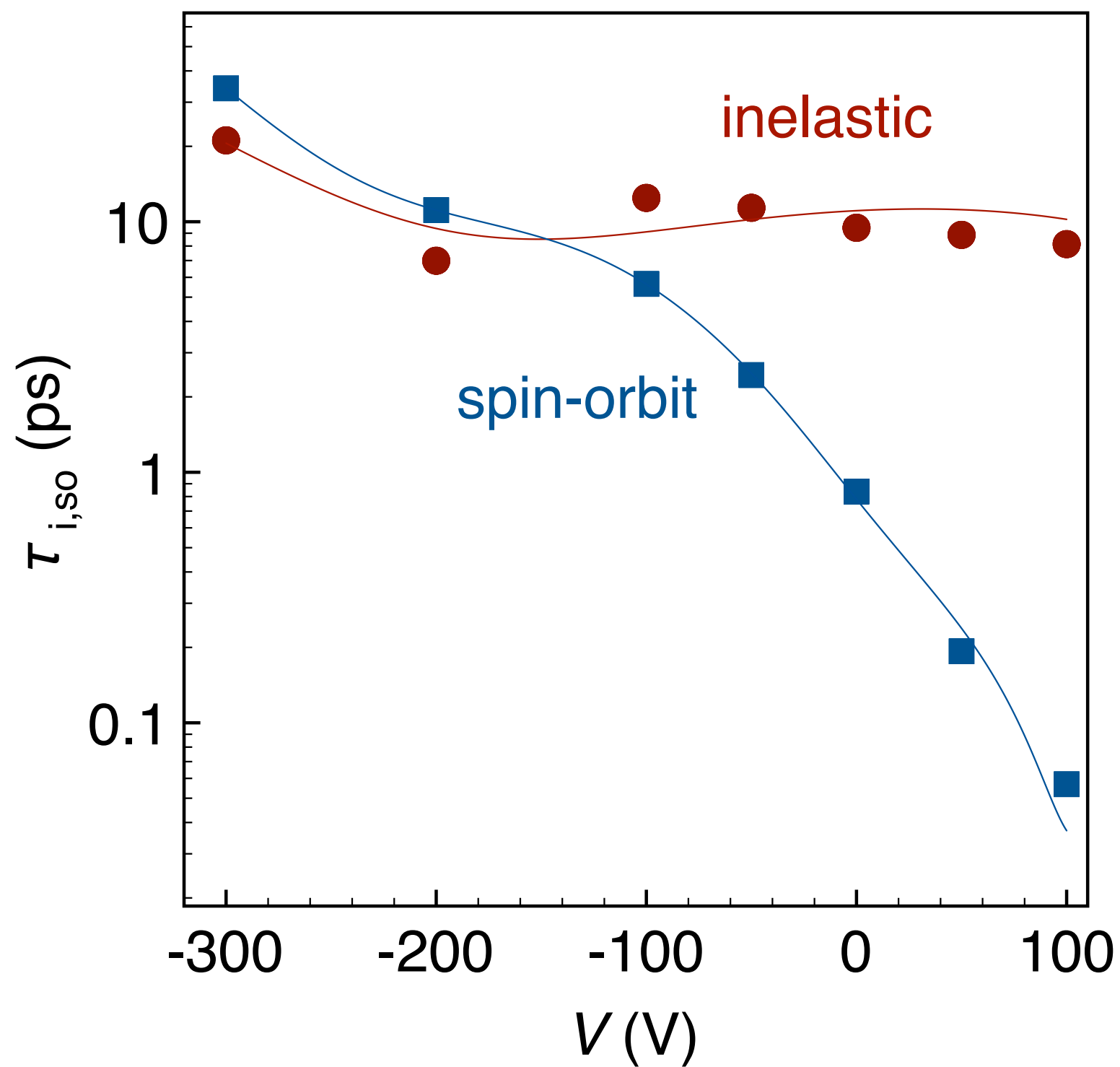
$$H_{so} = \hbar/4eD\tau_{so} \propto m^2$$

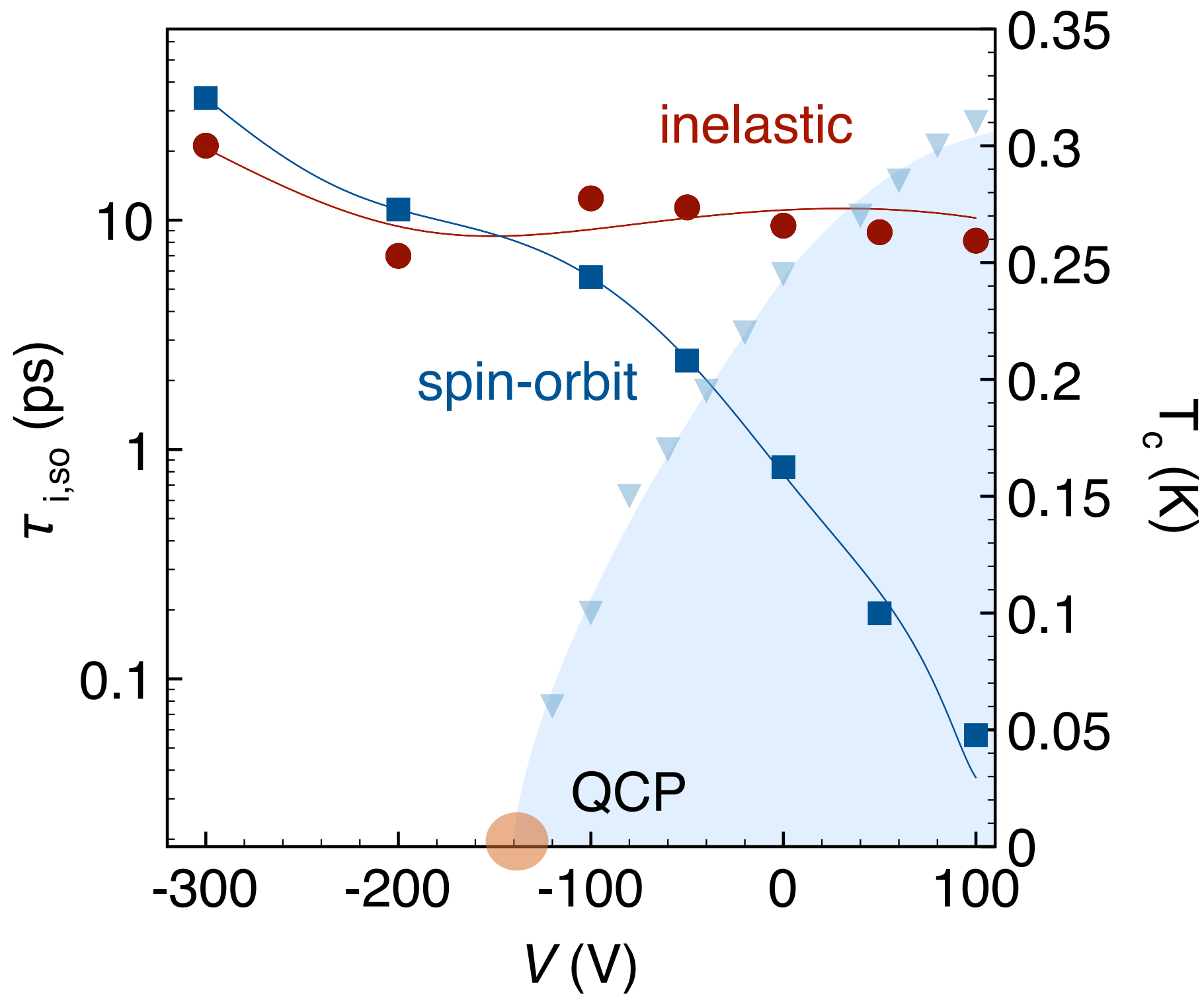
$$\gamma = g\mu_B H/4eDH_{so}$$

Nota bene: there is no
scale parameter.
It's a variation of
conductance on a
universal scale

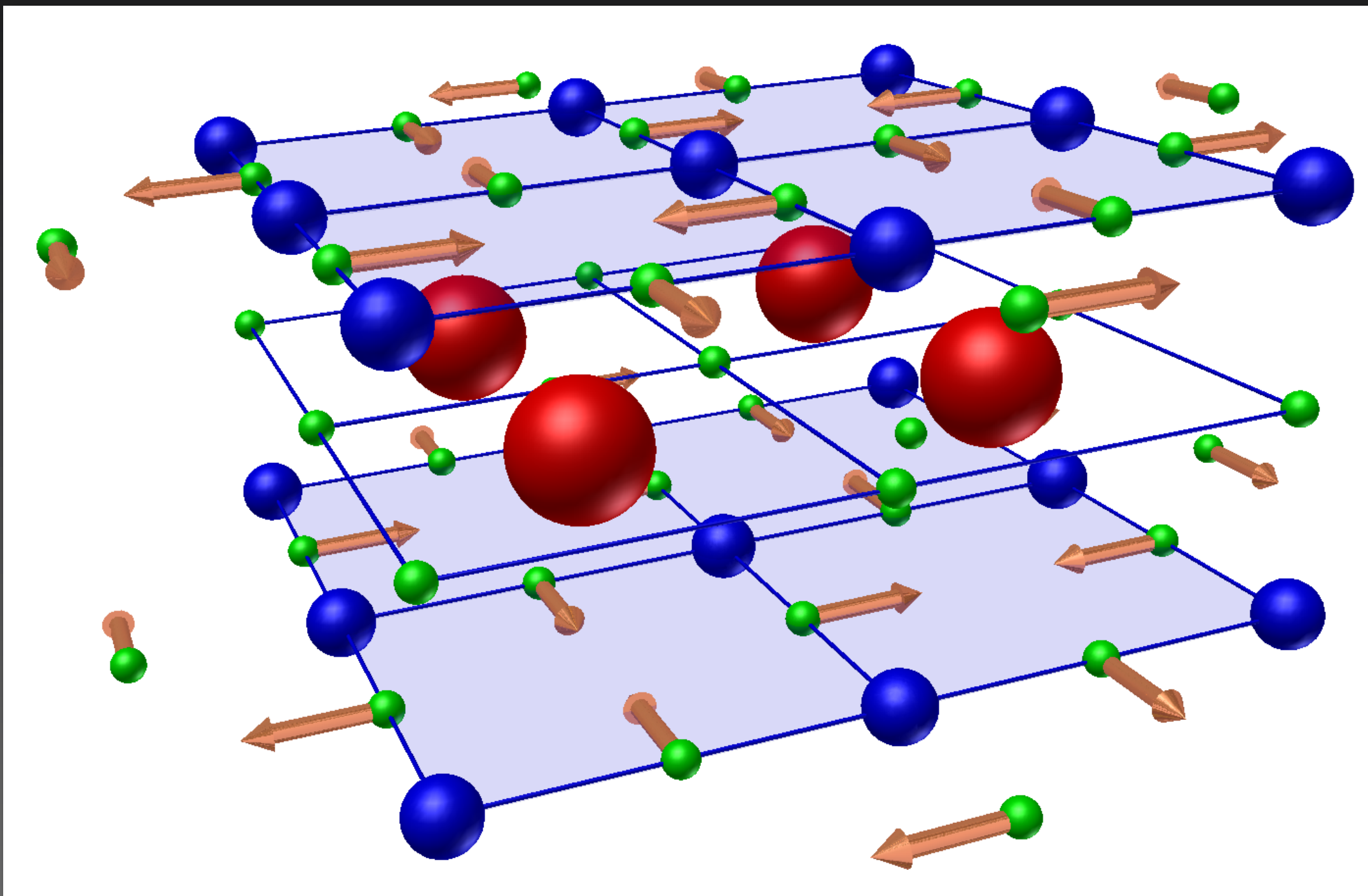


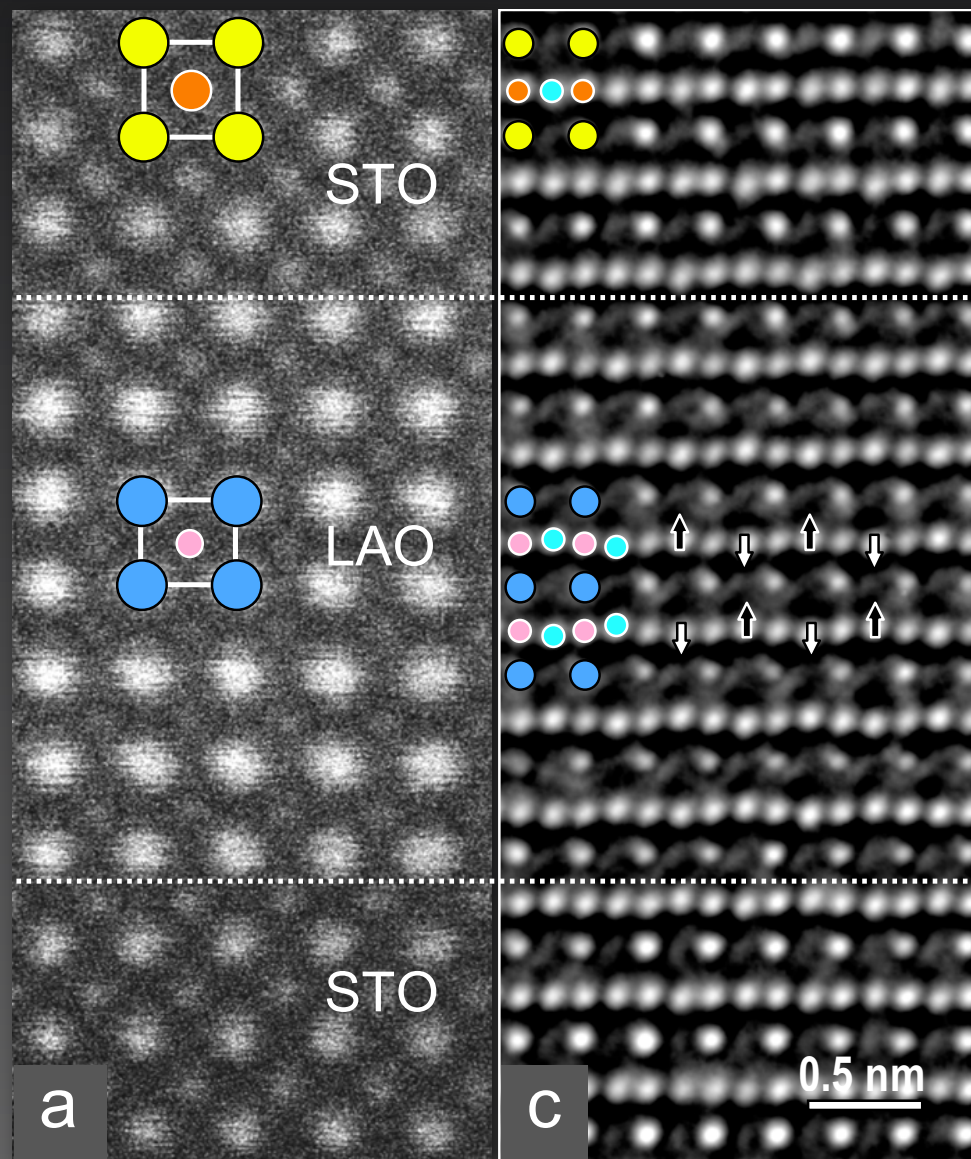






AFD mode

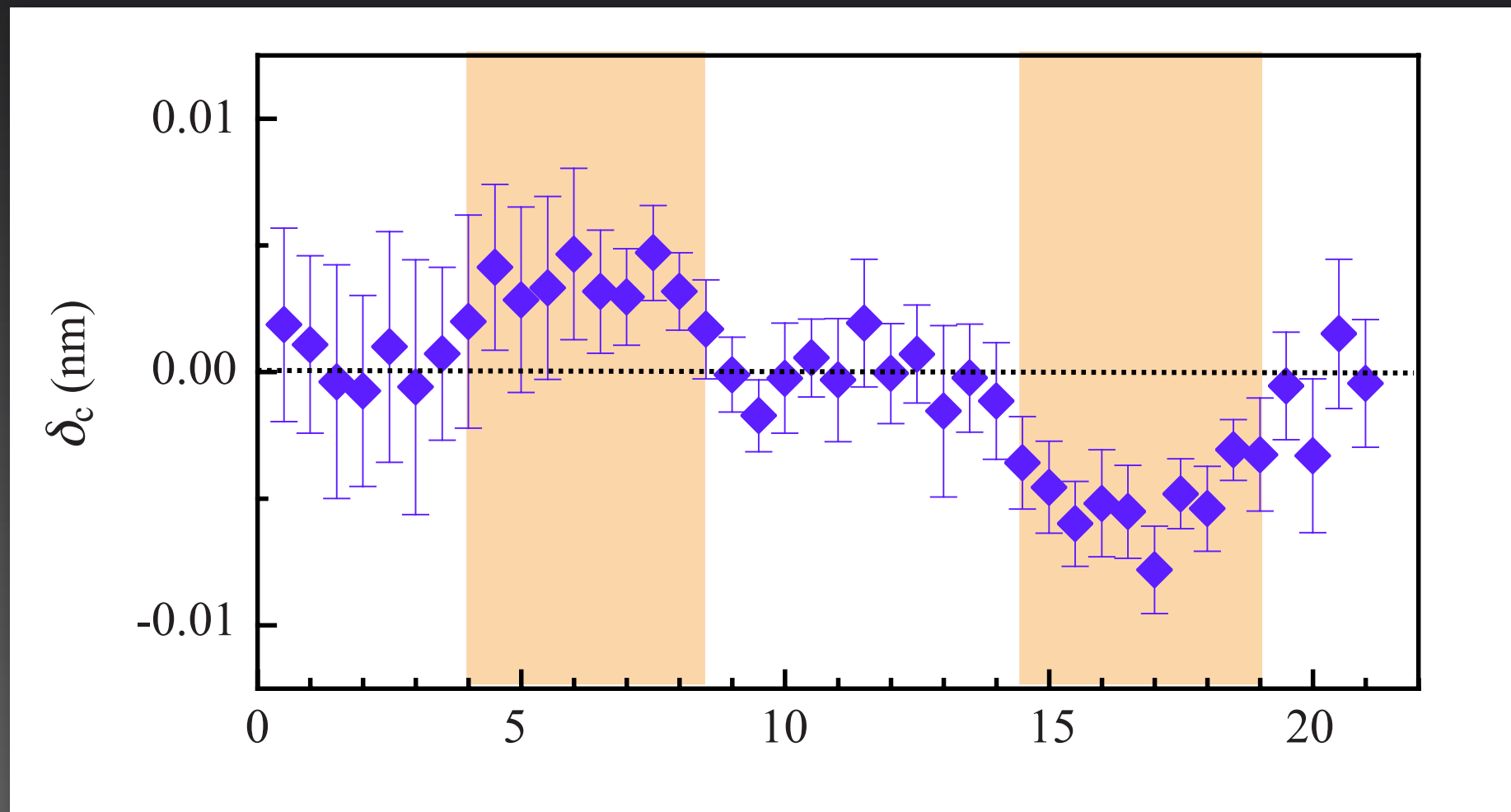




Structural distortion of
LAO propagates in several layers
of STO at the interface

C.L. Jia *et al.*, *PRB* **79**, 081405(R) (2009)

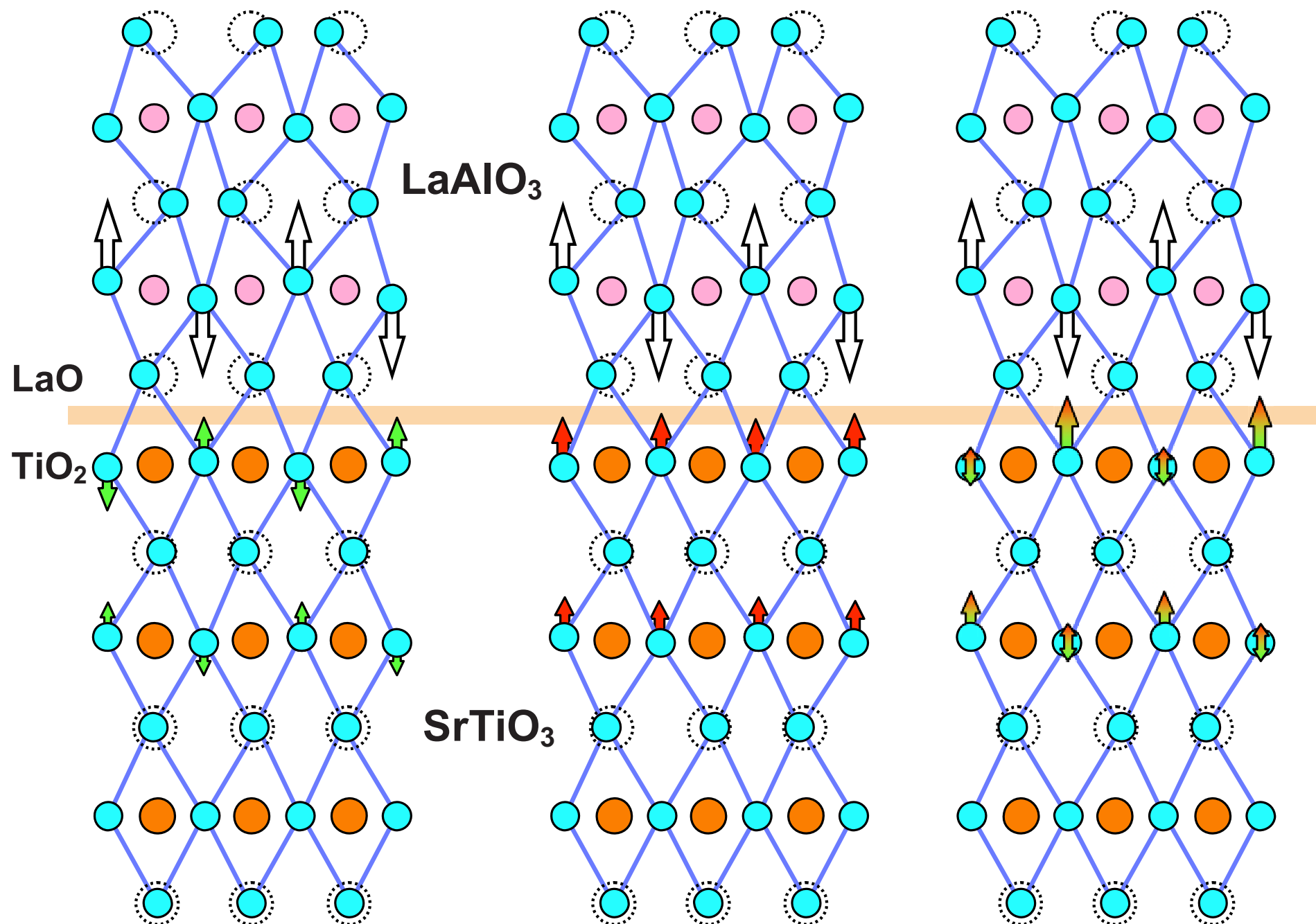
STO LAO STO



Shift of the center of the oxygen octahedron with respect to the appertaining cations

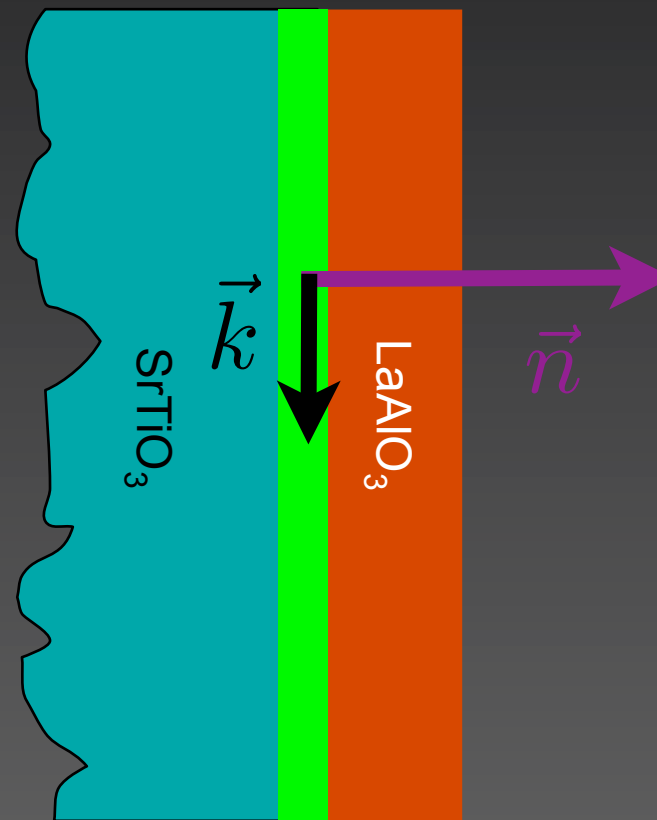
C.L. Jia *et al.*, *PRB* 79, 081405(R) (2009)

Polarisation



Combination of AFD and FE distortions

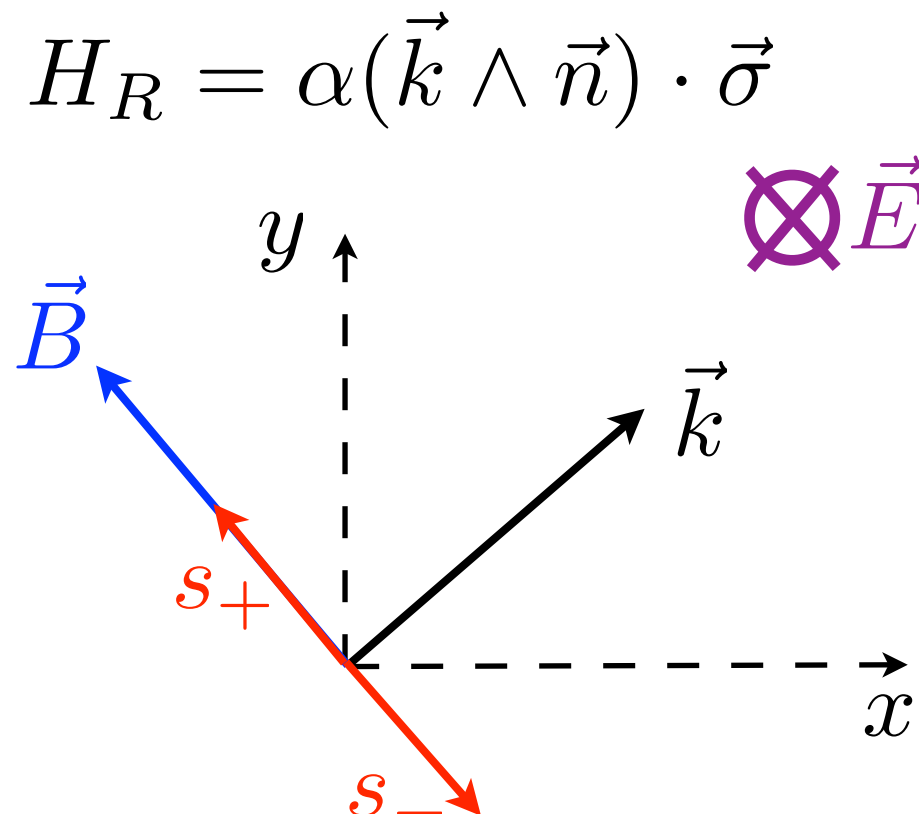
Rashba spin-orbit coupling



The electrons experience an internal magnetic field always oriented on the plane of the conduction

asymmetric
confining
potential

Rashba spin-orbit coupling



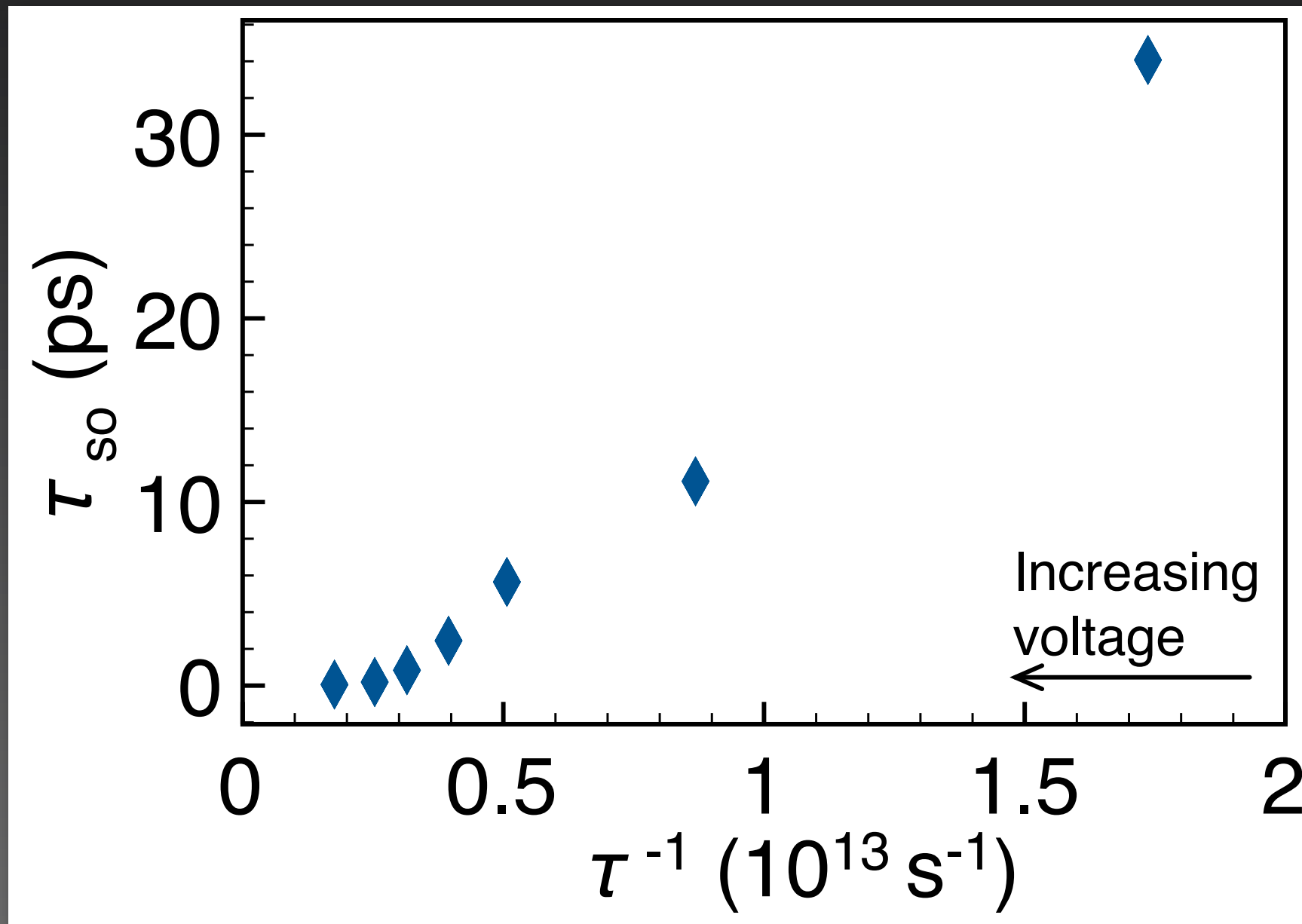
$$H_R = \alpha(\vec{k} \wedge \vec{n}) \cdot \vec{\sigma}$$

$$\tau_{\text{so}} = \hbar^4 / 4\alpha^2 m^2 2D$$

The electrons experience an internal magnetic field always oriented on the plane of the conduction

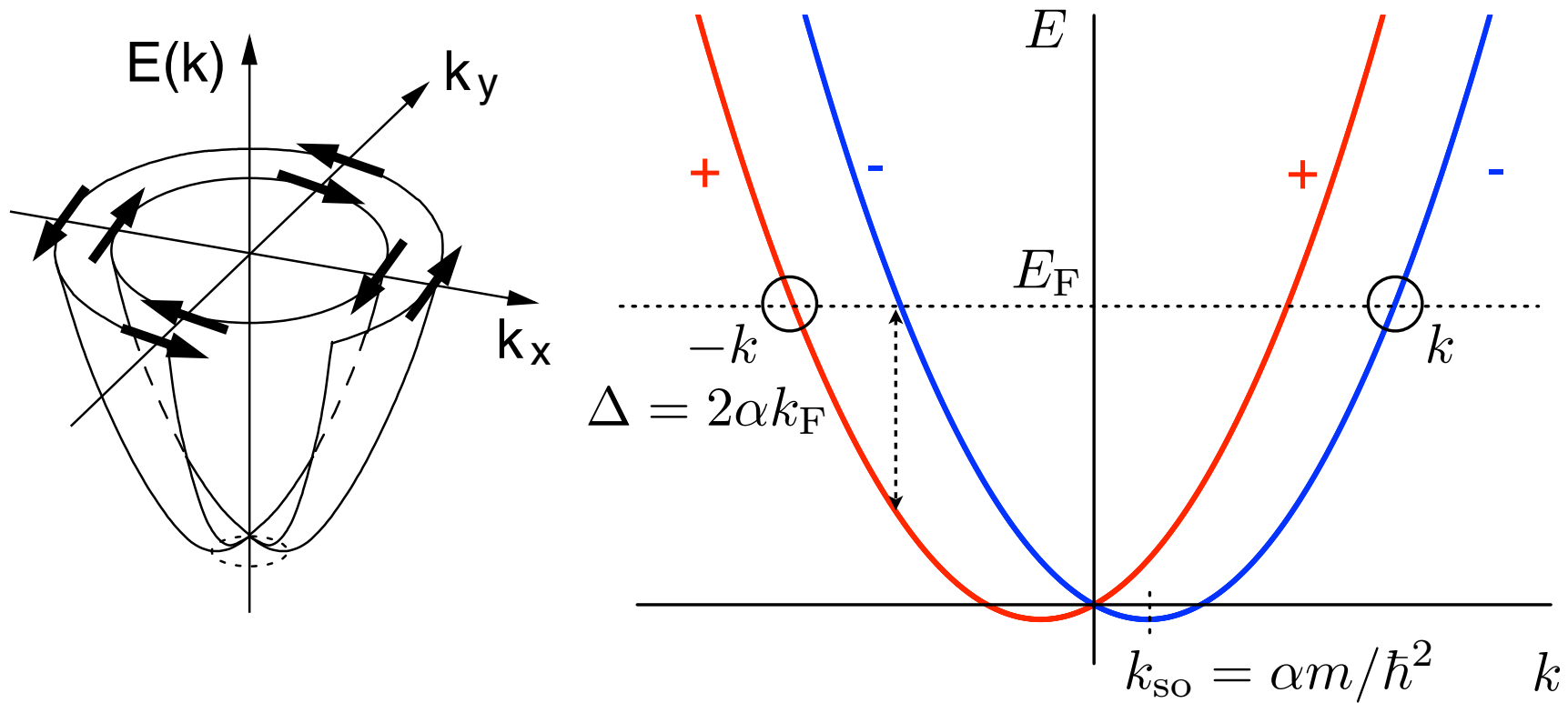
Rashba spin-orbit coupling

D'yakonov - Perel'

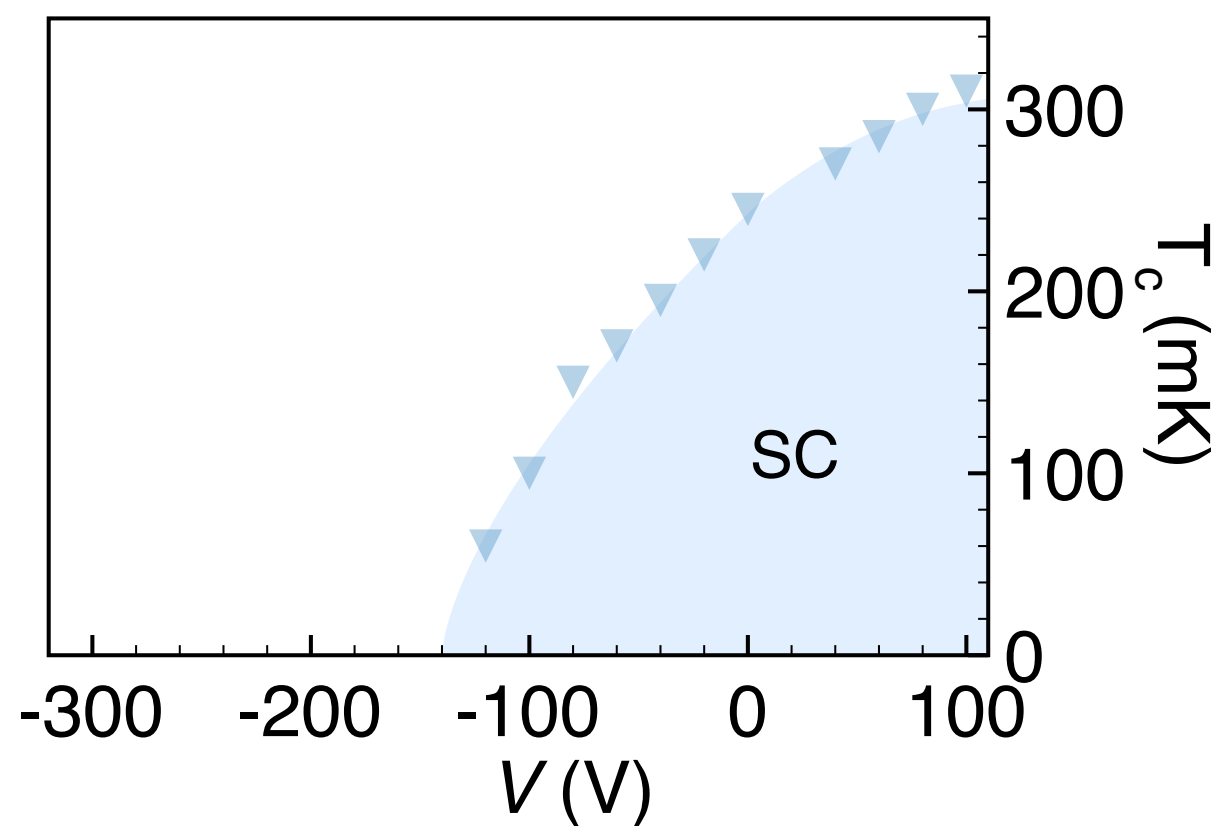
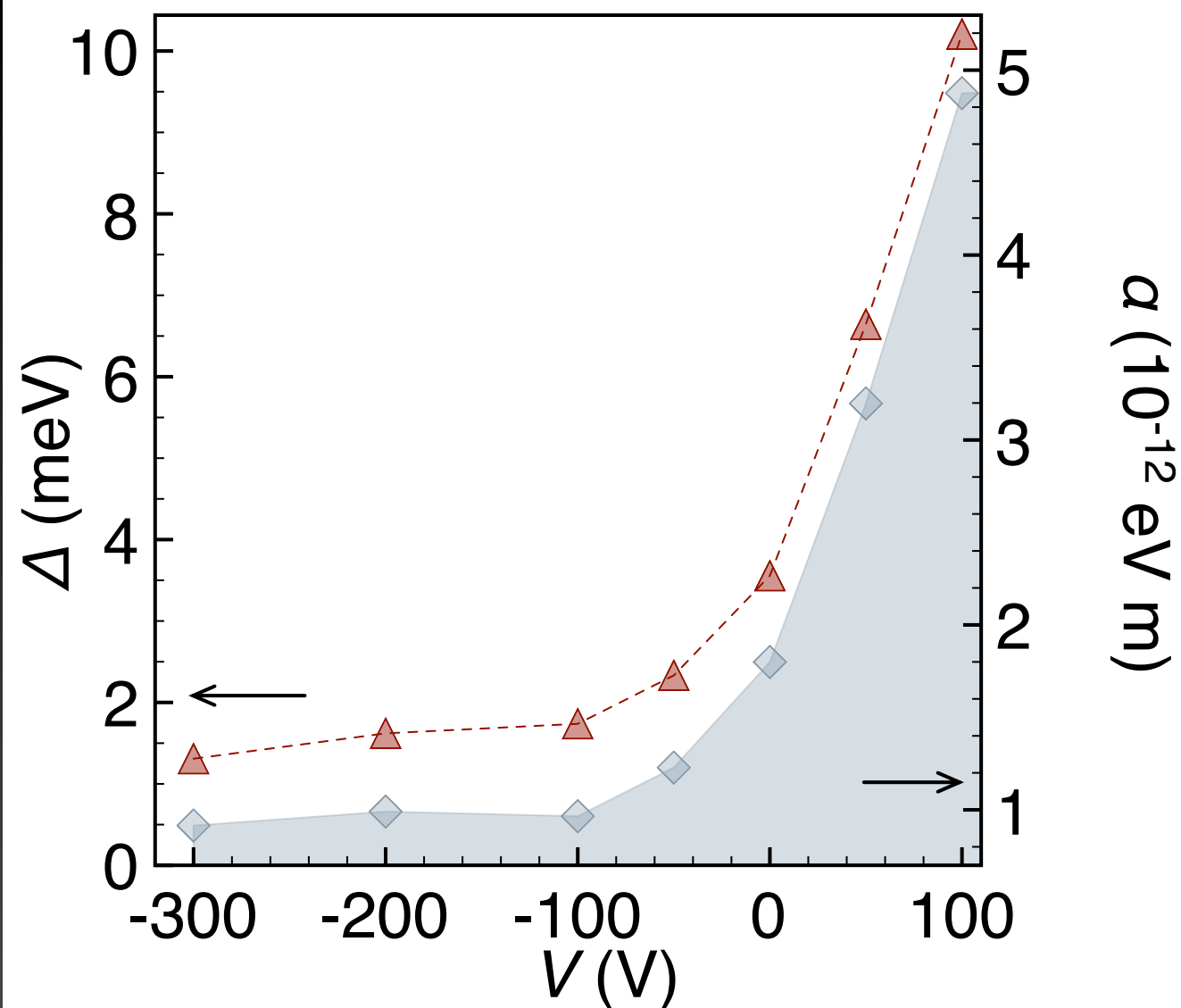


$$\tau_{so} \propto \tau^{-1}$$

Spin-splitting



$$\Delta = 2 k_F \alpha$$

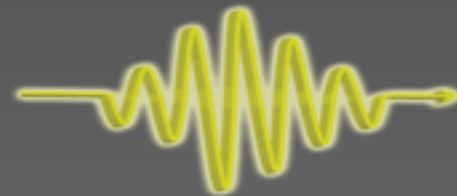


Limitations of the field effect approach

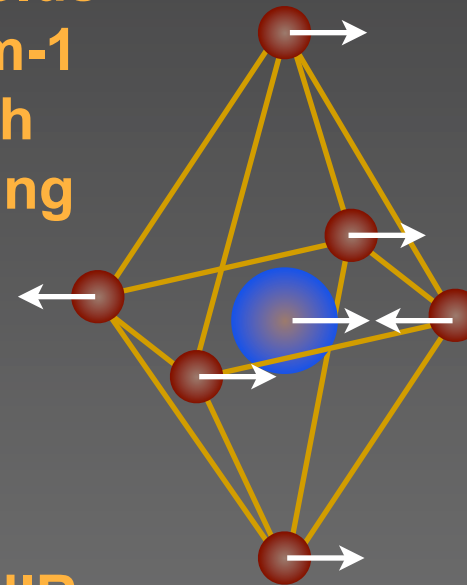
Need for a good gate dielectric
Limited to dc breakdown fields
Slow process

Ultrafast lattice excitation using light pulses

100 fs pulsed 10 MV/cm fields
16 μm wavelength, 600 cm^{-1}
modulation of bandwidth
non-linear phonon coupling

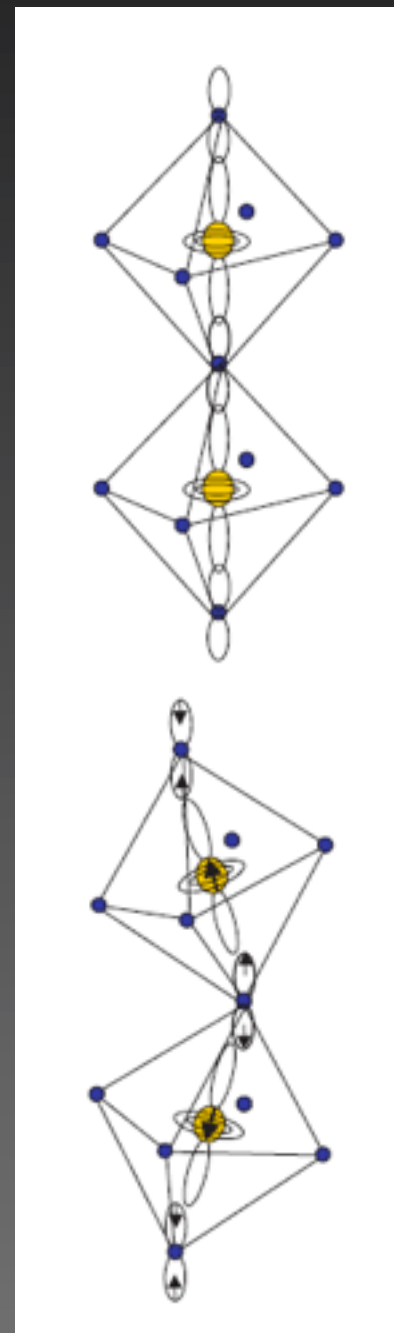
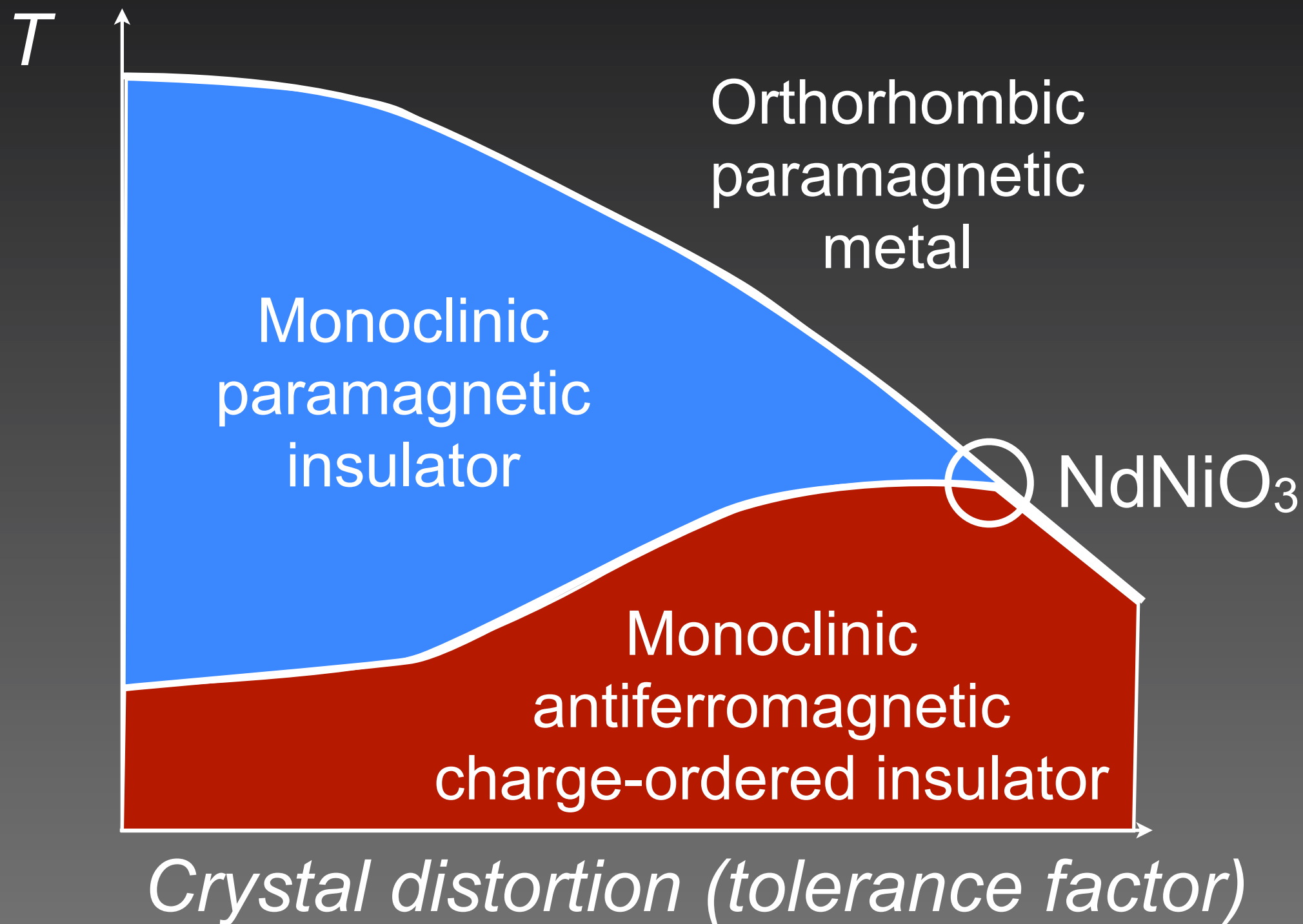


electronic properties
investigated by
pump and probe THz and NIR
spectroscopy



Nickelates phase diagram

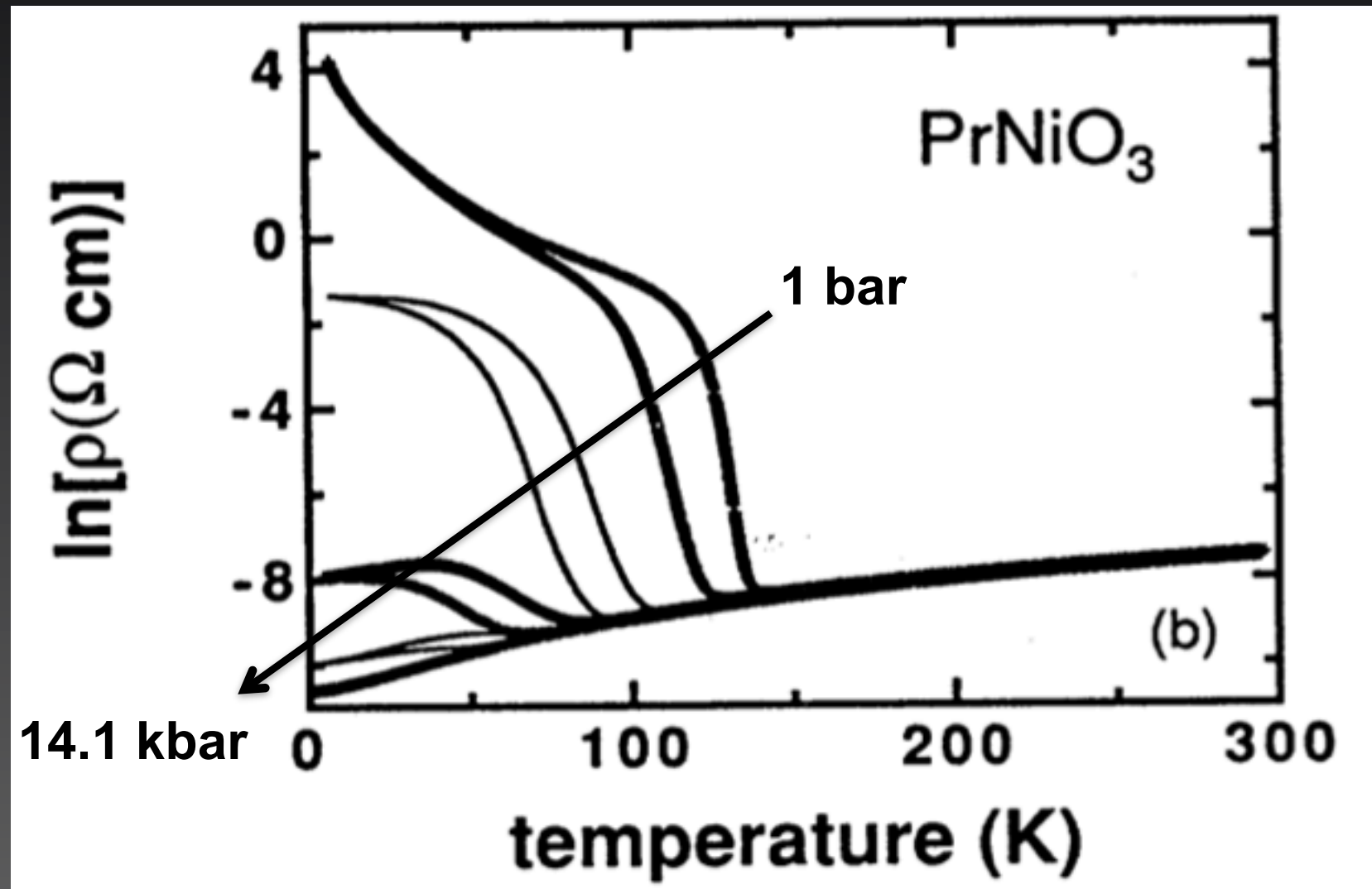
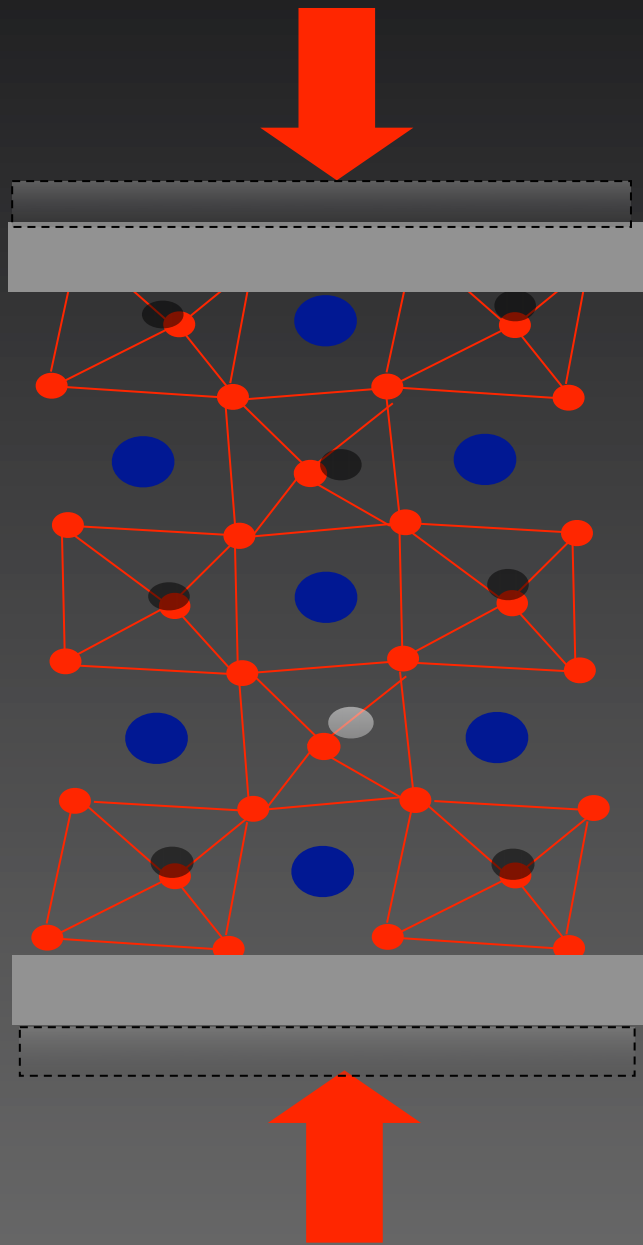
A story of coupling between electrons and lattice



M.L. Medarde, *J Phys Cond Matt* **9**, 1679 (1997)

G. Catalan, *Phase transitions* **81**, 729 (2008)

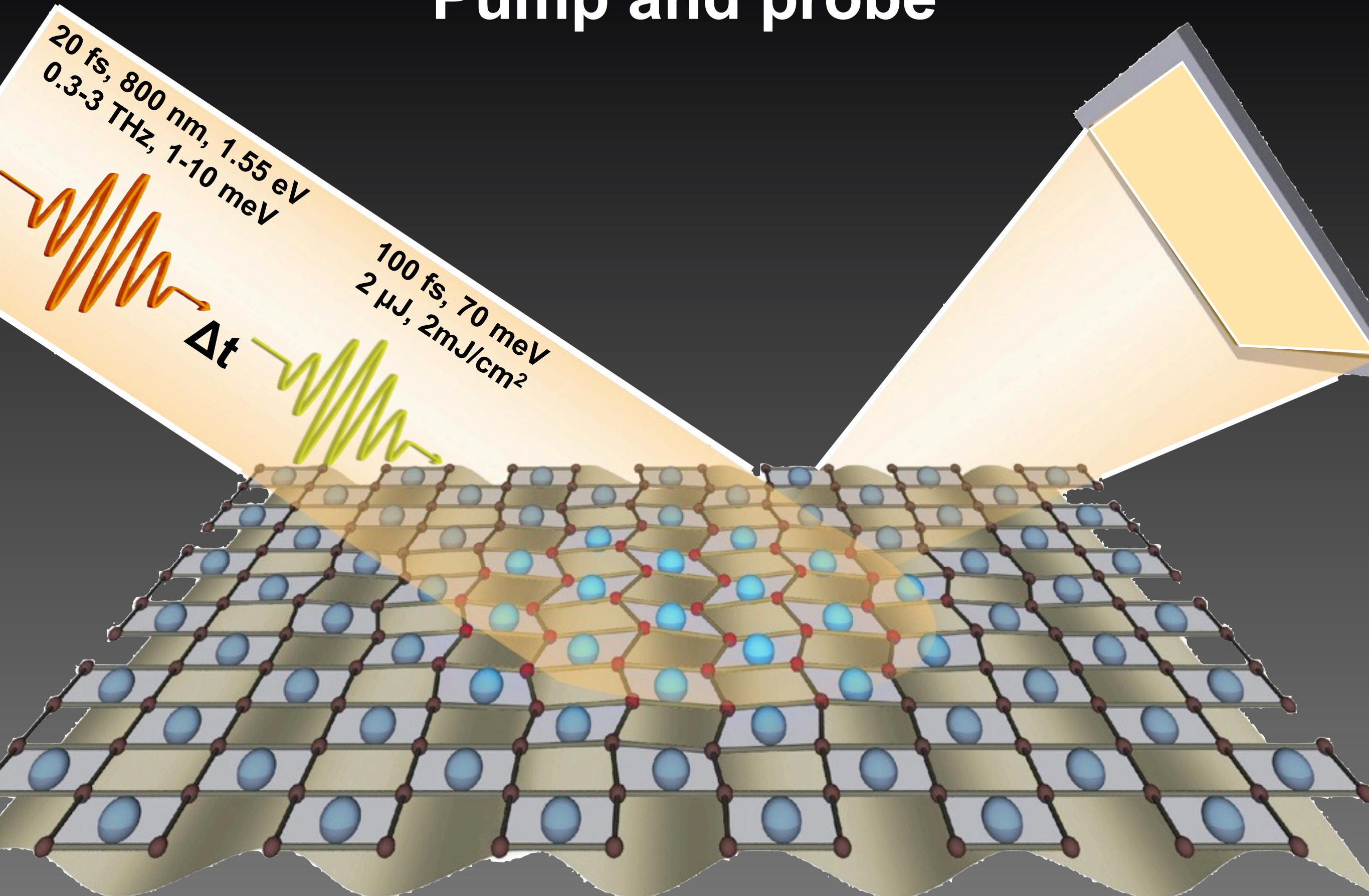
Static lattice control



P.C. Canfield et al. *PRB* 47, 12357 (1993)

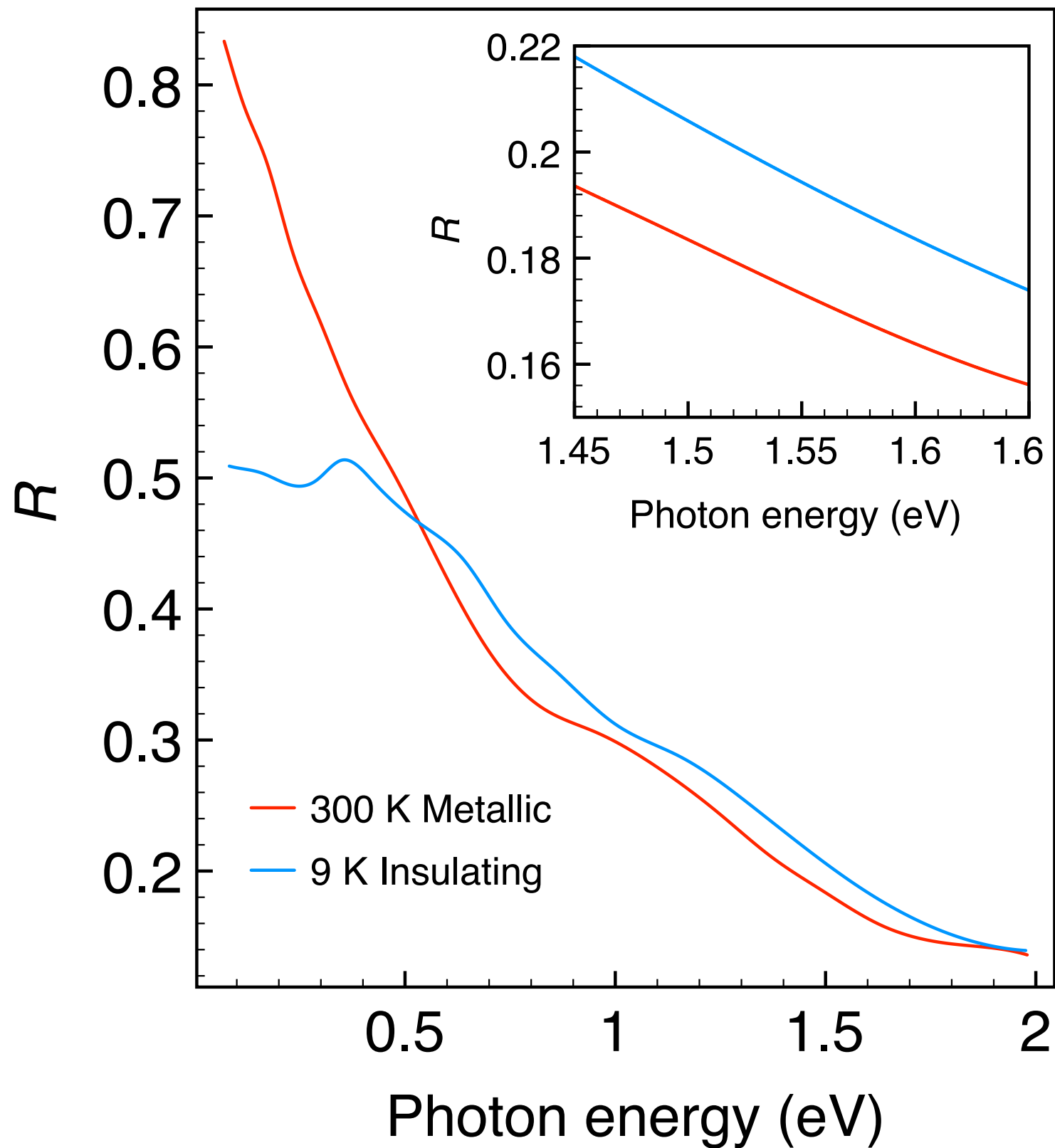
Can we drive this phase transition on the ultrafast time scale by exciting the lattice with light?

Pump and probe



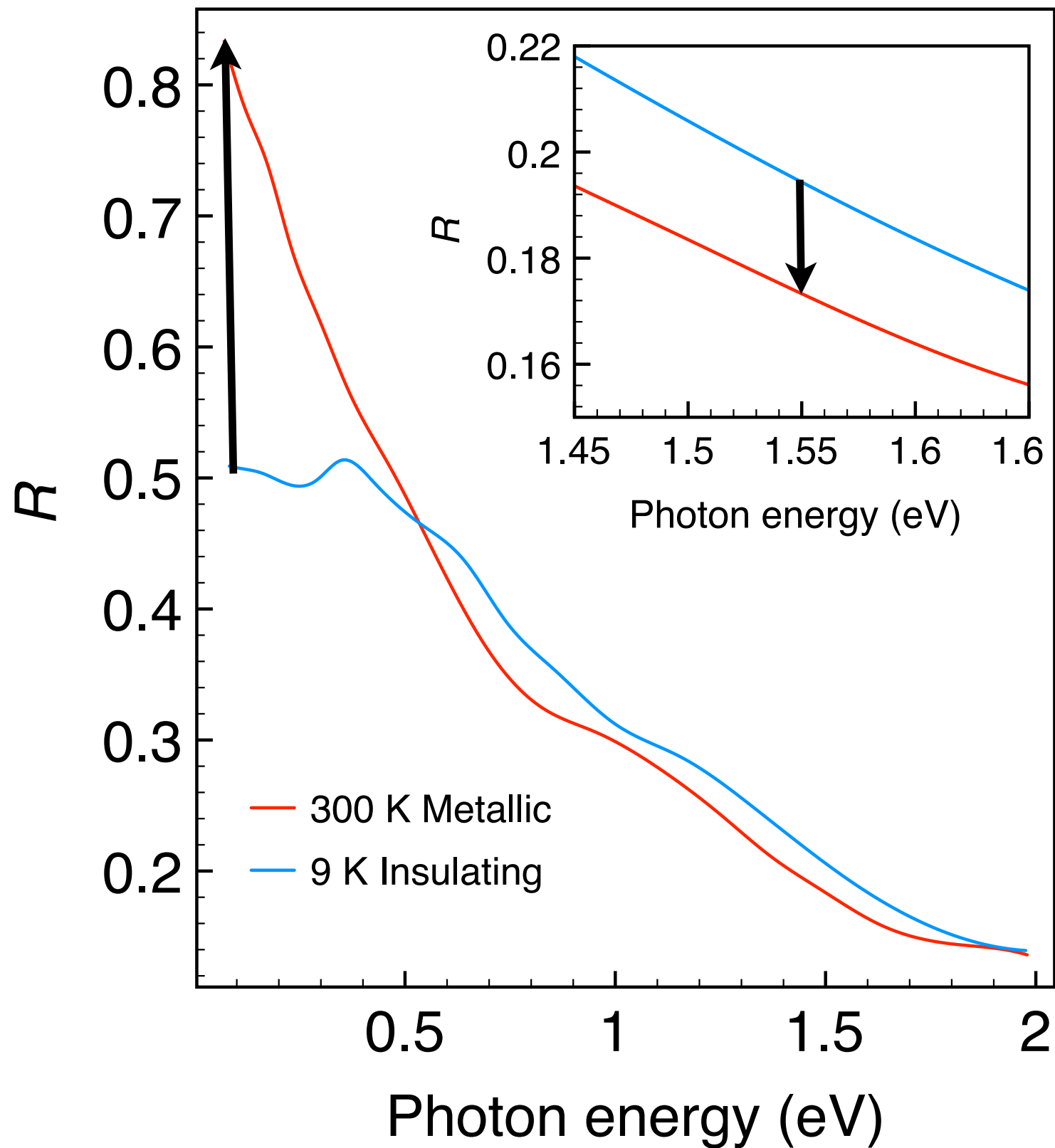
Rini et al., *Nature* 449, 72 (2007)

Static reflectivity



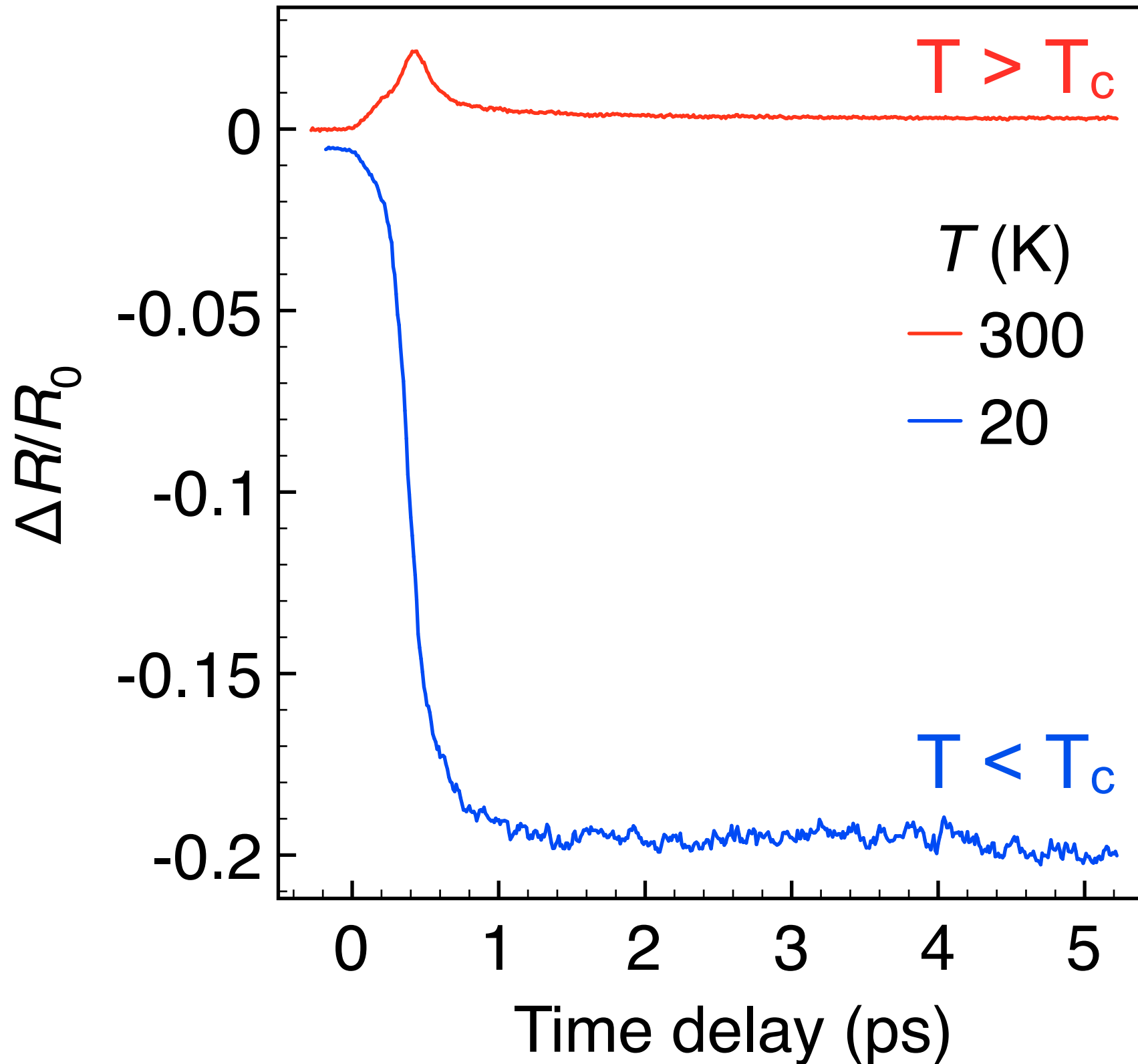
Katsufuji et al. *PRB* 51, 4830 (1995)

Static reflectivity

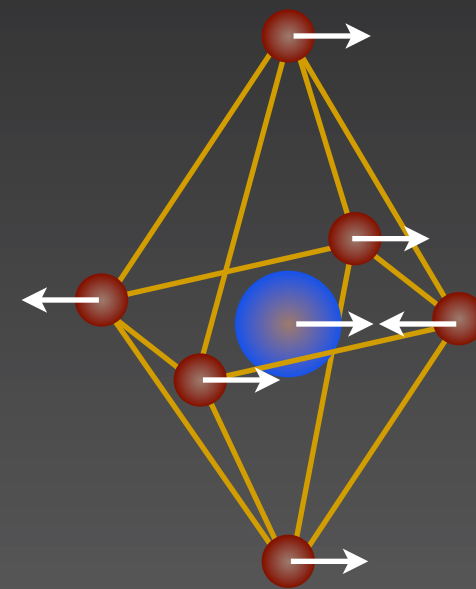
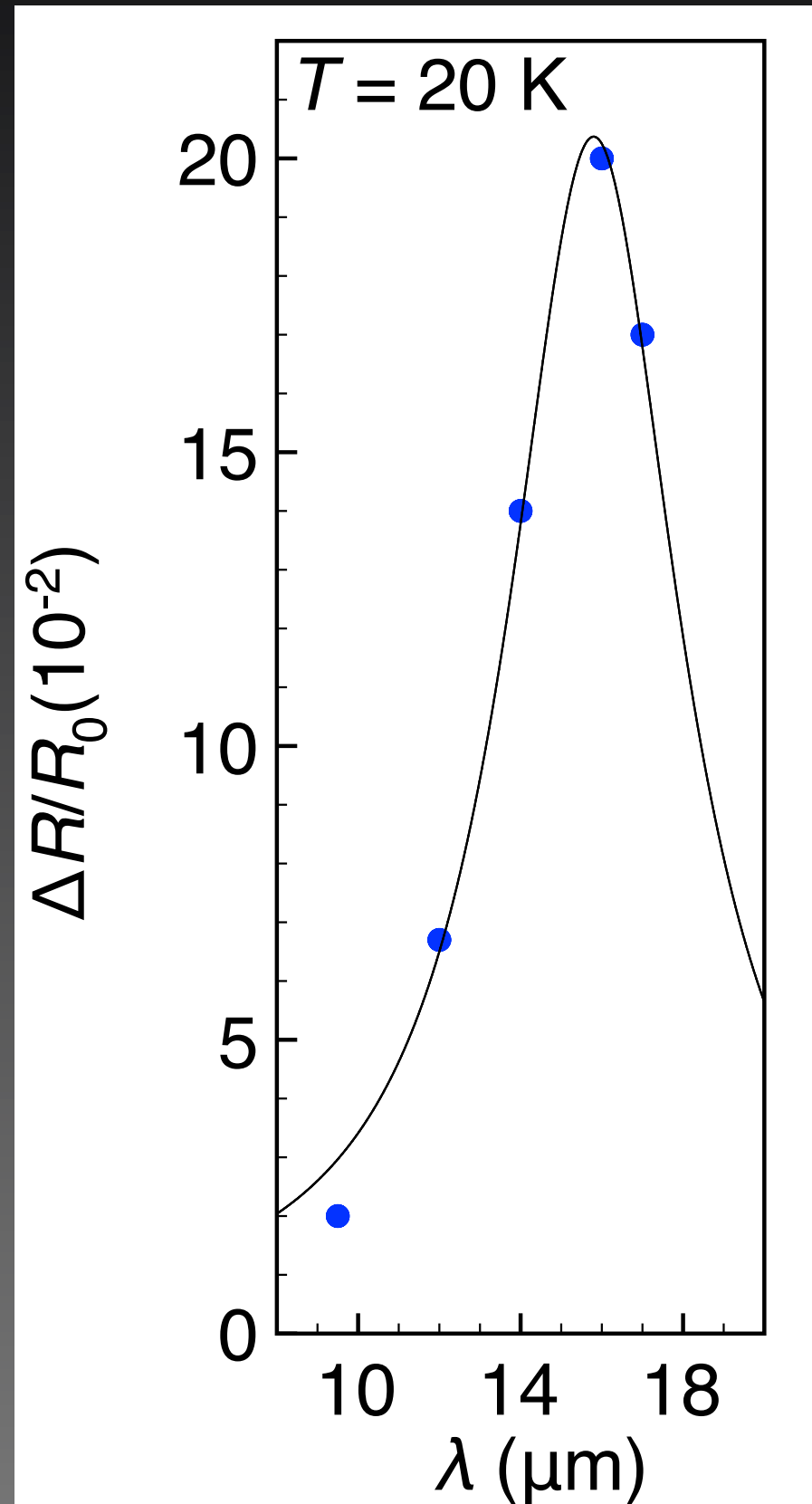


Katsufuji et al. *PRB* 51, 4830 (1995)

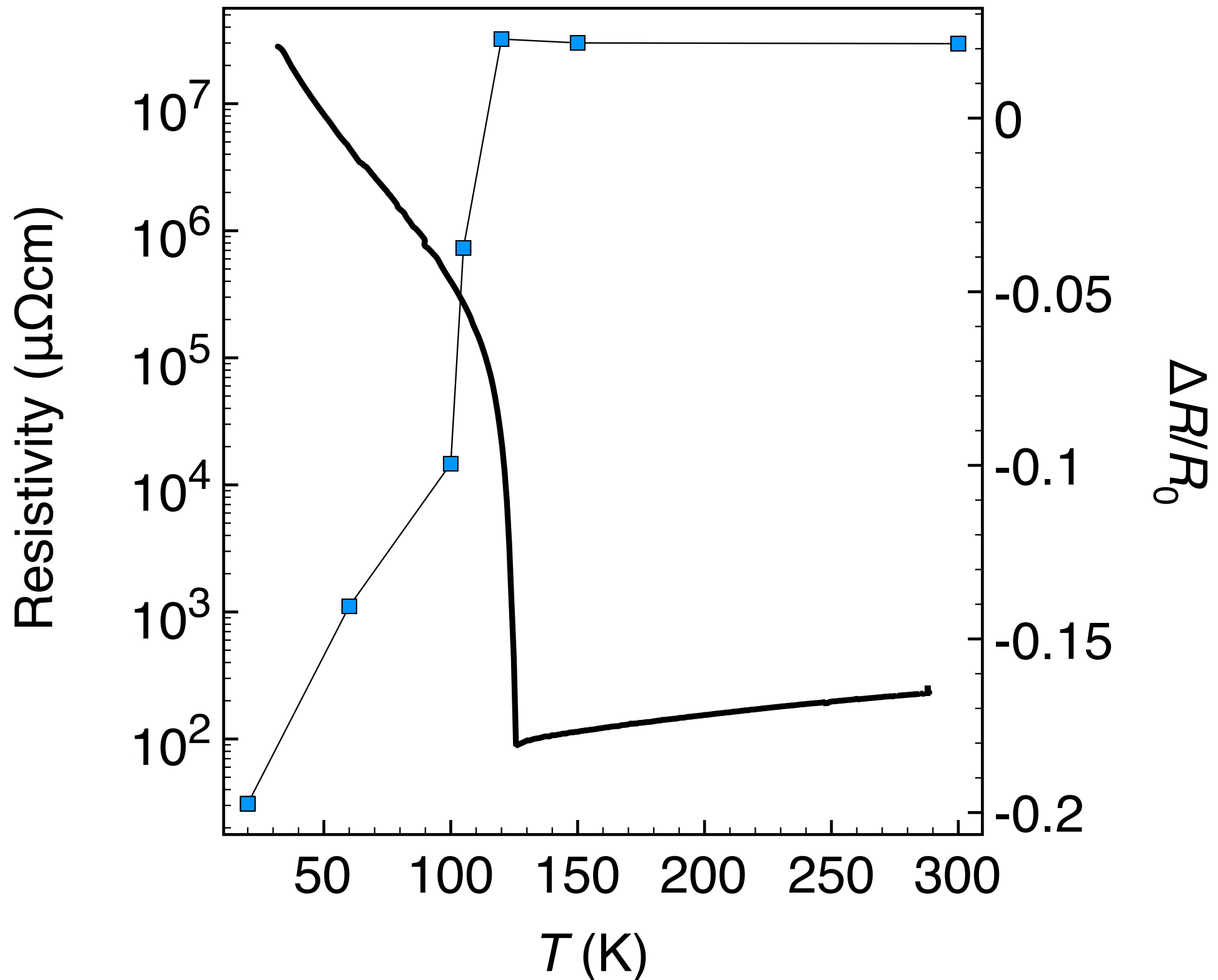
NIR probe



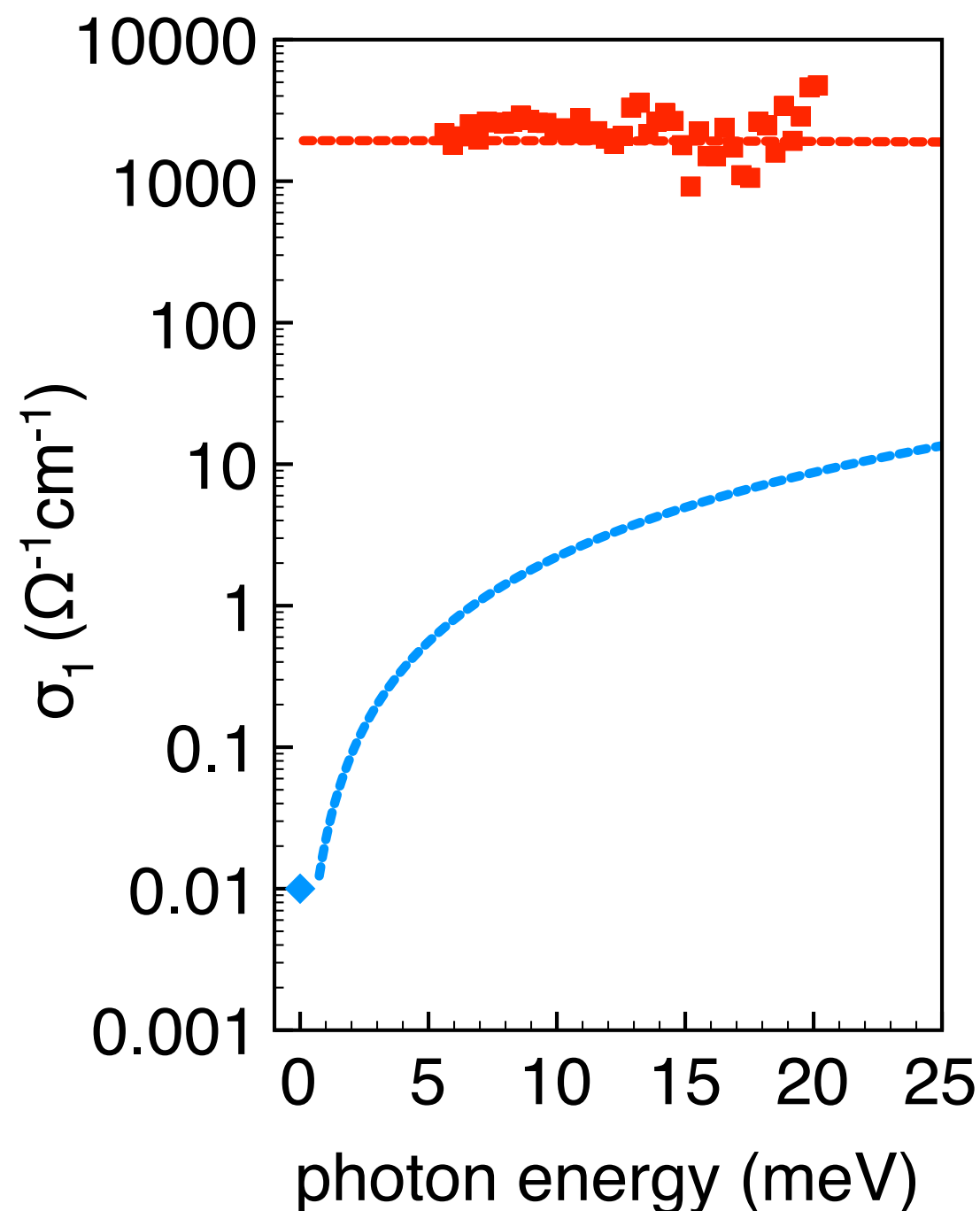
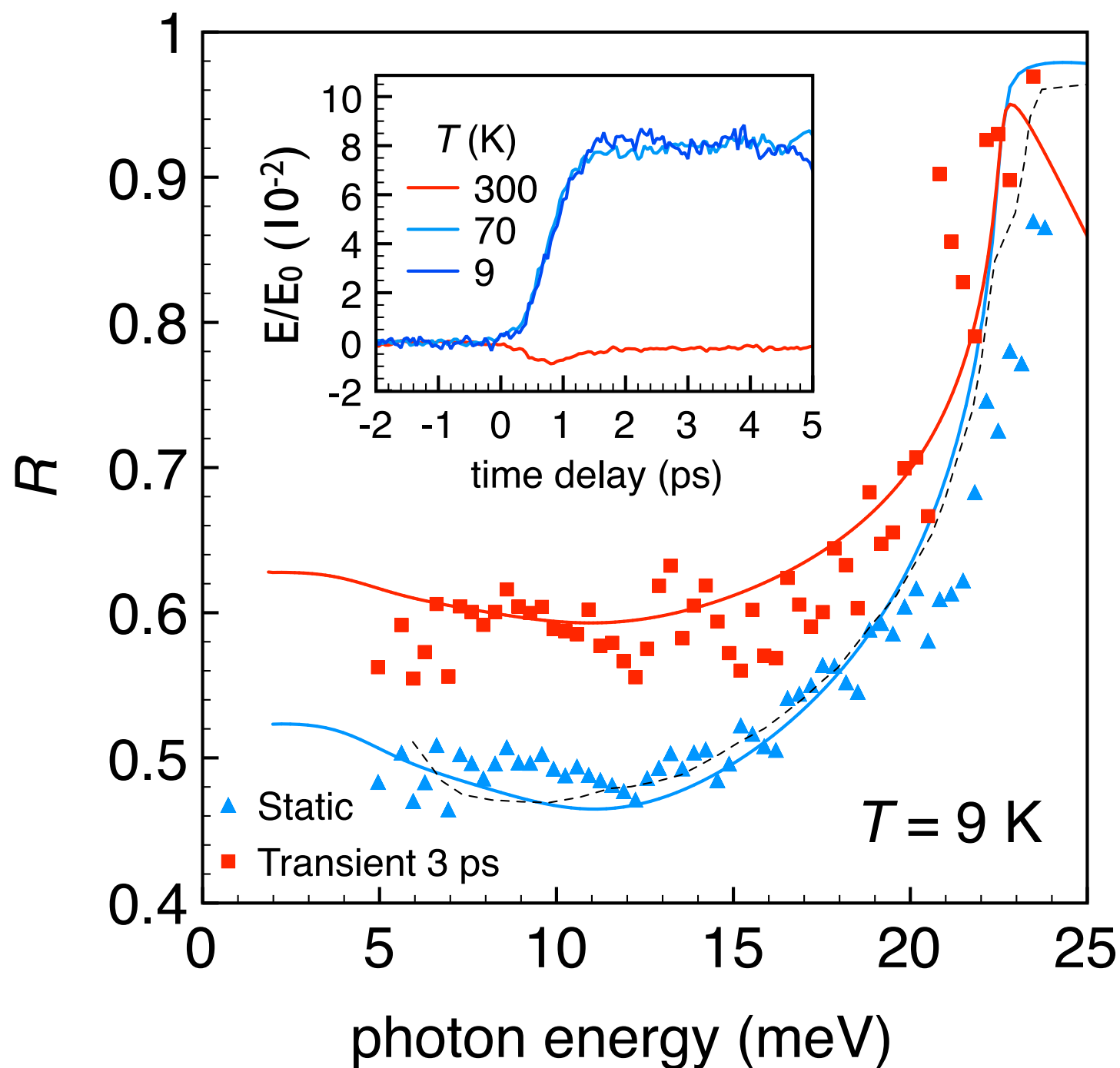
Wavelength dependence



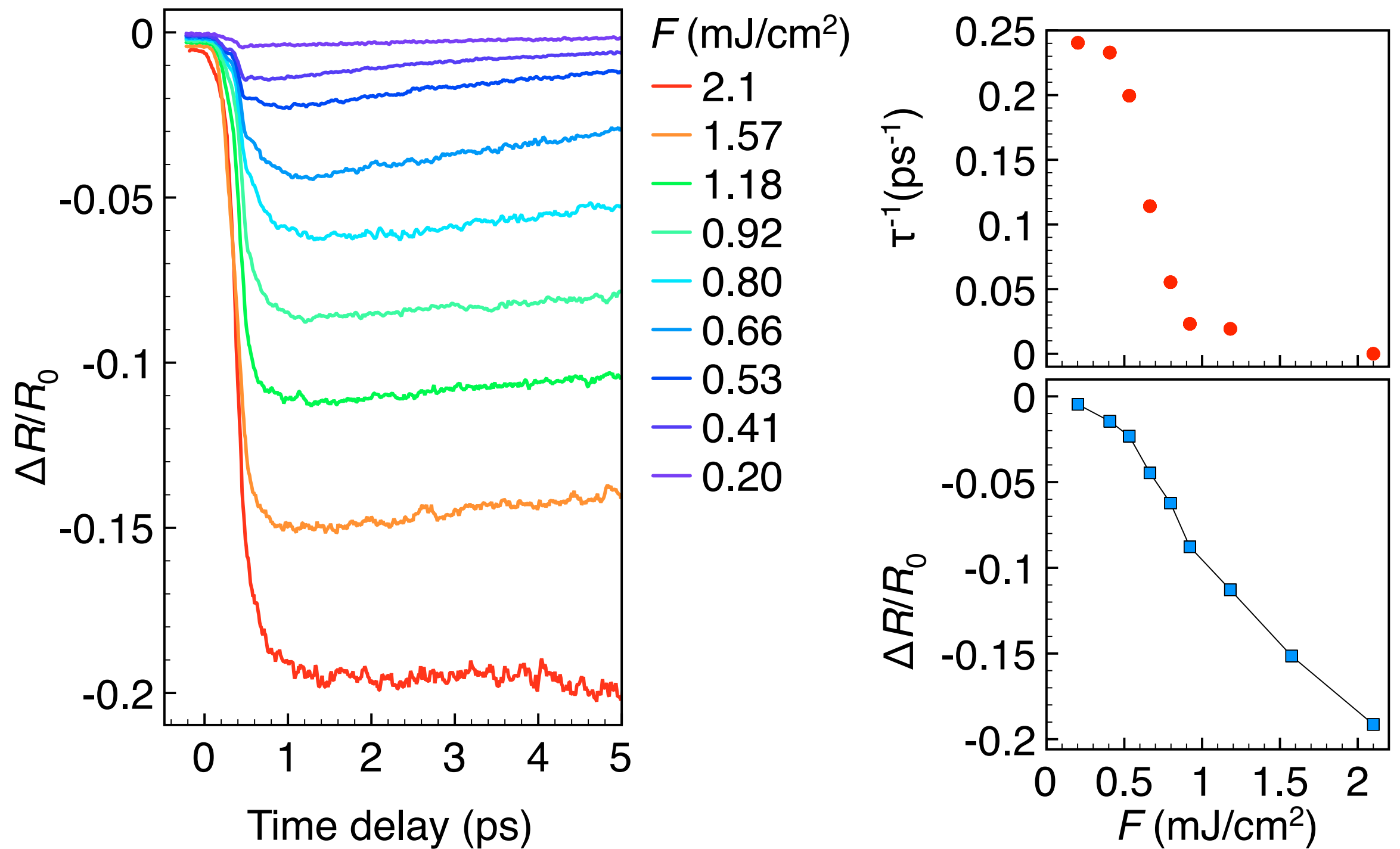
T dependence



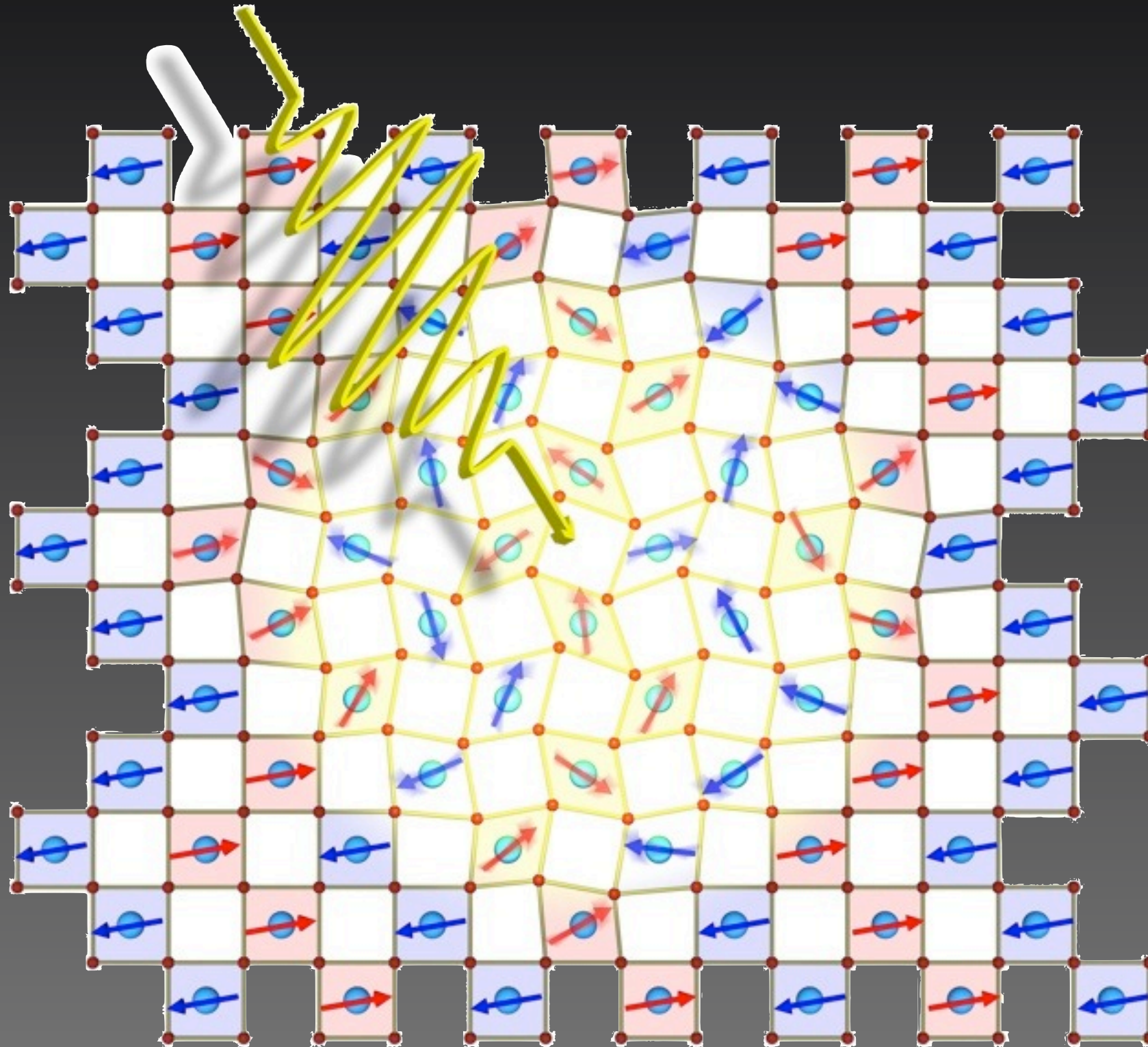
THz probe



Threshold

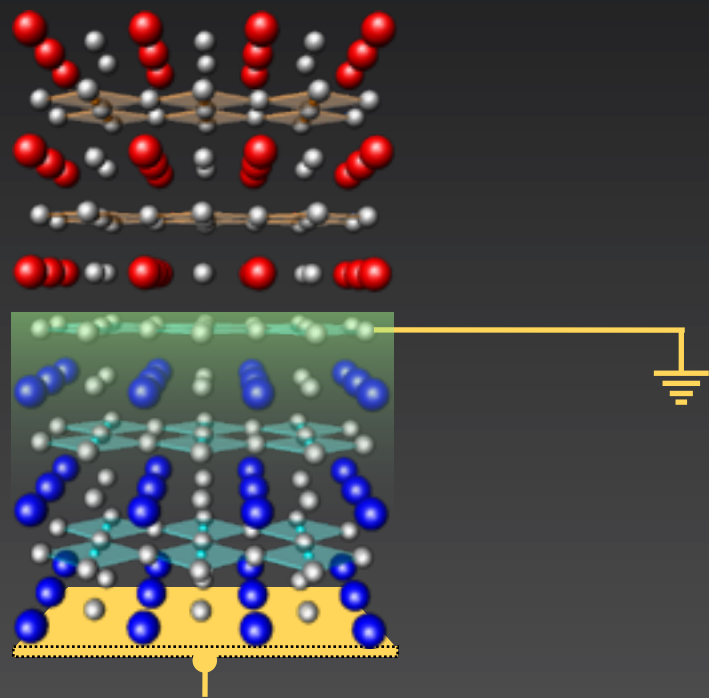


What about the electronic order?



Conclusions

Electrostatic field-effect

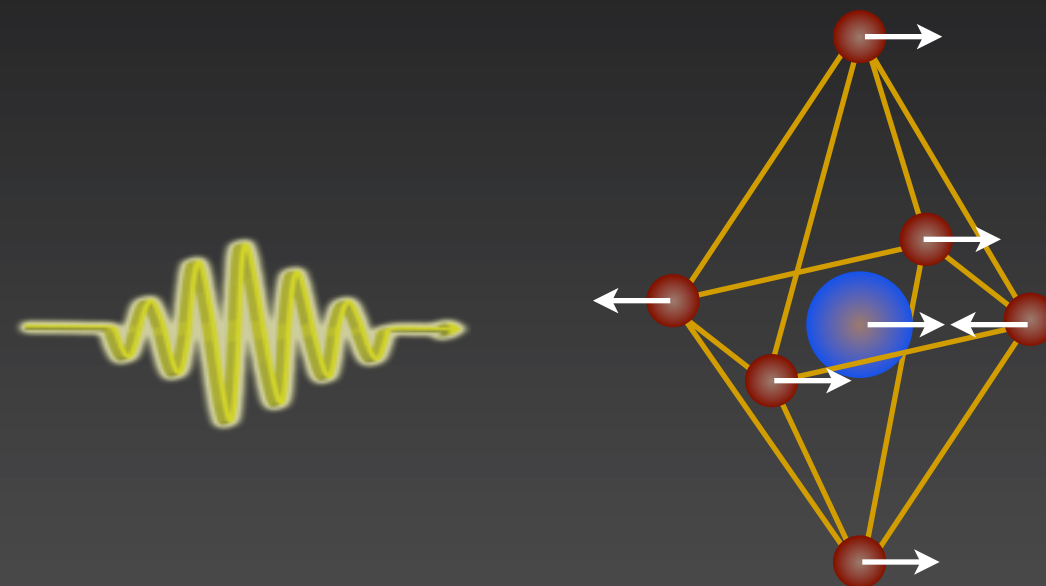


$\text{LaAlO}_3/\text{SrTiO}_3$

Superconductor - Insulator
Quantum Phase Transition

Tunable spin-orbit interaction

Ultrafast lattice excitation



$\text{NdNiO}_3/\text{LaAlO}_3$

Insulator - Metal
Non-equilibrium Phase Transition

5 orders of magnitude change in
dc conductivity

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**Financial support from the Swiss National Science Foundation
Fellowship**